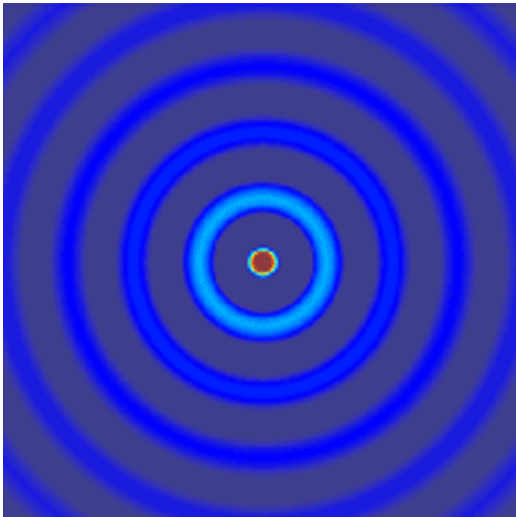


ECE 6341

Spring 2016

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ECE Dept.

Notes 28



Asymptotic Evaluation of Integral

Goal: Evaluate the following integral as Ω becomes large:

$$I(\Omega) = \int_a^b f(x) e^{j\Omega g(x)} dx$$

$g(x)$ is a real function here.

$$\Omega \rightarrow \infty$$

Case 1: $g'(x) \neq 0, x \in [a, b]$

Integration by Parts

Case 2: $g'(x_0) = 0, x_0 \in (a, b)$

Stationary-Phase Method

Integration by Parts

$$I(\Omega) = \int_a^b f(x) e^{j\Omega g(x)} dx \quad \Omega \rightarrow \infty$$

$$g'(x) \neq 0 \quad x \in [a, b]$$

Note that

$$\frac{d}{dx} \left[e^{j\Omega g(x)} \right] = (j\Omega g'(x)) e^{j\Omega g(x)}$$

So that we may write

$$I(\Omega) = \int_a^b f(x) \frac{1}{j\Omega g'(x)} \frac{d}{dx} \left[e^{j\Omega g(x)} \right] dx$$

Integration by Parts

$$I(\Omega) = \int_a^b \underbrace{f(x) \frac{1}{j\Omega g'(x)}}_u \frac{d}{dx} \underbrace{\left[e^{j\Omega g(x)} \right]}_v dx$$

$$\int_a^b u(x) \frac{d}{dx} v(x) dx = u(x) v(x) \Big|_a^b - \int_a^b v(x) \frac{d}{dx} u(x) dx$$

Hence

$$I(\Omega) = \frac{f(x)}{j\Omega g'(x)} e^{j\Omega g(x)} \Big|_a^b - \int_a^b e^{j\Omega g(x)} \left(\frac{1}{j\Omega} \right) \frac{d}{dx} \left[\frac{f(x)}{g'(x)} \right] dx$$

Integration by Parts (cont.)

Riemann-Lebesgue Lemma:

$$\int_a^b F(x) e^{j\Omega g(x)} dx \rightarrow 0 \quad \text{as } \Omega \rightarrow \infty$$

(The integrand oscillates faster, so the integral tends to zero)

Hence

$$I = \frac{f(x)}{j\Omega g'(x)} e^{j\Omega g(x)} \Big|_a^b + o(1/\Omega)$$

or

$$I \sim \frac{f(x)}{j\Omega g'(x)} e^{j\Omega g(x)} \Big|_a^b$$

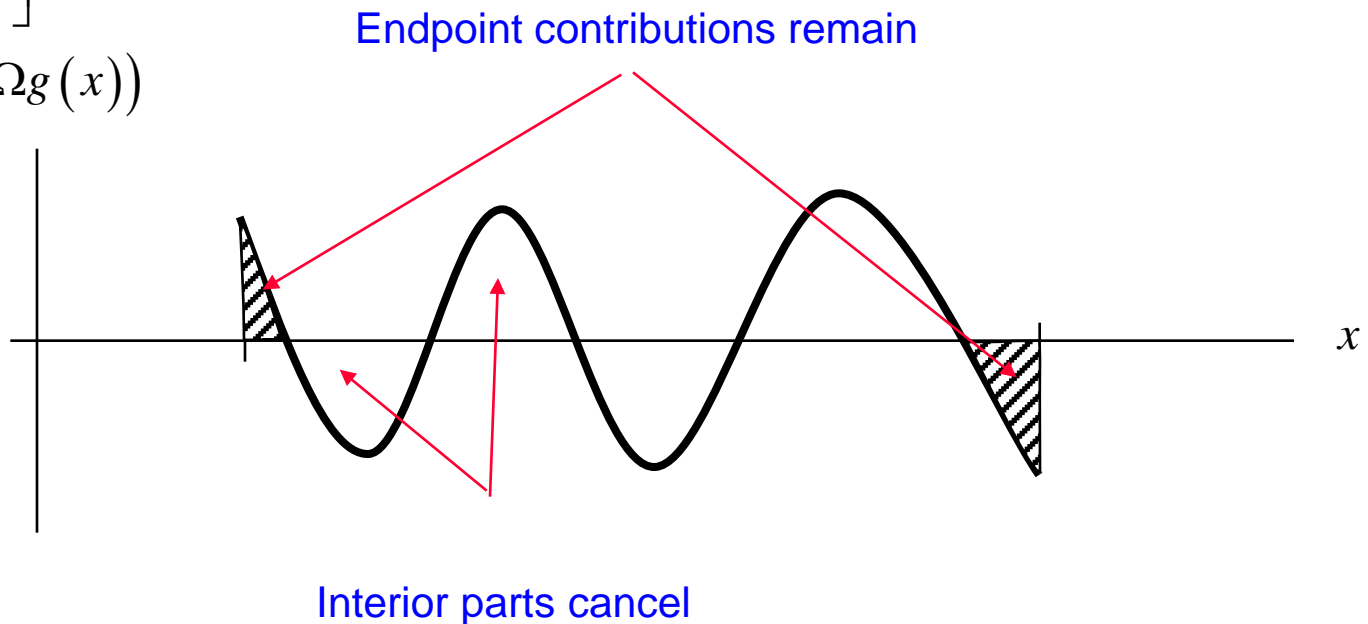
“Small o” notation:
“It decays faster than this as Ω gets large.”

Integration by Parts (cont.)

Physical interpretation

$$I(\Omega) = \int_a^b f(x) e^{j\Omega g(x)} dx \quad I \sim \frac{f(x)}{j\Omega g'(x)} e^{j\Omega g(x)} \Big|_a^b$$

$$\begin{aligned} \operatorname{Re}[f(x) e^{j\Omega g(x)}] \\ = f(x) \cos(\Omega g(x)) \end{aligned}$$



Stationary-Phase Method

$$I(\Omega) = \int_a^b f(x) e^{j\Omega g(x)} dx$$

$$g'(x_0) = 0 \quad x_0 \in (a, b)$$

Assume

$$f(x_0) \neq 0$$

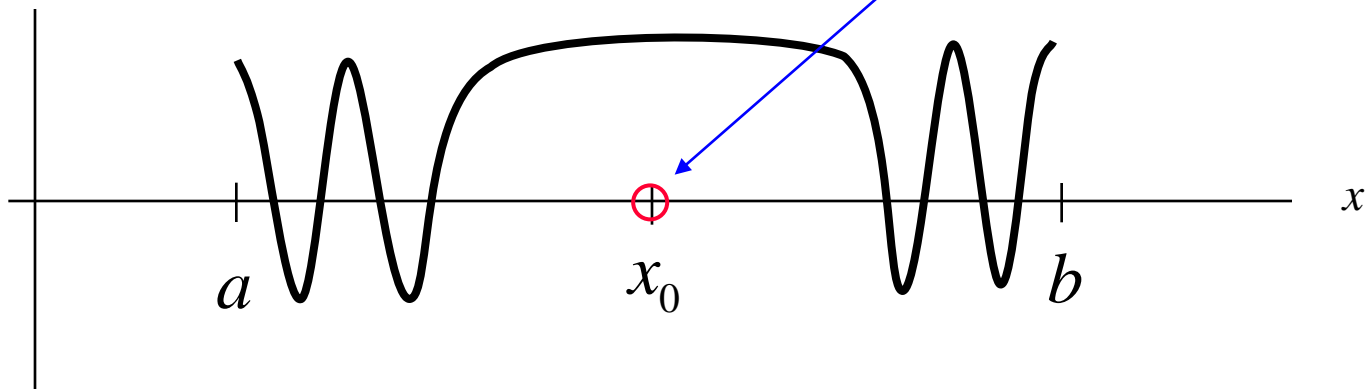
$$g''(x_0) \neq 0$$

$$\operatorname{Re}\left[f(x) e^{j\Omega g(x)} \right]$$

$$= f(x) \cos(\Omega g(x))$$

$$g(x) \approx g(x_0) + \cancel{g'(x_0)(x-x_0)} + \frac{1}{2} g''(x_0)(x-x_0)^2$$

Stationary-phase point



Stationary-Phase Method (cont.)

Conjecture:

$$I \sim \int_{x_0-\Delta}^{x_0+\Delta} f(x) e^{j\Omega g(x)} dx$$

If $\Delta \rightarrow 0$ (slow enough)

To ensure that the stationary phase point region dominates over the outside regions, we define:

$$I_{SPP}(\Omega) = \int_{x_0-\Delta}^{x_0+\Delta} f(x) e^{j\Omega g(x)} dx = \int_{x_0-\Delta}^{x_0+\Delta} f(x) \frac{1}{j\Omega g'(x)} \frac{d}{dx} \left[e^{j\Omega g(x)} \right] dx$$

$$I_{left}(\Omega) = \int_a^{x_0-\Delta} f(x) e^{j\Omega g(x)} dx = \int_a^{x_0-\Delta} f(x) \frac{1}{j\Omega g'(x)} \frac{d}{dx} \left[e^{j\Omega g(x)} \right] dx$$

$$I_{right}(\Omega) = \int_{x_0+\Delta}^b f(x) e^{j\Omega g(x)} dx = \int_{x_0+\Delta}^b f(x) \frac{1}{j\Omega g'(x)} \frac{d}{dx} \left[e^{j\Omega g(x)} \right] dx$$

Stationary-Phase Method (cont.)

We have that

$$|I_{left}| \leq \frac{A_1}{\Omega |g'(x_0 - \Delta)|} \quad |I_{right}| \leq \frac{A_2}{\Omega |g'(x_0 + \Delta)|}$$

For some constants A_1 and A_2 .

We will show later that $|I_{SPP}| = O\left(\frac{1}{\sqrt{\Omega}}\right)$ “Big O” notation:
“It behaves like this as Ω gets large.”

Stationary-Phase Method (cont.)

Hence, we require that

$$\left| I_{left} \right| \ll \frac{1}{\sqrt{\Omega}} \quad \left| I_{right} \right| \ll \frac{1}{\sqrt{\Omega}}$$

Approximating the first derivative of g by a Taylor series,

$$g'(x_0 - \Delta) \approx g'(x_0) + (-\Delta) g''(x_0)$$

We require:

$$\frac{1}{\Omega |g''(x_0) \Delta|} \ll \frac{1}{\sqrt{\Omega}}$$

Hence $\Delta \sqrt{\Omega} \gg 1$ (This needs to be one of our assumptions.)

Stationary-Phase Method (cont.)

Hence, we assume that:

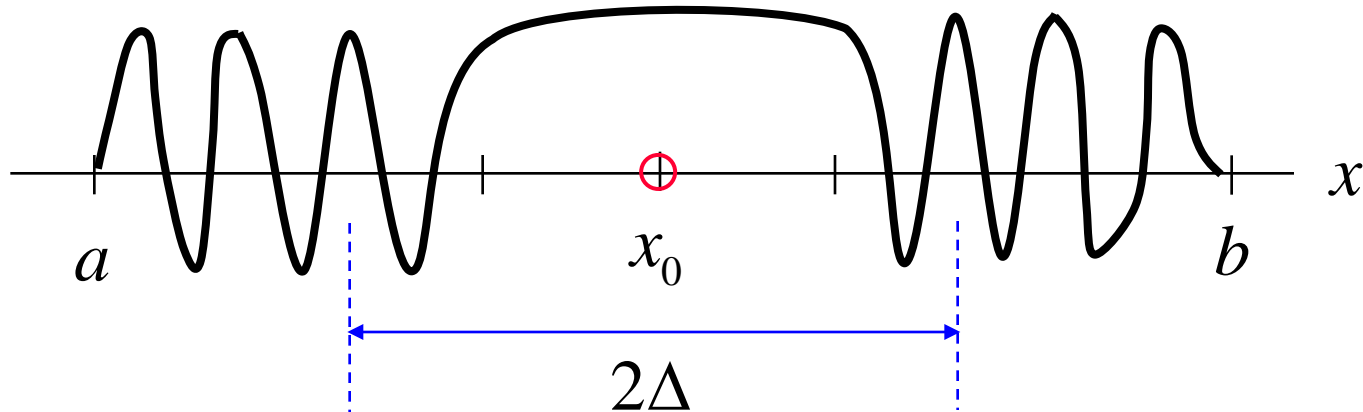
$$\Omega \rightarrow \infty$$

$$\Delta \rightarrow 0$$

$$\Delta\sqrt{\Omega} \rightarrow \infty$$

This ensures that

$$I \sim \int_{x_0-\Delta}^{x_0+\Delta} f(x) e^{j\Omega g(x)} dx$$



Stationary-Phase Method (cont.)

Since $\Delta \rightarrow 0$

Assume $f(x) \approx f(x_0)$

Hence $I \sim f(x_0) \int_{x_0-\Delta}^{x_0+\Delta} e^{j\Omega g(x)} dx$

Stationary-Phase Method (cont.)

$$g(x) \approx g(x_0) + \cancel{g'(x_0)}(x - x_0) + \frac{1}{2}g''(x_0)(x - x_0)^2$$

$$I(\Omega) \approx f(x_0)e^{j\Omega g(x_0)} \int_{x_0-\Delta}^{x_0+\Delta} e^{j\frac{1}{2}g''(x_0)\Omega(x-x_0)^2} dx$$

or

$$I(\Omega) \approx f(x_0)e^{j\Omega g(x_0)} \int_{x_0-\Delta}^{x_0+\Delta} e^{\pm j\frac{1}{2}|g''(x_0)|\Omega(x-x_0)^2} dx$$

where

$$g''(x_0) = \pm |g''(x_0)|$$

Stationary-Phase Method (cont.)

Let

$$s = (x - x_0) \sqrt{\frac{\Omega |g''(x_0)|}{2}} \quad ds = dx \sqrt{\frac{\Omega |g''(x_0)|}{2}}$$

Then

$$I(\Omega) \approx f(x_0) e^{j\Omega g(x_0)} \sqrt{\frac{2}{\Omega |g''(x_0)|}} \int_{-S_L}^{+S_L} e^{\pm js^2} ds$$

where

$$S_L = \Delta \sqrt{\frac{\Omega |g''(x_0)|}{2}}$$

As $\Omega \rightarrow \infty$, $S_L \rightarrow \infty$ since $\Delta \sqrt{\Omega} \rightarrow \infty$

Stationary-Phase Method (cont.)

Therefore

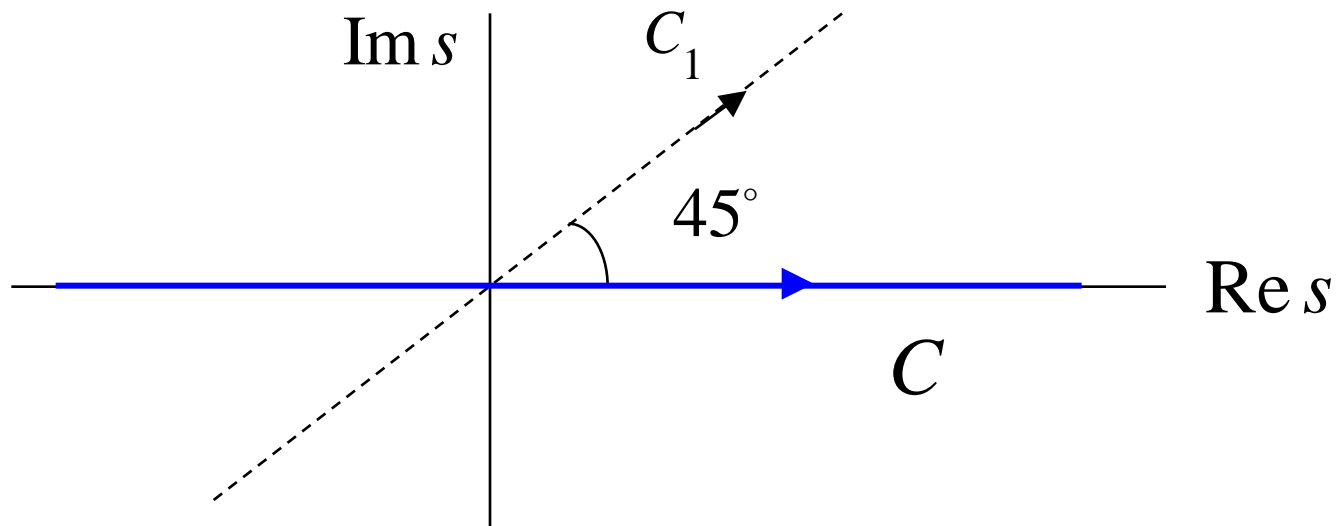
$$I_L \equiv \int_{-S_L}^{+S_L} e^{\pm js^2} ds \rightarrow \int_{-\infty}^{+\infty} e^{\pm js^2} ds \quad \Delta\sqrt{\Omega} \rightarrow \infty$$

Let

$$I_{L1} = \int_{-\infty}^{+\infty} e^{+js^2} ds = \int_C e^{+js^2} ds$$

Deform the path:

$$\int_C \rightarrow \int_{C_1}$$



Stationary-Phase Method (cont.)

Let

$$s = te^{j\frac{\pi}{4}} \quad \longrightarrow \quad \begin{cases} s^2 = t^2 e^{j\frac{\pi}{2}} = jt^2 \\ ds = dt e^{j\frac{\pi}{4}} \end{cases}$$

Hence

$$\begin{aligned} I_{L1} &= e^{j\frac{\pi}{4}} \int_{-\infty}^{+\infty} e^{j(jt^2)} dt \\ &= e^{j\frac{\pi}{4}} \int_{-\infty}^{+\infty} e^{-t^2} dt \\ &= e^{j\frac{\pi}{4}} \sqrt{\pi} \end{aligned}$$

Stationary-Phase Method (cont.)

Similarly,

$$I_{L2} \equiv \int_{-\infty}^{+\infty} e^{-js^2} ds = e^{-j\frac{\pi}{4}} \sqrt{\pi}$$

Then

$$I \approx f(x_0) e^{j\Omega g(x_0)} \sqrt{\frac{2}{\Omega |g''(x_0)|}} e^{\pm j\frac{\pi}{4}} \sqrt{\pi}$$

Note: $I = O\left(\frac{1}{\sqrt{\Omega}}\right)$

Stationary-Phase Method (cont.)

Hence

$$I \approx f(x_0) e^{j\Omega g(x_0)} \sqrt{\frac{2\pi}{\Omega |g''(x_0)|}} e^{\pm j\frac{\pi}{4}}$$

$$+, \quad g''(x_0) > 0$$

$$-, \quad g''(x_0) < 0$$

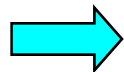
Example

$$J_0(\Omega) = \frac{1}{\pi} \int_0^\pi \cos(\Omega \sin \theta) d\theta$$
$$= \frac{1}{\pi} \operatorname{Re} \int_0^\pi e^{j\Omega \sin \theta} d\theta$$

$$f(\theta) = (1)$$

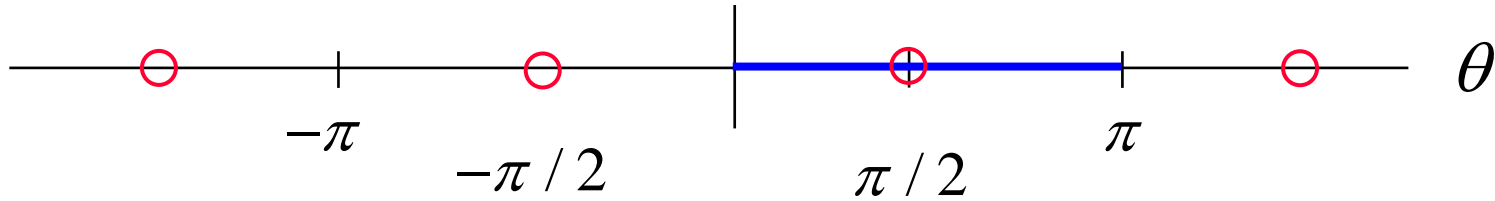
$$g(\theta) = \sin \theta$$

$$g'(\theta_0) = \cos \theta_0 = 0$$



$$\theta_0 = \frac{\pi}{2} + n\pi$$

Example (cont.)



$$g(\theta) = \sin \theta$$

$$g'(\theta) = \cos \theta$$

$$g''(\theta) = -\sin \theta$$

$$g''(\theta_0) = -\sin \theta_0 = -1 < 0$$

Hence

$$J_0(\Omega) \sim \frac{1}{\pi} \operatorname{Re} \left\{ e^{j\Omega(1)} \sqrt{\frac{2\pi}{\Omega|-1|}} e^{-j\frac{\pi}{4}} \right\}$$

Example (cont.)

$$J_0(\Omega) \sim \frac{1}{\pi} \operatorname{Re} \left\{ e^{j\Omega(1)} \sqrt{\frac{2\pi}{\Omega|-1|}} e^{-j\frac{\pi}{4}} \right\}$$

Hence

$$J_0(\Omega) \sim \sqrt{\frac{2}{\pi \Omega}} \cos\left(\Omega - \frac{\pi}{4}\right)$$

as $\Omega \rightarrow \infty$