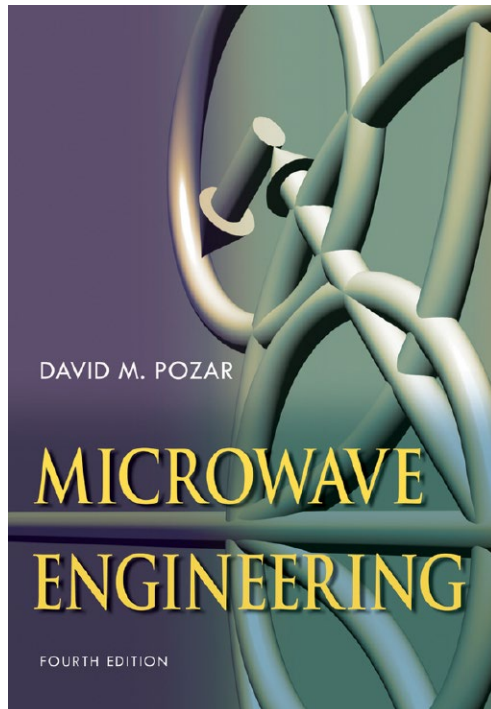


ECE 5317-6351

Microwave Engineering

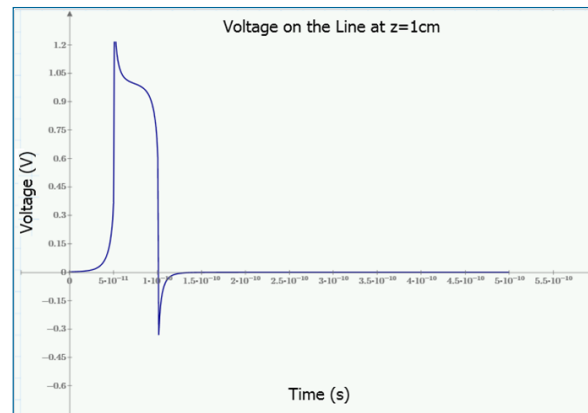
Fall 2019

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Dept. of ECE



Notes 13

Waveguiding Structures Part 8: Dispersion and Wave Velocities



Dispersion

Frequency dispersion: The phase velocity is not constant with frequency.

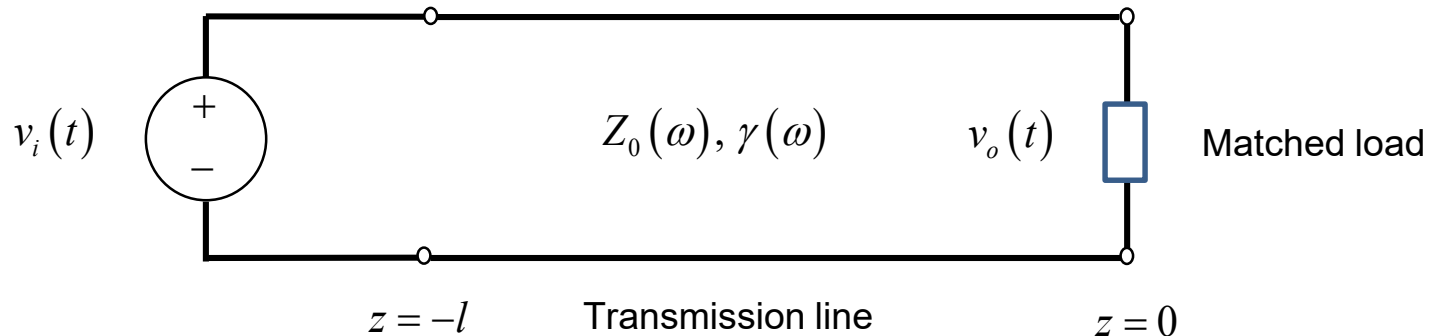
=> Different frequency components in the signal spectrum travel at different velocities.

In waveguiding structures, signal distortion is due to:

- Frequency dispersion
- Frequency-dependent attenuation
- Propagation of multiple modes that have different phase velocities

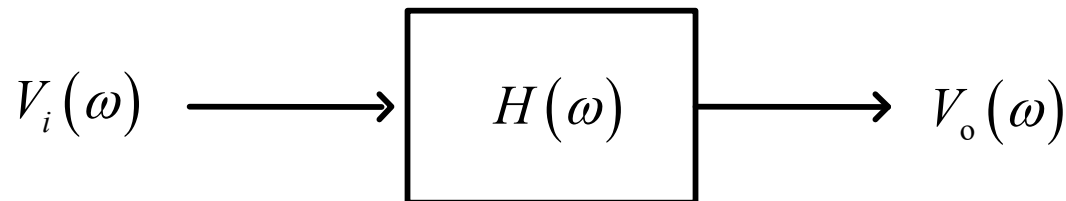
Voltage Signal on Transmission Line

Consider a signal applied to the input of a transmission line:



The transmission line as a transfer function in the phasor domain:

$$V_o(\omega) = H(\omega)V_i(\omega)$$



$$H(\omega) = e^{-\gamma(\omega)l} = A(\omega)e^{j\phi(\omega)} = e^{-\alpha(\omega)l} e^{-j\beta(\omega)l}$$

$$A(\omega) = e^{-\alpha(\omega)l}$$

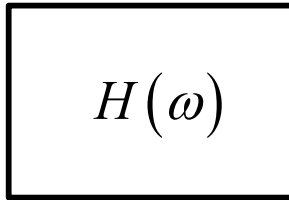
$$\phi(\omega) = -\beta(\omega)l$$

Signal Propagation

General System

Input signal

$$s_i(t)$$



Output signal

$$s_o(t)$$

$$H(\omega) \equiv \frac{S_o(\omega)}{S_i(\omega)}$$

$$\tilde{s}_i(\omega) = \int_{-\infty}^{\infty} s_i(t) e^{-j\omega t} dt$$
$$s_i(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega$$

Fourier transform pair

Property of real-valued signal:

$$\tilde{s}_i(-\omega) = \tilde{s}_i^*(\omega)$$

Note:
The signal will normally be assumed to represent the voltage.

(Please see the proof on the next slide.)

Signal Propagation (cont.)

Proof of Property:

$$\tilde{s}_i(-\omega) = \tilde{s}_i^*(\omega)$$

$$\tilde{s}_i(\omega) = \int_{-\infty}^{\infty} s_i(t) e^{-j\omega t} dt$$
$$s_i(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega$$

Real signal: $s_i(t) = s_i^*(t)$

$$\Rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} (\tilde{s}_i(\omega))^* e^{-j\omega t} d\omega$$

$$\Rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(-\omega') e^{-j\omega' t} d\omega' = \frac{1}{2\pi} \int_{-\infty}^{\infty} (\tilde{s}_i(\omega))^* e^{-j\omega t} d\omega \quad (\omega = -\omega')$$

$$\Rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(-\omega) e^{-j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} (\tilde{s}_i(\omega))^* e^{-j\omega t} d\omega \quad (\omega' \rightarrow \omega)$$

$$\Rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(-\omega) e^{j\omega t'} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} (\tilde{s}_i(\omega))^* e^{j\omega t'} d\omega \quad (t = -t')$$

$$\Rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(-\omega) e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} (\tilde{s}_i(\omega))^* e^{j\omega t} d\omega \quad (t' \rightarrow t)$$

$$\Rightarrow F^{-1} \{ \tilde{s}_i(-\omega) \} = F^{-1} \left\{ (\tilde{S}_i(\omega))^* \right\}$$

$$\Rightarrow \tilde{s}_i(-\omega) = (\tilde{s}_i(\omega))^*$$

Signal Propagation (cont.)

We can then show:

$$s_i(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega \quad \rightarrow \quad s_i(t) = \frac{1}{\pi} \operatorname{Re} \int_0^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega$$

(Please see the derivation on the next slide.)

The form on the right is convenient, since it only involves positive values of ω . (In this case, ω has the nice interpretation of being radian frequency: $\omega = 2\pi f$.)

Signal Propagation (cont.)

Derivation:

$$\begin{aligned} s_i(t) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega \\ &= \frac{1}{2\pi} \int_{-\infty}^0 \tilde{s}_i(\omega) e^{j\omega t} d\omega + \frac{1}{2\pi} \int_0^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega \quad (\text{split integral into two parts}) \\ &= \frac{1}{2\pi} (-1) \int_{\infty}^0 \tilde{s}_i(-\omega') e^{-j\omega' t} d\omega' + \frac{1}{2\pi} \int_0^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega \quad (\omega = -\omega' \text{ in first integral}) \\ &= \frac{1}{2\pi} \int_0^{\infty} \tilde{s}_i(-\omega') e^{-j\omega' t} d\omega' + \frac{1}{2\pi} \int_0^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega \quad (\text{reverse limits in first integral}) \\ &= \frac{1}{2\pi} \int_0^{\infty} \tilde{s}_i(-\omega) e^{-j\omega t} d\omega + \frac{1}{2\pi} \int_0^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega \quad (\omega' \rightarrow \omega \text{ in first integral}) \\ &= \frac{1}{2\pi} \int_0^{\infty} (\tilde{s}_i(\omega))^* e^{-j\omega t} d\omega + \frac{1}{2\pi} \int_0^{\infty} \tilde{s}_i(\omega) e^{j\omega t} d\omega \quad (\text{property of real - valued signal}) \\ &= \frac{1}{2\pi} \left[\int_0^{\infty} (\tilde{s}_i(\omega) e^{j\omega t})^* + (\tilde{s}_i(\omega) e^{j\omega t}) \right] d\omega \quad (\text{conjugate property}) \\ &= \frac{1}{\pi} \int_0^{\infty} \text{Re} \{ \tilde{s}_i(\omega) e^{j\omega t} \} d\omega \quad (\text{add the two integrands together}) \end{aligned}$$

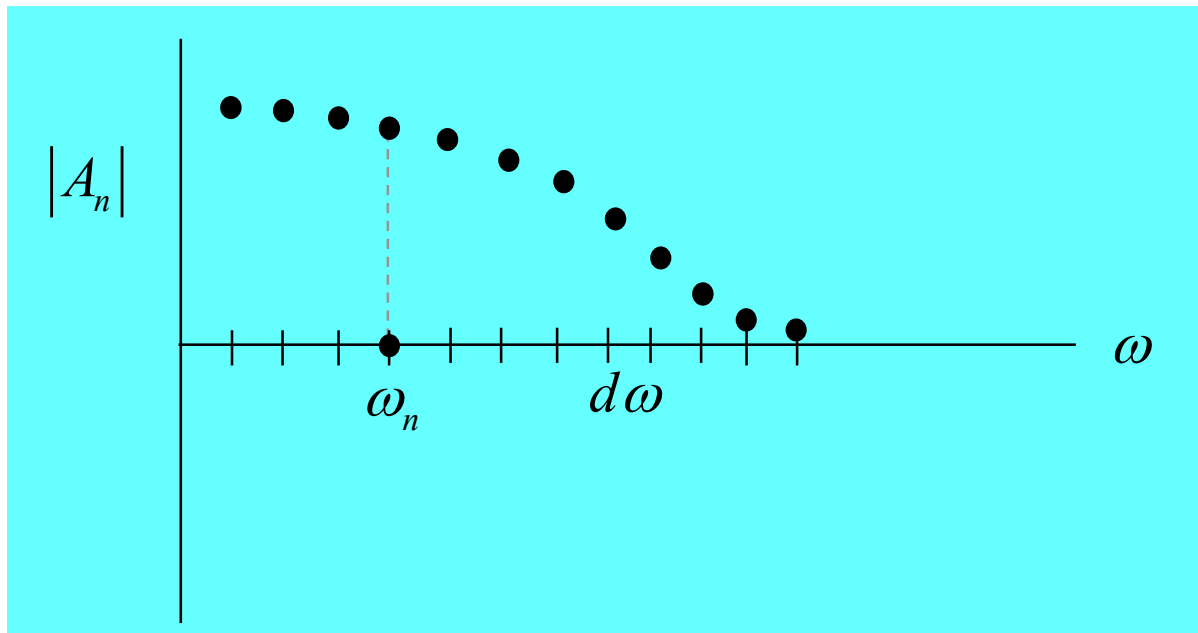
Signal Propagation (cont.)

Hence, we have

$$s_i(t) = \int_0^{\infty} \operatorname{Re} \left\{ \left(\frac{1}{\pi} \tilde{s}_i(\omega) d\omega \right) e^{j\omega t} \right\}$$

Interpreted as a **phasor**.

$$A_n = \text{complex amplitude of phasor } \#n = \frac{1}{\pi} \tilde{s}_i(\omega) d\omega$$



Signal Propagation (cont.)

$$s_i(t) = \int_0^{\infty} \operatorname{Re} \left\{ \left(\frac{1}{\pi} \tilde{s}_i(\omega) d\omega \right) e^{j\omega t} \right\}$$

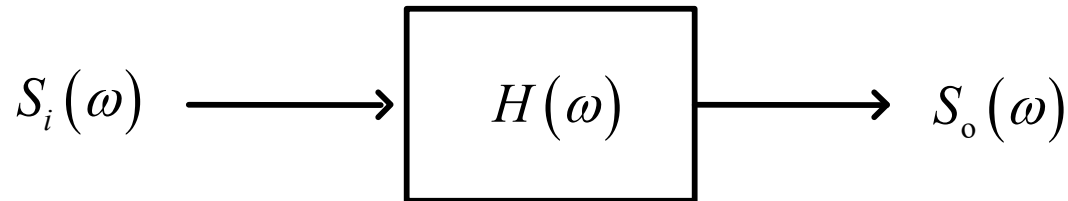
Using the transfer function, we have:

$$\begin{aligned} s_o(t) &= \int_0^{\infty} \operatorname{Re} \left\{ H(\omega) \left(\frac{1}{\pi} \tilde{s}_i(\omega) d\omega \right) e^{j\omega t} \right\} \\ &= \frac{1}{\pi} \operatorname{Re} \int_0^{\infty} H(\omega) \tilde{s}_i(\omega) e^{j\omega t} d\omega \end{aligned}$$

Signal Propagation (cont.)

Summary

$$s_o(t) = \frac{1}{\pi} \operatorname{Re} \int_0^{\infty} H(\omega) \tilde{s}_i(\omega) e^{j\omega t} d\omega$$



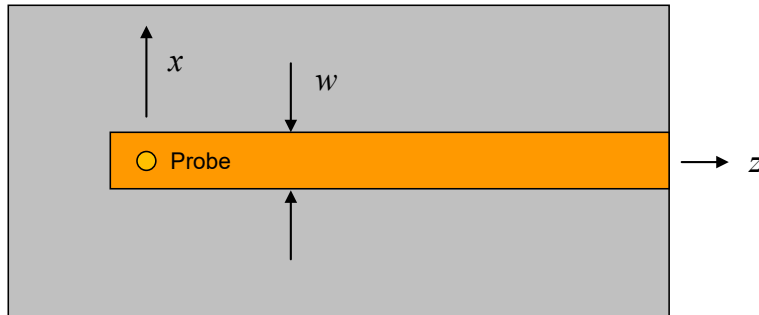
Phasor domain

Waveguiding system:

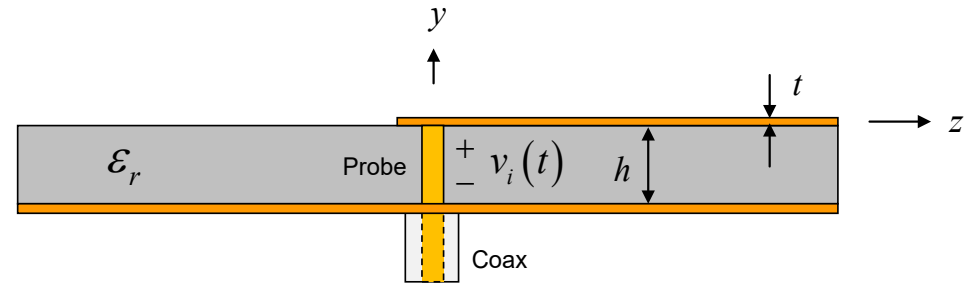
$$H(\omega) = A(\omega) e^{j\phi(\omega)} = e^{-\alpha(\omega)l} e^{-j\beta(\omega)l}$$

Signal Propagation (cont.)

Propagation on a microstrip line



TOP VIEW



SIDE VIEW

$$\epsilon_r = 2.33$$

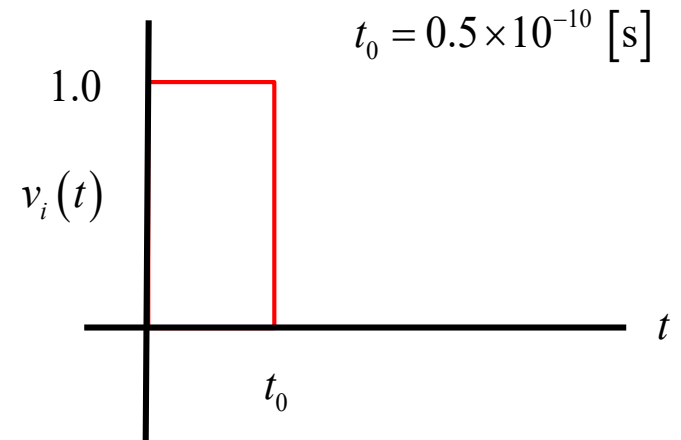
$$\tan \delta = 0.001$$

$$h = 0.787 \text{ [mm]} \text{ (31 mils)}$$

$$w = 2.35 \text{ [mm]}$$

$$t = 0.0175 \text{ [mm]} \text{ ("half oz" copper cladding)}$$

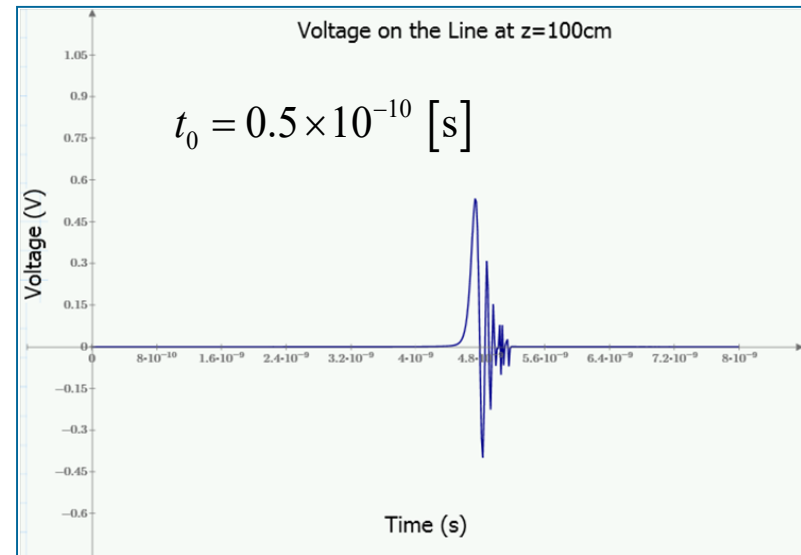
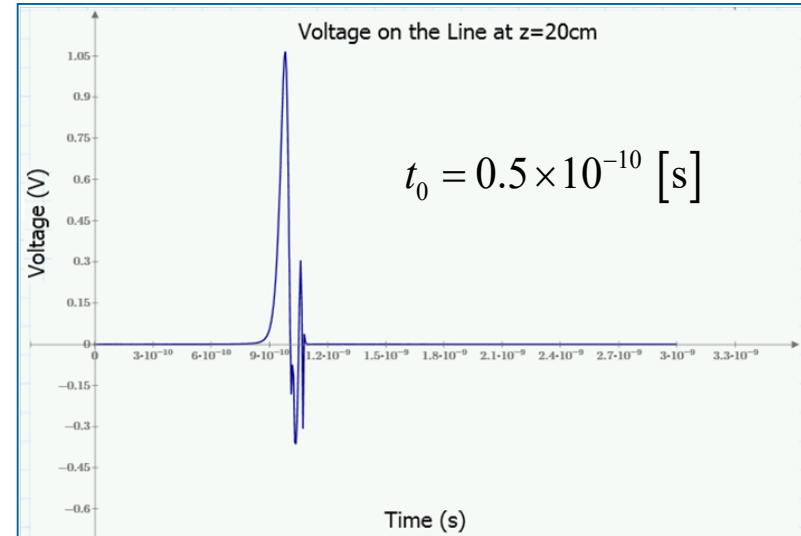
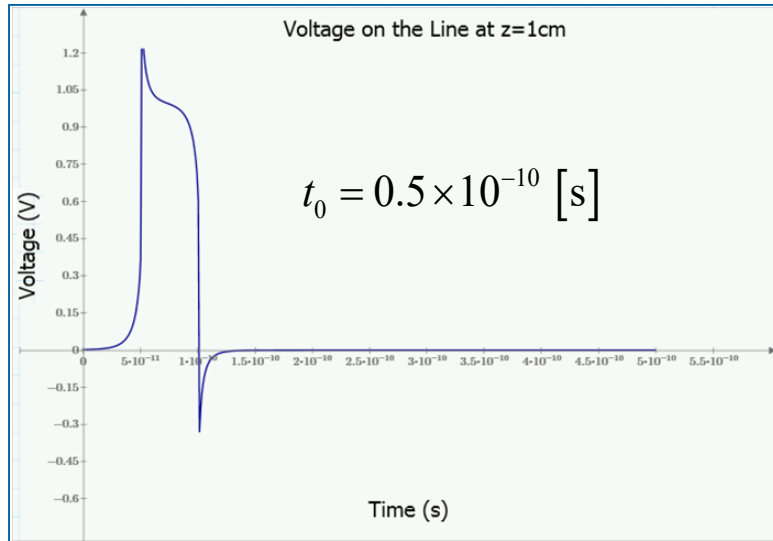
$$\sigma = 3.0 \times 10^7 \text{ [S/m]}$$



CAD formulas are used to get $\alpha(\omega)$, $\beta(\omega)$.

Signal Propagation (cont.)

Propagation on a microstrip line



$$\epsilon_r = 2.33$$

$$\tan \delta = 0.001$$

$$h = 0.787 \text{ [mm]} \text{ (31 mils)}$$

$$w = 2.35 \text{ [mm]}$$

$$t = 0.0175 \text{ [mm]} \text{ ("half oz" copper cladding)}$$

$$\sigma = 3.0 \times 10^7 \text{ [S/m]}$$

Dispersionless System

Dispersionless System with Constant Attenuation

$$H(\omega) = A(\omega)e^{-j\beta(\omega)l}$$

$$A(\omega) \approx \text{constant} \equiv A_0 = e^{-\alpha(\omega_0)l}$$

$$\Rightarrow s_o(t) = \frac{1}{\pi} \int_0^{\infty} \text{Re} \left\{ A_0 e^{-j\beta(\omega)l} \tilde{s}_i(\omega) e^{j\omega t} \right\} d\omega$$

$$= A_0 \frac{1}{\pi} \text{Re} \int_0^{\infty} \tilde{s}_i(\omega) e^{j\omega \left(t - \frac{l}{v_{p0}} \right)} d\omega$$

$$v_p \equiv \frac{\omega}{\beta} = \text{constant} = v_{p0}$$

No dispersion:
The phase velocity is constant
(not a function of frequency).

$$s_o(t) = A_0 s_i \left(t - \frac{l}{v_{p0}} \right)$$

The output is a delayed and
scaled version of the input.

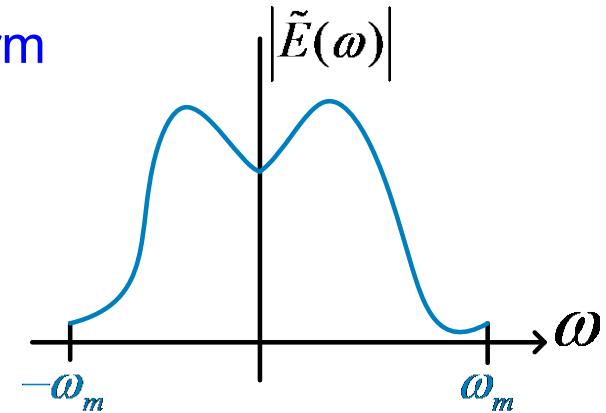
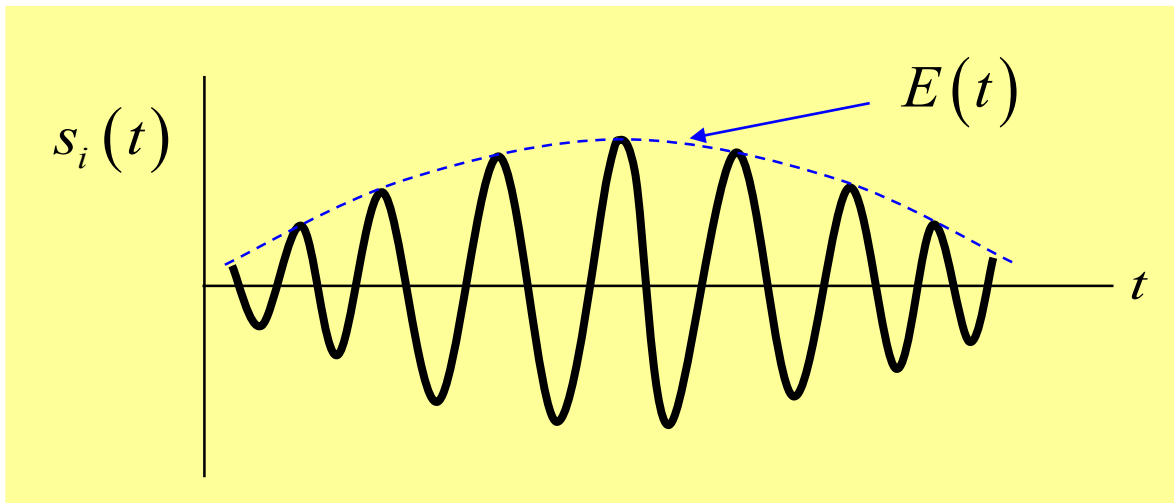
The output signal has no distortion.

Narrow-Band Signal

Low-Loss System with Dispersion and Narrow-Band Signal

Now consider a narrow-band input signal of the form

$$s_i(t) = E(t) \cos(\omega_0 t)$$



Narrow band

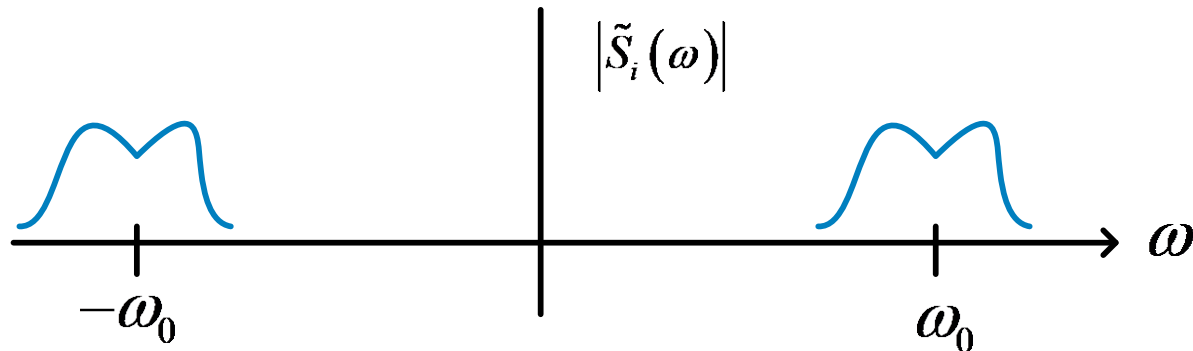
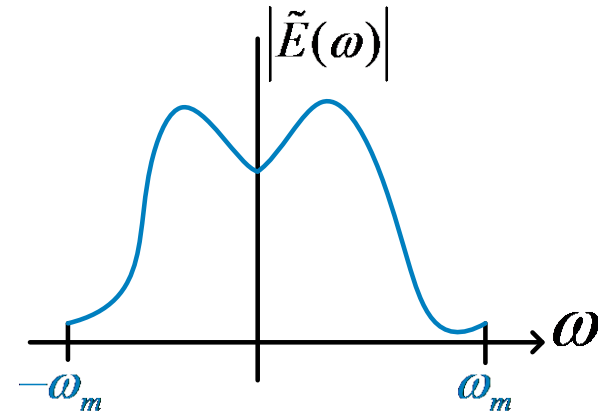
$$\Rightarrow \omega_m \ll \omega_0$$

Physically, the envelope is slowing varying compared with the carrier.

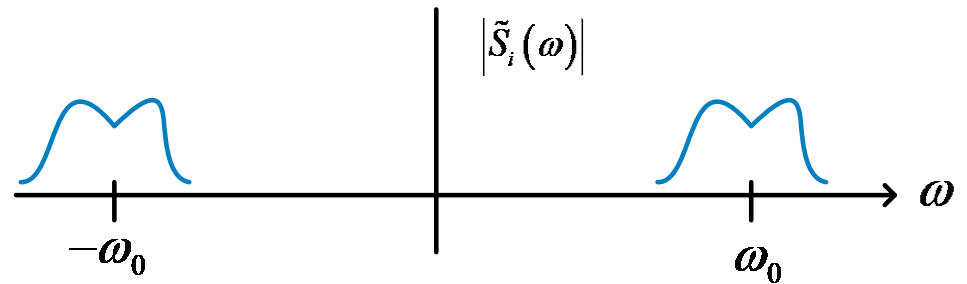
Narrow-Band Signal (cont.)

$$s_i(t) = E(t) \cos(\omega_0 t)$$

$$\begin{aligned}\tilde{s}_i(\omega) &= F\{E(t) \cos(\omega_0 t)\} \\ &= \frac{1}{2} F\{E(t) e^{j\omega_0 t}\} + \frac{1}{2} F\{E(t) e^{-j\omega_0 t}\} \\ &= \frac{1}{2} [\tilde{E}(\omega - \omega_0) + \tilde{E}(\omega + \omega_0)]\end{aligned}$$



Narrow-Band Signal (cont.)



$$\tilde{s}_i(\omega) = \frac{1}{2} [\tilde{E}(\omega - \omega_0) + \tilde{E}(\omega + \omega_0)]$$

Hence, we have:

$$\begin{aligned}
 s_o(t) &= \frac{1}{\pi} \operatorname{Re} \int_0^{\infty} A(\omega) \tilde{s}_i(\omega) e^{-j\beta(\omega)l} e^{j\omega t} d\omega && \left(A(\omega) = e^{-\alpha(\omega)l} \right) \\
 &\approx \frac{1}{\pi} \operatorname{Re} \int_0^{\infty} A(\omega) \frac{1}{2} \tilde{E}(\omega - \omega_0) e^{j(\omega t - \beta(\omega)l)} d\omega && \text{(Ignore } \tilde{E}(\omega + \omega_0)) \\
 &= \frac{1}{2\pi} \operatorname{Re} \int_0^{\infty} A(\omega) \tilde{E}(\omega - \omega_0) e^{j(\omega t - \beta(\omega)l)} d\omega
 \end{aligned}$$

Narrow-Band Signal (cont.)

Since the signal is narrow band, using a Taylor series expansion about ω_0 results in:

$$\beta(\omega) \approx \underbrace{\beta(\omega_0)}_{\beta_0} + \underbrace{\left. \frac{d\beta}{d\omega} \right|_{\omega_0}}_{\beta'_0} (\omega - \omega_0) + \dots \approx \beta_0 + \beta'_0 (\omega - \omega_0)$$

$$A(\omega) \approx \underbrace{A(\omega_0)}_{A_0} + \underbrace{\left. \frac{dA}{d\omega} \right|_{\omega_0}}_{\text{neglect}} (\omega - \omega_0) + \dots \approx A_0$$

↑
Low loss assumption

$$A(\omega) = e^{-\alpha(\omega)z}$$

$$\frac{dA(\omega)}{d\omega} = -z \frac{d\alpha(\omega)}{d\omega} e^{-\alpha(\omega)z}$$

Small ↙

Narrow-Band Signal (cont.)

Thus,

$$\begin{aligned}
 s_o(t) &\approx \frac{1}{2\pi} \operatorname{Re} \int_0^{\infty} A(\omega) \tilde{E}(\omega - \omega_0) e^{j(\omega t - \beta(\omega)l)} d\omega \\
 &\approx \frac{1}{2\pi} \operatorname{Re} \int_0^{\infty} A_0 \tilde{E}(\omega - \omega_0) e^{j\omega t} e^{-j\beta_0 l} e^{-j\beta'_0(\omega - \omega_0)l} d\omega \quad (\text{use Taylor approximation}) \\
 &= \frac{A_0}{2\pi} \operatorname{Re} \left[e^{-j\beta_0 l} e^{j\omega_0 t} \int_0^{\infty} \tilde{E}(\omega - \omega_0) e^{-j\beta'_0(\omega - \omega_0)l} e^{j(\omega - \omega_0)t} d\omega \right] \quad (\text{multiply and divide by } e^{j\omega_0 t}) \\
 &= \frac{A_0}{2\pi} \operatorname{Re} \left[e^{j(\omega_0 t - \beta_0 l)} \int_{-\omega_0}^{\infty} \tilde{E}(\omega_s) e^{-j\beta'_0 \omega_s l} e^{j\omega_s t} d\omega_s \right] \quad (\omega_s \equiv \omega - \omega_0) \\
 &\approx \frac{A_0}{2\pi} \operatorname{Re} \left[e^{j(\omega_0 t - \beta_0 l)} \int_{-\infty}^{\infty} \tilde{E}(\omega_s) e^{-j\beta'_0 \omega_s l} e^{j\omega_s t} d\omega_s \right] \quad (\text{extend lower limit to } -\infty)
 \end{aligned}$$

The spectrum of $E(t)$ is concentrated near $\omega = 0$.

Narrow-Band Signal (cont.)

$$\begin{aligned} s_0(t) &\approx \frac{A_0}{2\pi} \operatorname{Re} \left[e^{j(\omega_0 t - \beta_0 l)} \int_{-\infty}^{\infty} \tilde{E}(\omega_s) e^{j\omega_s(t - \beta'_0 l)} d\omega_s \right] \\ &= A_0 \operatorname{Re} \left[e^{j(\omega_0 t - \beta'_0 l)} E(t - \beta'_0 l) \right] \\ &= A_0 \cos \left(\omega_0 \left(t - \frac{\beta_0}{\omega_0} l \right) \right) E(t - \beta'_0 l) \end{aligned}$$

Define:

$$v_p \equiv \frac{\omega_0}{\beta_0} \quad \text{phase velocity @ } \omega_0$$

Define:

$$v_g \equiv \frac{1}{\beta'_0} = \left. \left(\frac{d\omega}{d\beta} \right) \right|_{\omega_0} \quad \text{group velocity @ } \omega_0$$

Narrow-Band Signal (cont.)

Summary

$$s_o(t) \approx A_0 E \left(t - \frac{l}{v_g} \right) \cos \left(\omega_0 \left(t - \frac{l}{v_p} \right) \right)$$

Envelope travels with group velocity

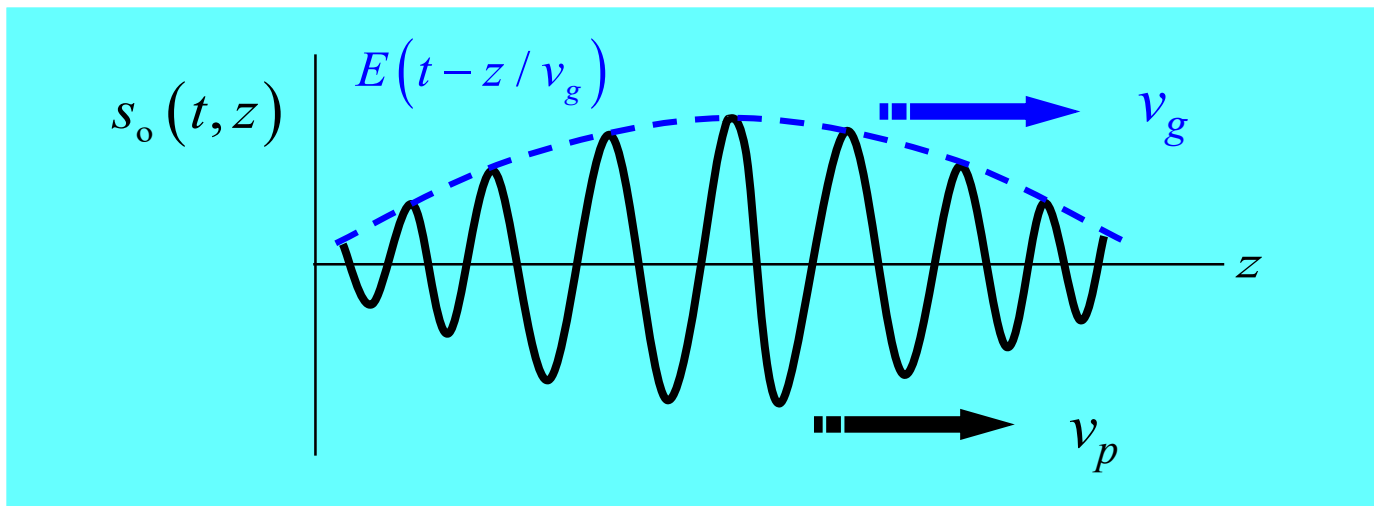
Carrier travels with phase velocity

$$v_p \equiv \frac{\omega_0}{\beta_0}$$

$$v_g \equiv \left. \left(\frac{d\omega}{d\beta} \right) \right|_{\omega_0}$$

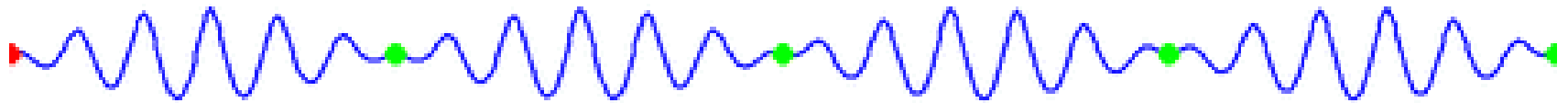
Narrow-Band Signal (cont.)

$$s_o(t) \approx A_0 E \left(t - \frac{l}{v_g} \right) \cos \left(\omega_0 \left(t - \frac{l}{v_p} \right) \right)$$



Narrow-Band Signal (cont.)

Example from Wikipedia (view in full-screen mode with pptx)



Red dot: phase velocity

Green dot: group velocity

Phase velocity $>$ group velocity

http://en.wikipedia.org/wiki/Group_velocity

Narrow-Band Signal (cont.)

Note on dispersion

Assume:

No dispersion

$$\Rightarrow v_p = v_g$$

In this case the envelope and carrier are delayed the same.

Proof

Assume :

$$v_p = \frac{\omega}{\beta} = \text{constant} = c_1$$

$$\Rightarrow \omega = c_1 \beta$$

$$\Rightarrow \frac{d\omega}{d\beta} = c_1$$

$$\Rightarrow v_g = c_1$$

Example: lossless transmission line

Example: TE₁₀ Mode of Rectangular Waveguide

Recall:
$$\beta = \sqrt{\omega^2 \mu \epsilon - \left(\frac{\pi}{a}\right)^2}$$

Phase velocity:
$$v_p = \frac{\omega}{\beta}$$

Group velocity:
$$v_g = \frac{d\omega}{d\beta} = \left(\frac{d\beta}{d\omega}\right)^{-1}$$

Observation:

$$v_p v_g = \frac{1}{\mu \epsilon} = c_d^2$$

After simple calculation:

$$v_p = \frac{\omega}{\sqrt{\omega^2 \mu \epsilon - \left(\frac{\pi}{a}\right)^2}}$$

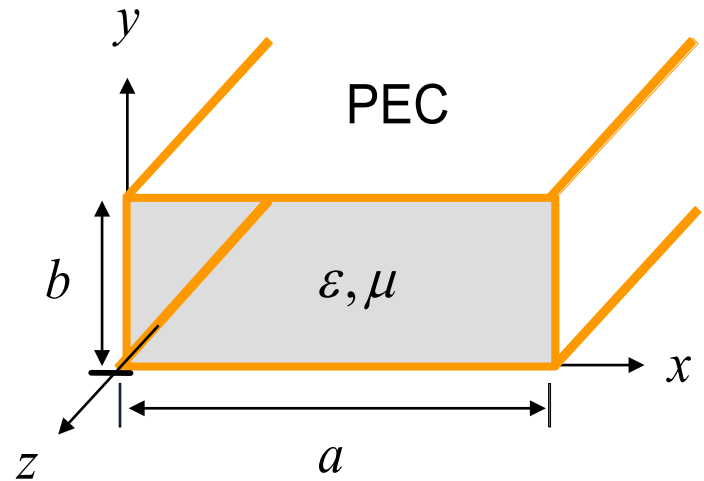
$$v_g = \frac{1}{\omega \mu \epsilon} \sqrt{\omega^2 \mu \epsilon - \left(\frac{\pi}{a}\right)^2}$$

(This final result is valid for any mode of a lossless waveguide.)

Example (cont.)

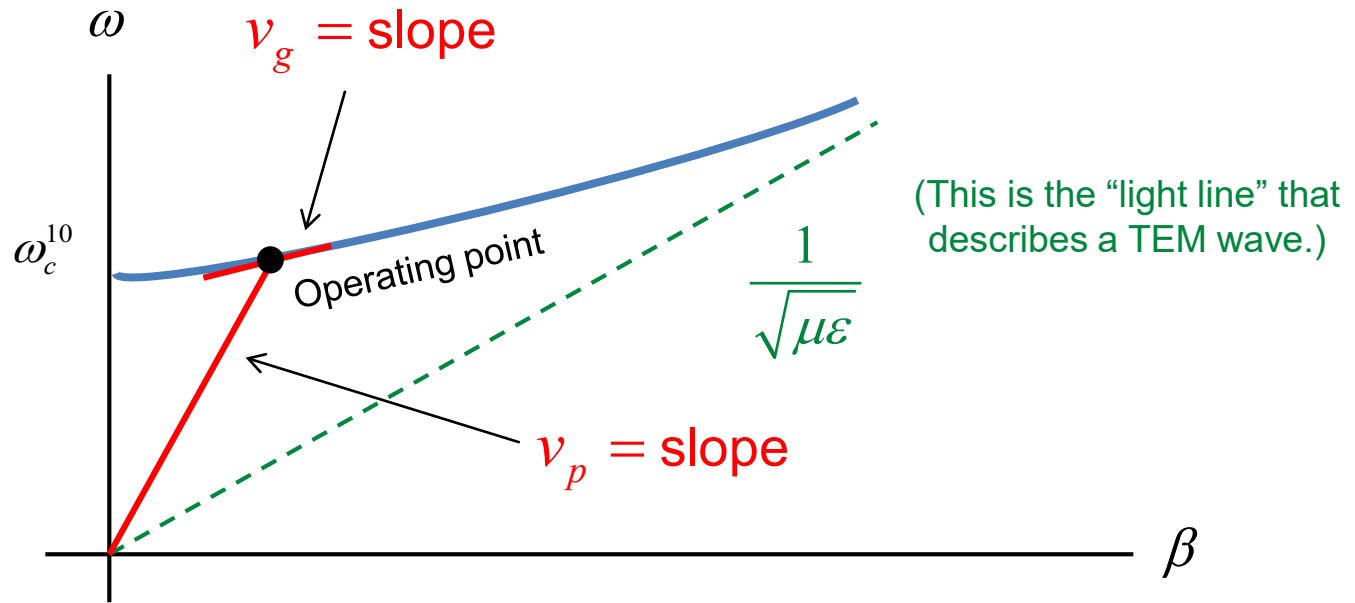
Lossless Case ($\epsilon_c = \epsilon = \epsilon'$)

$$f > f_c$$



$$v_p = \frac{\omega}{\beta}$$

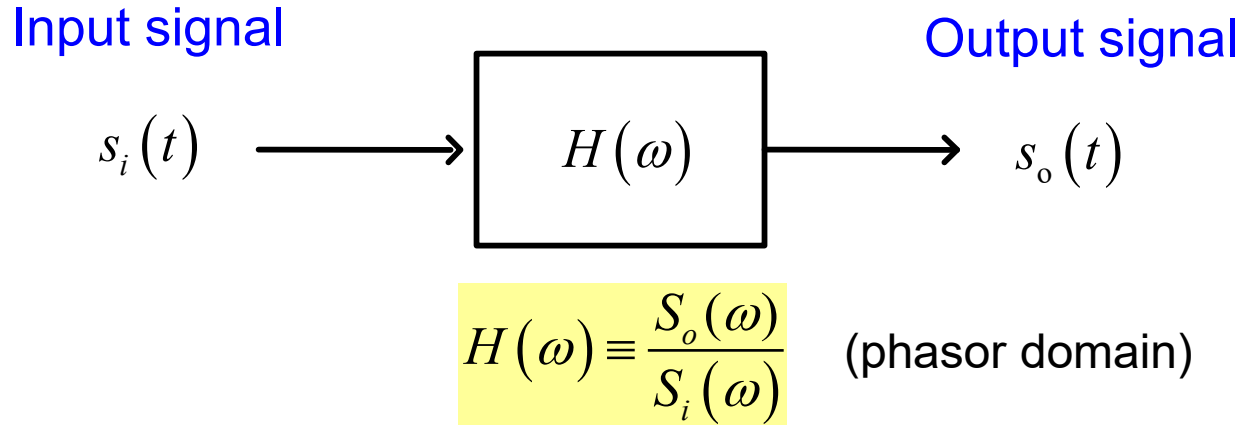
$$v_g = \frac{d\omega}{d\beta}$$



$$v_p > c_d$$

$$v_g < c_d$$

Filter Response



What we have done also applies to a filter, but here we use the transfer function phase directly, and do not introduce a phase constant β .

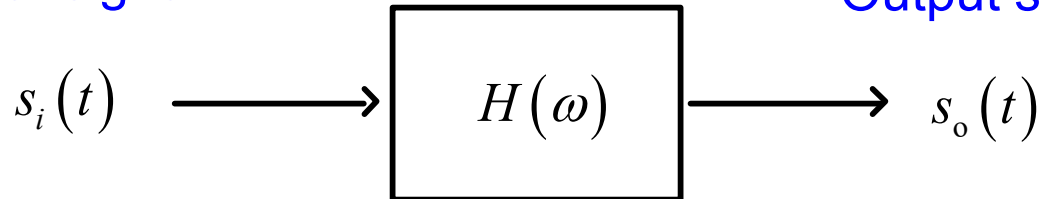
$$H(\omega) = A(\omega)e^{j\phi(\omega)}$$

From the previous results, we have

$$s_o(t) = \frac{1}{\pi} \int_0^{\infty} \text{Re} \left\{ A(\omega) e^{j\phi(\omega)} \tilde{s}_i(\omega) e^{j\omega t} \right\} d\omega$$

Filter Response (cont.)

Input signal



Output signal

$$H(\omega) = A(\omega) e^{j\phi(\omega)}$$

Let $\beta l \rightarrow -\phi$

$$s_o(t) \approx A_0 E(t - \beta_0' l) \cos\left(\omega_0 \left(t - \frac{\beta_0 l}{\omega_0}\right)\right)$$

Assume we have our modulated input signal:

$$s_i(t) = E(t) \cos(\omega_0 t)$$

The output is:

$$s_o(t) \approx A_0 E(t + \phi_0') \cos\left(\omega_0 \left(t + \frac{\phi_0}{\omega_0}\right)\right)$$

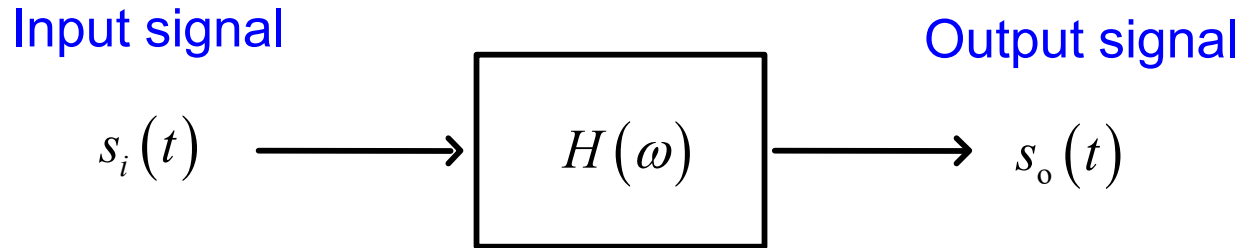
where

$$A_0 = A(\omega_0)$$

$$\phi_0 = \phi(\omega_0)$$

$$\phi_0' = \left. \frac{d\phi}{d\omega} \right|_{\omega_0}$$

Filter Response (cont.)



$$s_o(t) \approx A_0 E(t + \phi'_0) \cos\left(\omega_0 \left(t + \frac{\phi_0}{\omega_0}\right)\right)$$

This motivates the following definitions:

Phase delay:

$$\tau_p \equiv -\frac{\phi_0}{\omega_0}$$

Group delay:

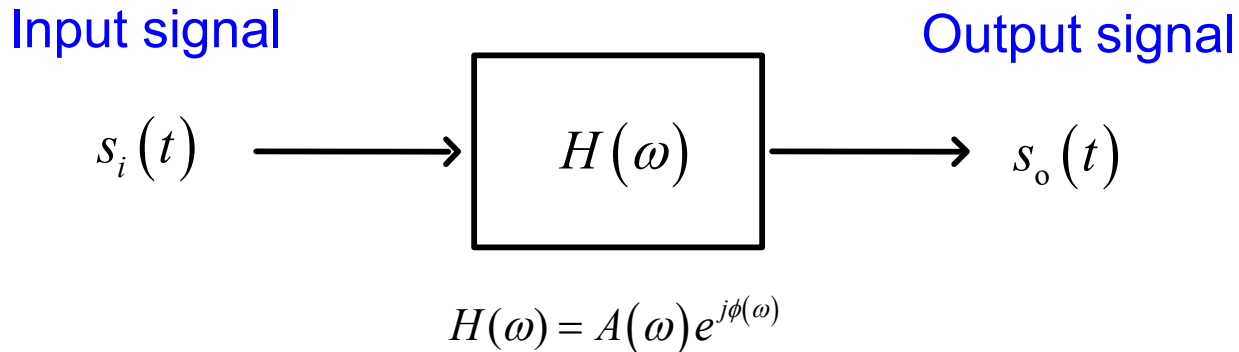
$$\tau_g \equiv -\left.\frac{d\phi}{d\omega}\right|_{\omega_0}$$

(The units are seconds.)

→ $s_o(t) \approx A_0 E(t - \tau_g) \cos\left(\omega_0 (t - \tau_p)\right)$

Filter Response (cont.)

Summary



$$s_i(t) = E(t) \cos(\omega_0 t)$$

$$s_o(t) \approx A_0 E(t - \tau_g) \cos\left(\omega_0 (t - \tau_p)\right)$$

Phase delay:

$$\tau_p \equiv -\frac{\phi_0}{\omega_0}$$

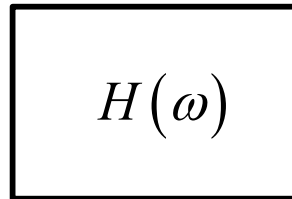
Group delay:

$$\tau_g \equiv -\left. \frac{d\phi}{d\omega} \right|_{\omega_0}$$

Linear-Phase Filter Response

Input signal

$$s_i(t)$$



Output signal

$$s_o(t)$$

$$H(\omega) = A(\omega)e^{j\phi(\omega)}$$

Ideal linear phase filter:

$$\phi(\omega) = -\tau\omega$$

$$A(\omega) \approx A_0$$

The attenuation of the ideal filter is constant, at least over the bandwidth of the filter.

Recall:

$$s_o(t) = \frac{1}{\pi} \int_0^{\infty} \text{Re} \left\{ A(\omega) e^{j\phi(\omega)} \tilde{s}_i(\omega) e^{j\omega t} \right\} d\omega$$

Hence

$$s_o(t) = \frac{1}{\pi} \int_0^{\infty} \text{Re} \left\{ \left(A_0 e^{-j\tau\omega} \right) \tilde{s}_i(\omega) e^{j\omega t} \right\} d\omega$$

$$\phi(\omega) = -\tau\omega$$

$$\Rightarrow -\frac{\phi}{\omega} = -\frac{d\phi}{d\omega} = \tau$$

$$\Rightarrow \tau_p = \tau_g = \tau = \text{constant}$$

The envelope and carrier are delayed the same.

Linear-Phase Filter Response (cont.)

We then have:

$$\begin{aligned} s_o(t) &= \frac{1}{\pi} \int_0^{\infty} \operatorname{Re} \left\{ \left(A_0 e^{-j\tau\omega} \right) \tilde{s}_i(\omega) e^{j\omega t} \right\} d\omega \\ &= \frac{1}{\pi} A_0 \operatorname{Re} \left\{ \int_0^{\infty} \tilde{s}_i(\omega) e^{j\omega(t-\tau)} d\omega \right\} \\ &= A_0 s_i(t - \tau) \end{aligned}$$

so

$$s_o(t) = A_0 s_i(t - \tau)$$

An ideal linear-phase filter does not distort the signal.

It may be desirable to have a filter maintain a linear phase, at least over the bandwidth of the filter. This will tend to minimize signal distortion.