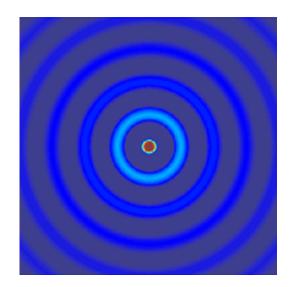


#### Spring 2016

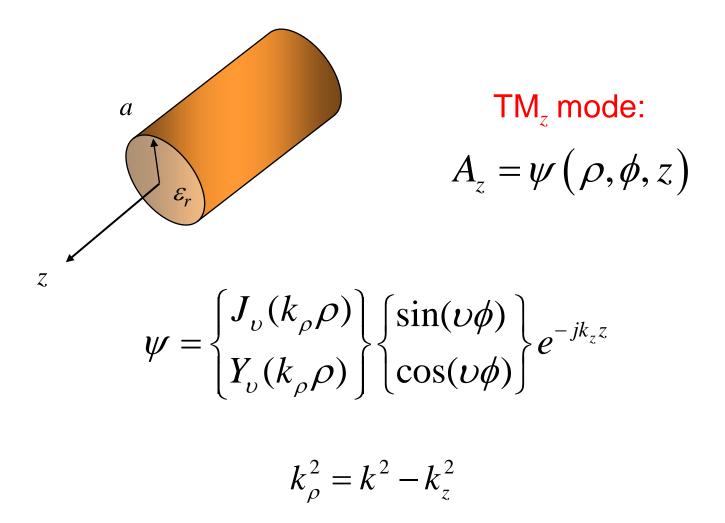
Prof. David R. Jackson ECE Dept.

Notes 9



## **Circular Waveguide**

The waveguide is homogeneously filled, so we have independent  $TE_z$  and  $TM_z$  modes.



(1) 
$$\phi$$
 variation  $\phi \in [0, 2\pi]$   
 $\psi(\rho, \phi + 2\pi, z) = \psi(\rho, \phi, z)$  (uniqueness of solution)  
 $\implies \upsilon = n$ 

Choose  $\cos(n\phi)$ 

$$\Psi = \begin{cases} J_n(k_\rho \rho) \\ Y_n(k_\rho \rho) \end{cases} \cos(n\phi) \ e^{-jk_z z}$$

(2) The field should be finite on the *z* axis  $\psi(0, \phi, z)$ 

$$\implies Y_n(k_\rho \rho)$$
 is not allowed

$$\psi = \cos(n\phi) J_n(k_\rho \rho) e^{-jk_z z}$$

$$k_{\rho}^2 = k^2 - k_z^2$$

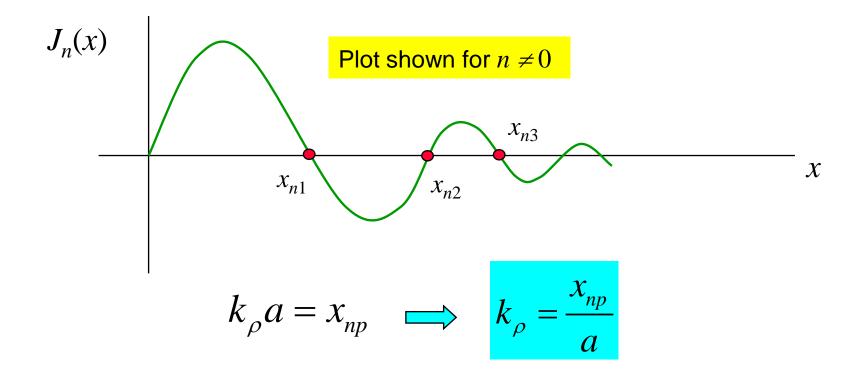
(3) B.C.'s: 
$$E_z(a, \phi, z) = 0$$

$$\begin{split} E_{z} &= \frac{1}{j\omega\mu\varepsilon} \left( k^{2} + \frac{\partial^{2}}{\partial z^{2}} \right) \psi \\ &= \frac{1}{j\omega\mu\varepsilon} \left( k^{2} - k_{z}^{2} \right) \psi \\ &= \left( \frac{k_{\rho}^{2}}{j\omega\mu\varepsilon} \right) \psi \end{split}$$

so 
$$\psi(a,\phi,z) = 0$$

Hence 
$$J_n(k_\rho a) = 0$$

$$J_n(k_\rho a) = 0$$



Note:  $x_{n0} = 0$  is not included since  $J_n\left(x_{np}\frac{\rho}{a}\right) = 0$  (trivial soln.)

 $\mathsf{TM}_{np}$  mode:

$$A_{z} = \cos(n\phi) J_{n}\left(x_{np}\frac{\rho}{a}\right)e^{-jk_{z}z} \quad n = 0, 1, 2...$$

$$k_{z} = \left(k^{2} - \left(\frac{x_{np}}{a}\right)^{2}\right)^{1/2} \qquad p = 1, 2, 3, \dots$$

#### Cutoff Frequency: TM<sub>z</sub>

$$k_z^2 = k^2 - k_\rho^2$$

$$k_z = 0$$
  $\longrightarrow$   $k = k_\rho = \frac{x_{np}}{a}$ 

$$2\pi f_c \sqrt{\mu\varepsilon} = \frac{x_{np}}{a}$$

$$f_c^{TM} = \left(\frac{c}{2\pi a\sqrt{\varepsilon_r}}\right) x_{np}$$

Cutoff Frequency: TM<sub>z</sub> (cont.)

$x_{np}$ values									
$p \setminus n$	0	1	2	3	4	5			
1	2.405	3.832	5.136	6.380	7.588	8.771			
2	5.520	7.016	8.417	9.761	11.065	12.339			
3	8.654	10.173	11.620	13.015	14.372				
4	11.792	13.324	14.796						

 $TM_{01}, TM_{11}, TM_{21}, TM_{02}, \dots$ 

**TE**<sub>z</sub> Modes

$$F_z = \psi(\rho, \phi, z)$$

$$\psi = \cos(n\phi) J_n(k_\rho \rho) e^{-jk_z z}$$

$$H_z = \frac{k_\rho^2}{j\omega\mu\varepsilon}\psi$$

Note: 
$$\psi(a, \phi, z) \neq 0$$

#### TE<sub>z</sub> Modes (cont.)

Set  $E_{\phi}(a,\phi,z) = 0$ 

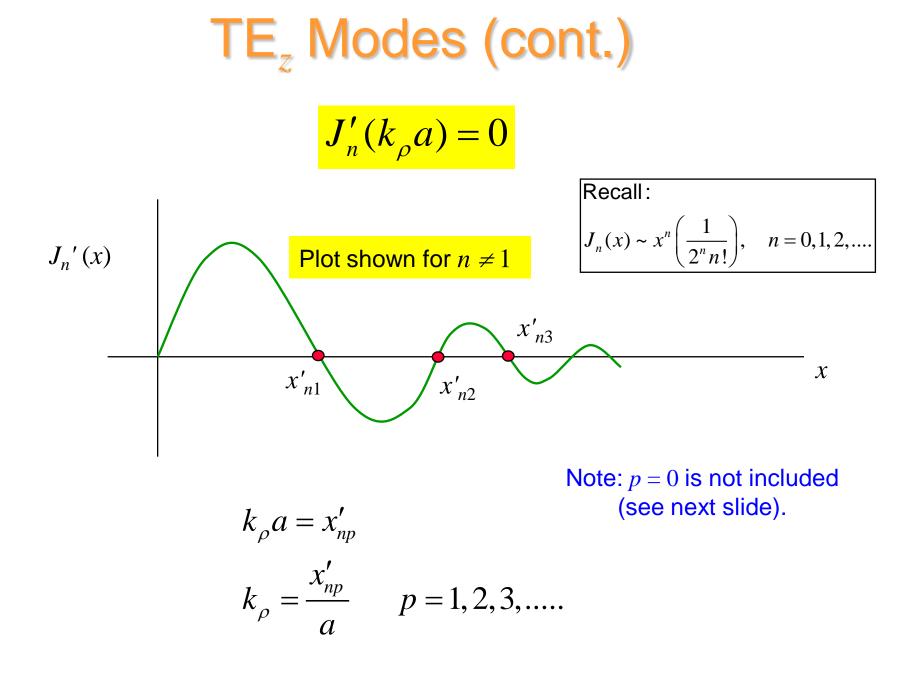
$$E_{\phi} = \frac{1}{\varepsilon} \frac{\partial \psi}{\partial \rho}$$

**SO** 

$$\frac{\partial \psi}{\partial \rho}\Big|_{\rho=a} = 0$$

Hence

$$J'_n(k_\rho a) = 0$$



#### $TE_z$ Modes (cont.)

$$\psi = \cos(n\phi) J_n\left(x'_{np}\frac{\rho}{a}\right)e^{-jk_z z}$$
  $p = 1, 2, ...$ 

Note: If p = 0  $x'_{np} = 0$ 

$$n \neq 0$$
  $J_n\left(x'_{np}\frac{\rho}{a}\right) = J_n\left(0\right) = 0$  (trivial soln.)

$$n = 0 \qquad J_0\left(x'_{np}\frac{\rho}{a}\right) = J_0\left(0\right) = 1$$

$$\psi = e^{-jk_z z} = e^{-jkz}$$

(trivial fields)

#### **Cutoff Frequency: TE**<sub>z</sub>

$$k_z^2 = k^2 - k_\rho^2$$

$$k_z = 0$$
  $\implies$   $k_\rho = k = \frac{x'_{np}}{a}$ 

$$2\pi f_c \sqrt{\mu\varepsilon} = \frac{x'_{np}}{a}$$

$$f_c^{TE} = \left(\frac{c}{2\pi a\sqrt{\varepsilon_r}}\right) x'_{np}$$

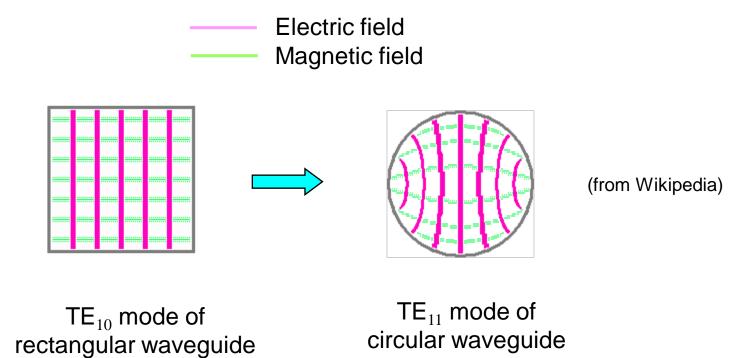
# **Cutoff Frequency:TE**<sub>z</sub>

$x'_{np}$ values									
$p \setminus n$	0	1	2	3	4	5			
1	3.832	1.841	3.054	4.201	5.317	5.416			
2	7.016	5.331	6.706	8.015	9.282	10.520			
3	10.173	8.536	9.969	11.346	12.682	13.987			
4	13.324	11.706	13.170						

 $TE_{11}, TE_{21}, TE_{01}, TE_{31}, \dots$ 



#### The dominant mode of circular waveguide is the $TE_{11}$ mode.



The  $TE_{11}$  mode can be thought of as an evolution of the  $TE_{10}$  mode of rectangular waveguide as the boundary changes shape.

## Attenuation Property of TE<sub>01</sub> Mode

**Goal:** We wish to study the high-frequency dependence of attenuation on frequency for circular waveguide modes, and show the interesting behavior of the  $TE_{01}$  mode (the loss <u>decreases</u> as frequency increases).

$$\alpha_c = \frac{\langle \mathcal{P}_d \rangle}{2 \langle \mathcal{P}_f \rangle}$$

 $\langle \mathscr{P}_{d} \rangle = \oint \frac{1}{2} R_{s} \left| \underline{H}_{t} \right|^{2} dl$ 

$$R_{s} = \frac{1}{\sigma\delta} = \frac{1}{\sigma\sqrt{\frac{2}{\omega\mu\sigma}}}$$
$$= \sqrt{\frac{\omega\mu}{2\sigma}}$$
$$= \mathcal{O}\left(\sqrt{\omega}\right)$$

 $TE_z$  Mode:

$$H_{\phi} = \frac{1}{j\omega\mu\varepsilon} \frac{1}{\rho} \frac{\partial^2 F_z}{\partial\phi\partial z}$$

 $H_{z} = \frac{1}{i\omega\mu\varepsilon}\kappa_{\rho}r_{z}$ 

Assume that 
$$F_z$$
 is  
order 1 as the  
frequency increases.

Recall that  $k_{\rho}$  is a constant.

$$\begin{split} H_{z} &= \mathscr{O}\left(\frac{1}{\omega}\right) & \text{Note:} \\ H_{\phi} &= \begin{cases} \mathscr{O}(1) & n \neq 0 \\ 0 & n = 0 \end{cases} & k_{z} &= \mathscr{O}(k) = \mathscr{O}(\omega) \end{split}$$

$$n \neq 0 \qquad n = 0$$

$$< \mathscr{P}_{d} >= \mathscr{O}\left(\sqrt{\omega}\right) \left[ \mathscr{O}\left(\frac{1}{\omega^{2}}\right) + \mathscr{O}(1) \right] \qquad < \mathscr{P}_{d} >= \mathscr{O}\left(\sqrt{\omega}\right) \left[ \mathscr{O}\left(\frac{1}{\omega^{2}}\right) + 0 \right]$$

$$= \mathscr{O}\left(\omega^{1/2}\right) \qquad = \mathscr{O}\left(\omega^{-3/2}\right)$$

From the  $TE_z$  table:

$$\underline{E} = \mathscr{O}(1) \quad \text{e.g.} \quad E_{\phi} = \frac{1}{\varepsilon} \frac{\partial F_z}{\partial \rho}$$
$$\underline{H} = \mathscr{O}(1) \quad \text{e.g.} \quad H_{\rho} = \frac{1}{j \omega \mu \varepsilon} \frac{\partial^2 F_z}{\partial \rho \partial z}$$

Hence 
$$\left\langle \mathscr{P}_{f} \right\rangle = \mathscr{O}\left(1\right)$$

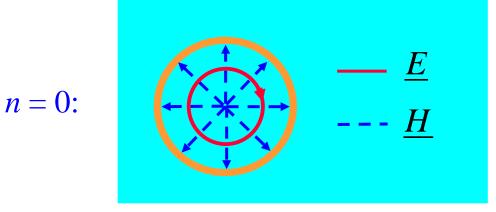
Hence

If 
$$n \neq 0$$
:  $\alpha_c = \mathcal{O}(\omega^{1/2})$ 

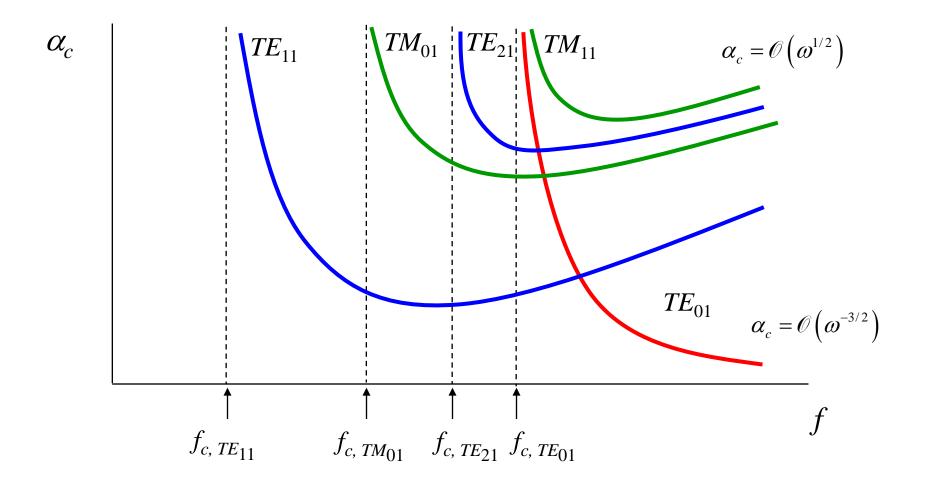
Usual behavior for rectangular waveguides

If 
$$n = 0$$
:  $\alpha_c = \mathcal{O}(\omega^{-3/2})$ 

Decreases with frequency!



Note: The mode  $TE_{0p}$  mode can be supported by a series of concentric rings, since there is no longitudinal (*z*-directed) current ( $H_{\phi} = 0$ ).



The  $TE_{01}$  mode was studied extensively as a candidate for long-range communications – but was not competitive with antennas. Also, fiber-optic cables eventually became available with lower loss than the  $TE_{01}$  mode. It is still useful for some applications (e.g., high power).

From the beginning, the most obvious application of waveguides had been as a communications medium. It had been determined by both Schelkunoff and Mead, independently, in July 1933, that an axially symmetric electric wave (TE<sub>01</sub>) in circular waveguide would have an attenuation factor that decreased with increasing frequency [44]. This unique characteristic was believed to offer a great potential for wide-band, multichannel systems, and for many years to come the development of such a system was a major focus of work within the waveguide group at BTL. It is important to note, however, that the use of waveguide as a long transmission line never did prove to be practical, and Southworth eventually began to realize that the role of waveguide would be somewhat different than originally expected. In a memorandum dated October 23, 1939, he concluded that microwave radio with highly directive antennas was to be preferred to long transmission lines. "Thus," he wrote, "we come to the conclusion that the hollow, cylindrical conductor is to be valued primarily as a new circuit element, but not yet as a new type of toll cable" [45]. It was as a circuit element in military radar that waveguide technology was to find its first major application and to receive an enormous stimulus to both practical and theoretical advance.

K. S. Packard, "The Origins of Waveguide: A Case of Multiple Rediscovery," *IEEE Trans. MTT*, pp. 961-969, Sept. 1984.

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