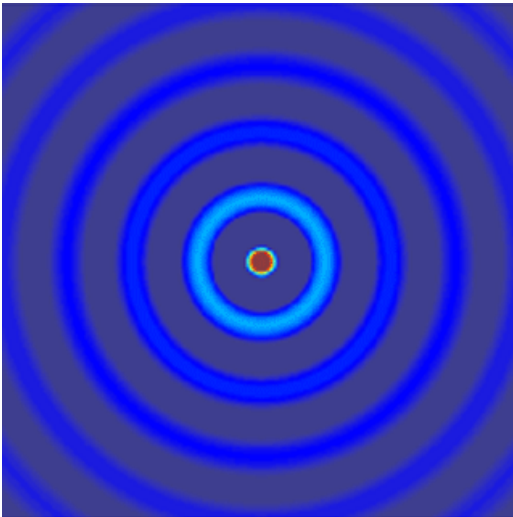


# ECE 6341

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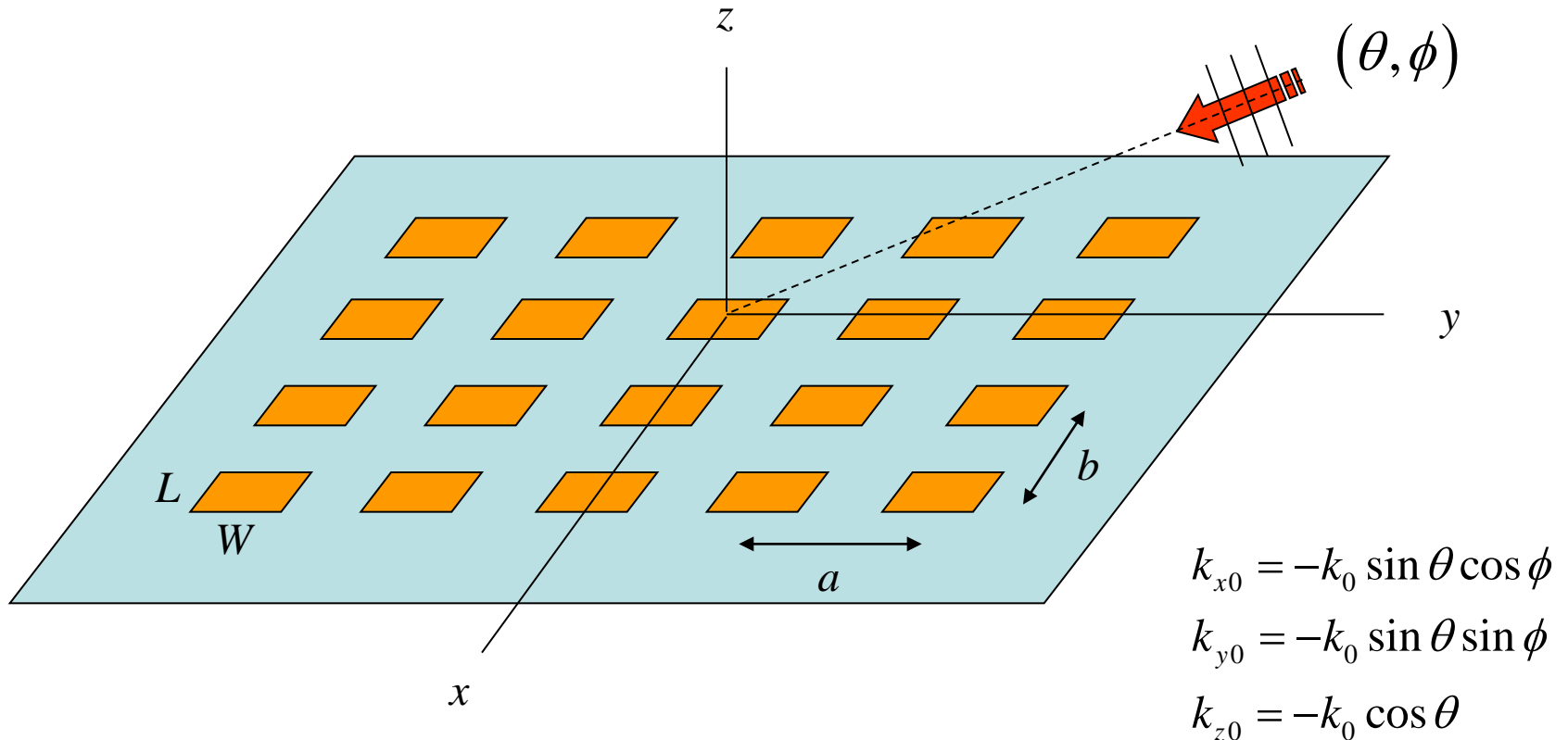


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# 2D Frequency Selective Surface (FSS)

Scattering from a 2-D array of metal patches (FSS structure)

$$\underline{E}^i(x, z) = \underline{E}_0 e^{-j(k_{x0}x + k_{y0}y + k_{z0}z)}$$



# 2D FSS (cont.)

$$\text{TM}_z: \quad \psi^i = A_z^i = e^{-j(k_{x0}x + k_{y0}y + k_{z0}z)}$$

$$\text{TE}_z: \quad \psi^i = F_z^i = e^{-j(k_{x0}x + k_{y0}y + k_{z0}z)}$$

Note: The **scattered** field will have both  $A_z$  and  $F_z$ , regardless of the polarization of the incident wave.

Denote  $\psi^s(x, y, 0) = e^{-j(k_{x0}x + k_{y0}y)} P(x, y)$

$P(x, y)$  = periodic function in  $x$  and  $y$

$$P(x + a, y) = P(x, y)$$

$$P(x, y + b) = P(x, y)$$

# 2D FSS (cont.)

Use a 2D Fourier series:

$$P(x, y) = \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} A_{mn} e^{-j\left(\frac{2\pi m}{a}\right)x} e^{-j\left(\frac{2\pi n}{b}\right)y}$$

Define:

$$k_{xm} = k_{x0} + \frac{2\pi m}{a}$$

$$k_{yn} = k_{y0} + \frac{2\pi n}{b}$$

# 2D FSS (cont.)

Define 2D Floquet waves:

$$\psi^+(x, y, z) = e^{-j(k_{xm}x + k_{yn}y)} e^{-jk_{zmn}z}$$

$$\psi^-(x, y, z) = e^{-j(k_{xm}x + k_{yn}y)} e^{+jk_{zmn}z}$$

where

$$k_{zmn} = \left( k_0^2 - k_{xm}^2 - k_{yn}^2 \right)^{1/2}$$

# 2D FSS (cont.)

$z > 0$ :

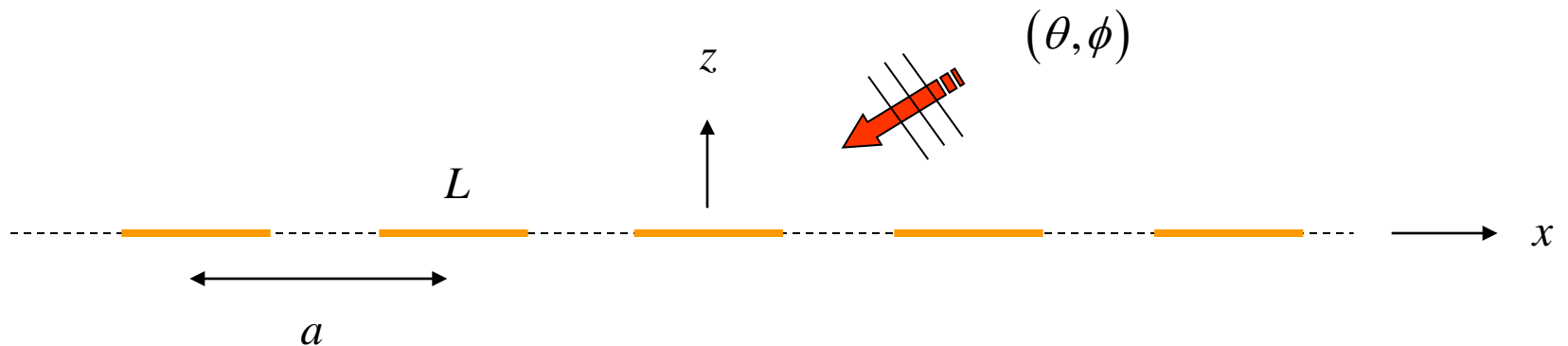
$$A_z^s(x, y, z) = \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} A_{mn} \psi_{mn}^+(x, y, z)$$

$$F_z^s(x, y, z) = \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} B_{mn} \psi_{mn}^+(x, y, z)$$

$z < 0$ :

$$A_z^s(x, y, z) = \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} C_{mn} \psi_{mn}^-(x, y, z)$$

$$F_z^s(x, y, z) = \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} D_{mn} \psi_{mn}^-(x, y, z)$$

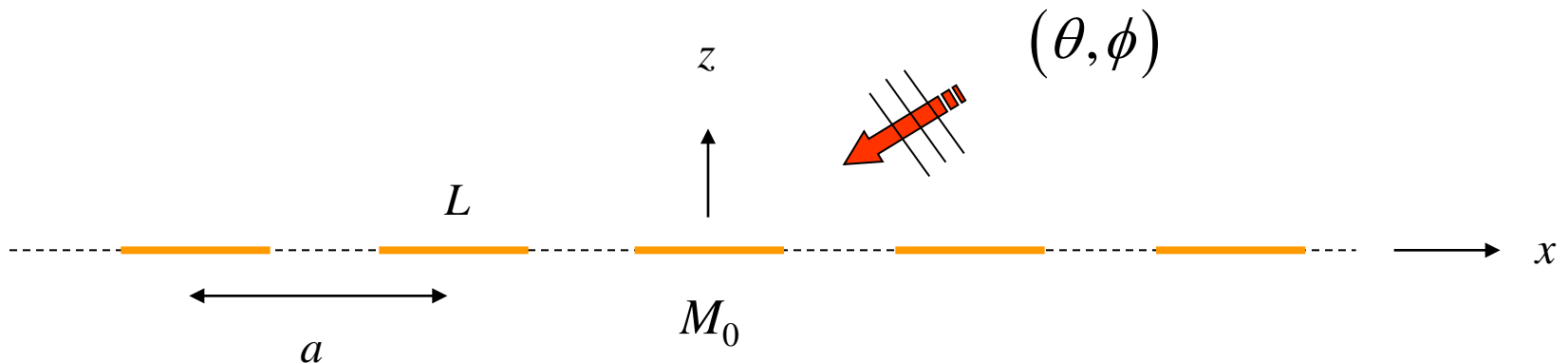


# 2D FSS (cont.)

Boundary condition (PEC patches):

$$\begin{aligned} E_x^s(x, y, 0) &= -E_x^i(x, y, 0) \\ E_y^s(x, y, 0) &= -E_y^i(x, y, 0) \end{aligned} \quad (x, y) \in M_0$$

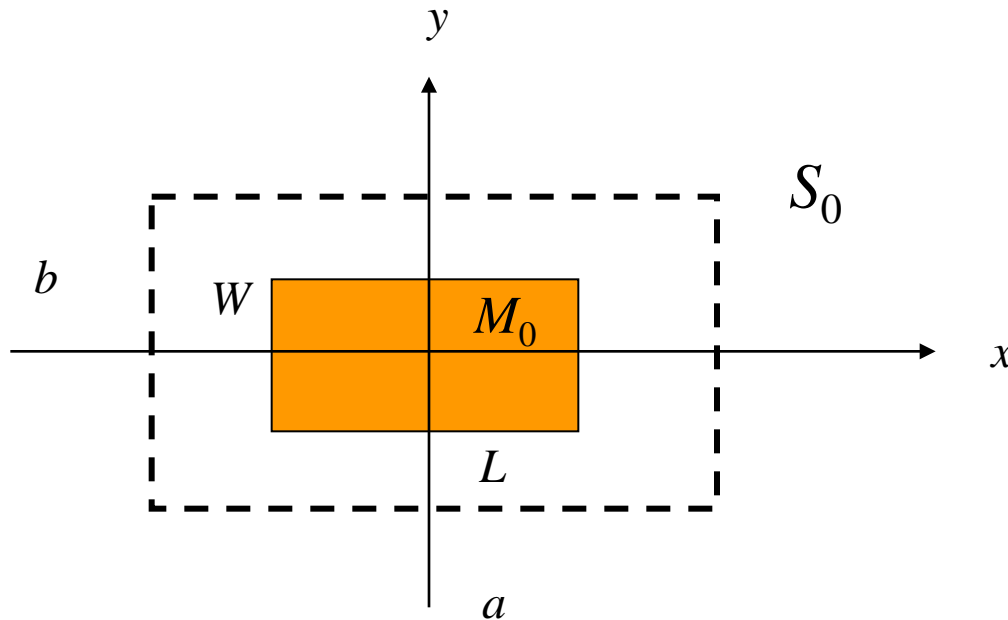
$M_0$  = metal surface of (0,0) patch



# 2D FSS (cont.)

Orthogonality of Floquet waves:

$$\int_{S_0} \psi_{mn}(x, y) \psi_{m'n'}^*(x, y) dx dy = \int_{S_0} e^{-j\left(\frac{2\pi}{a}\right)(m-m')x} e^{-j\left(\frac{2\pi}{b}\right)(n-n')y} dx dy$$
$$= \begin{cases} 0, & (m, n) \neq (m', n') \\ A, & (m, n) = (m', n') \end{cases}$$



$S_0 =$  unit cell

$A = ab =$  area of unit cell



# 2D FSS (cont.)

## OUTLINE OF SOLUTION STEPS

(1) Expand current on  $(0,0)$  patch in terms of basis functions:

$$\underline{J}_s(x, y) = \sum_{j=1}^N c_j \underline{B}_j(x, y) \quad (x, y) \in S_0$$

(2) Express the unknown Floquet coefficients  $A_{mn}$ ,  $B_{mn}$ ,  $C_{mn}$ ,  $D_{mn}$  in terms of the unknown current coefficients  $c_j$ , using the orthogonality of the Floquet waves over  $S_0$ .

- Enforce continuity of tangential electric field in  $S_0$ .
- Enforce jump condition on tangential magnetic field in  $S_0$ .

# 2D FSS (cont.)

## OUTLINE OF SOLUTION STEPS (cont.)

### (3) Enforce EFIE on $M_0$

- Use  $N$  testing functions.
- This leads to an  $N \times N$  system of linear equations for the unknown current coefficients  $c_j$ .

# Grating Waves

To avoid grating waves:

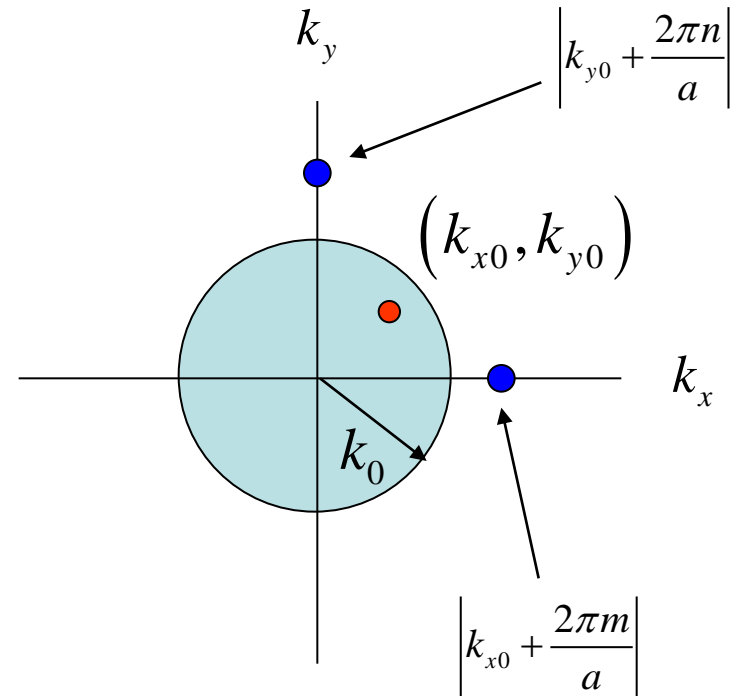
$$\left| k_{x0} + \frac{2\pi m}{a} \right| > k_0, \quad m \neq 0$$

$$\left| k_{y0} + \frac{2\pi n}{b} \right| > k_0, \quad n \neq 0$$

Set  $m = 1$ , in first equation,  
 $n = 1$  in second equation:

$$k_{x0} + \frac{2\pi}{a} > k_0$$

$$k_{y0} + \frac{2\pi}{b} > k_0$$



# Grating Waves (cont.)

Worst-case:

$$-k_0 + \frac{2\pi}{a} > k_0$$

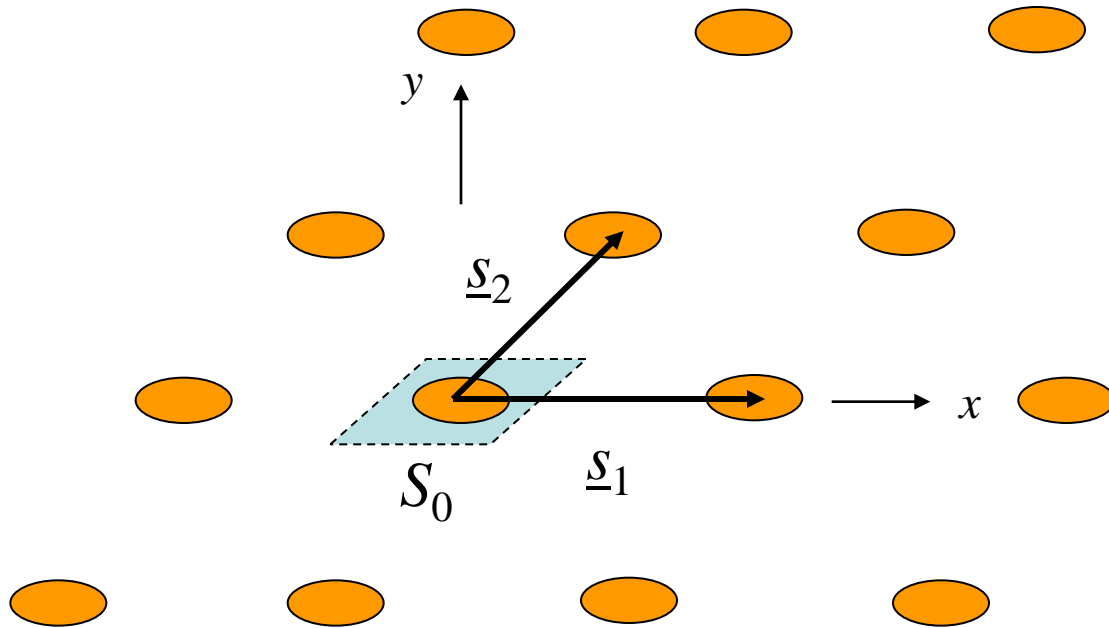
$$-k_0 + \frac{2\pi}{b} > k_0$$

This yields

$$\frac{a}{\lambda_0} < \frac{1}{2}$$

$$\frac{b}{\lambda_0} < \frac{1}{2}$$

# Skewed Lattice



Assume:

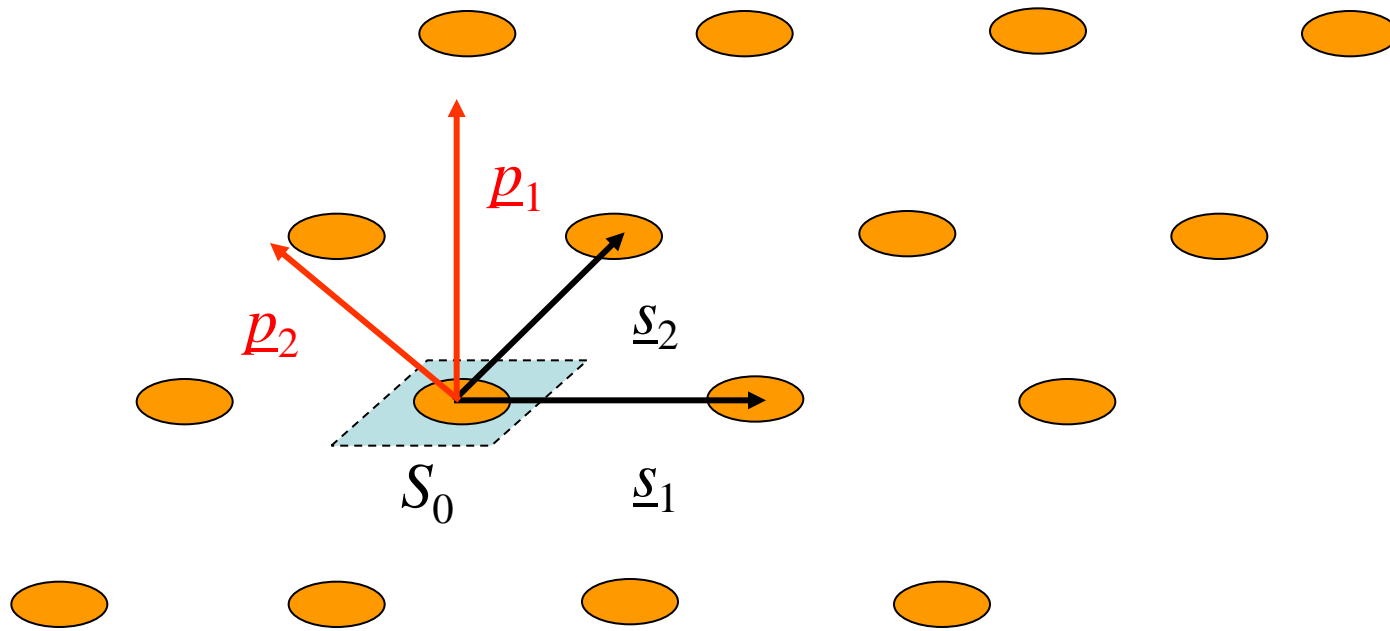
$$P_{mn}(\underline{r}) = e^{-j\mathbf{k}_{tmn}^p \cdot \underline{r}}$$

$$P_{mn}(\underline{r} + \underline{s}_1) = P_{mn}(\underline{r})$$

$$P_{mn}(\underline{r} + \underline{s}_2) = P_{mn}(\underline{r})$$

$$\begin{aligned} \psi^s(x, y, 0) &= e^{-j(k_{x0}x + k_{y0}y)} \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} P_{mn}(x, y) \\ &= e^{-j(\mathbf{k}_{tmn}^0 \cdot \underline{r})} \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} P_{mn}(x, y) \end{aligned}$$

# Skewed Lattice (cont.)



Define perpendicular vectors:

$$\underline{p}_1 = \hat{\underline{z}} \times \underline{s}_1$$

$$\underline{p}_2 = \hat{\underline{z}} \times \underline{s}_2$$

# Skewed Lattice (cont.)

$$P_{mn}(\underline{r}) = e^{-jk_{tmn}^p \cdot \underline{r}}$$

Assume  $\underline{k}_{tmn}^p = m(\alpha_1 \underline{p}_2) + n(\alpha_2 \underline{p}_1)$

Enforce:  $\underline{k}_{tmn}^p \cdot (\underline{r} + \underline{s}_1) = \underline{k}_{tmn}^p \cdot (\underline{r}) + 2\pi m$

$\longrightarrow \underline{k}_{tmn}^p \cdot (\underline{s}_1) = 2\pi m$

$\longrightarrow [m(\alpha_1 \underline{p}_2) + n(\alpha_2 \underline{p}_1)] \cdot (\underline{s}_1) = 2\pi m$

$\longrightarrow [m(\alpha_1 \underline{p}_2)] \cdot (\underline{s}_1) = 2\pi m$

# Skewed Lattice (cont.)

Hence we have

$$\alpha_1 = \frac{2\pi}{\underline{p}_2 \cdot \underline{s}_1}$$

Similarly, we have

$$\underline{k}_{tmn}^p \cdot (\underline{r} + \underline{s}_2) = \underline{k}_{tmn}^p \cdot (\underline{r}) + 2\pi n$$

$$\longrightarrow \underline{k}_{tmn}^p \cdot (\underline{s}_2) = 2\pi n$$

$$\longrightarrow \left[ m(\alpha_1 \underline{p}_2) + n(\alpha_2 \underline{p}_1) \right] \cdot (\underline{s}_2) = 2\pi n$$



# Skewed Lattice (cont.)

$$\left[ m(\alpha_1 \underline{p}_2) + n(\alpha_2 \underline{p}_1) \right] \cdot (\underline{s}_2) = 2\pi n$$

$$\longrightarrow \left[ n(\alpha_2 \underline{p}_1) \right] \cdot (\underline{s}_2) = 2\pi n$$

Hence, we have

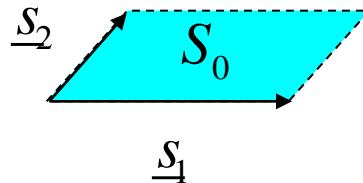
$$\alpha_2 = \frac{2\pi}{\underline{p}_1 \cdot \underline{s}_2}$$

# Skewed Lattice (cont.)

Simplify these results:

$$\alpha_1 = \frac{2\pi}{\underline{p}_2 \cdot \underline{s}_1}$$

$$\underline{p}_2 \cdot \underline{s}_1 = (\hat{z} \times \underline{s}_2) \cdot \underline{s}_1 = -\hat{z} \cdot (\underline{s}_1 \times \underline{s}_2) = -\hat{z} \cdot (\hat{z} |\underline{s}_1 \times \underline{s}_2|) = -\hat{z} \cdot (\hat{z} A) = -A$$



$A =$  area of unit cell

$$A = |\underline{s}_1 \times \underline{s}_2|$$

Hence the result is

$$\alpha_1 = -\frac{2\pi}{A}$$

# Skewed Lattice (cont.)

Similarly,

$$\alpha_2 = \frac{2\pi}{\underline{p}_1 \cdot \underline{s}_2}$$

$$\underline{p}_1 \cdot \underline{s}_2 = A$$

$$\alpha_2 = \frac{2\pi}{A}$$

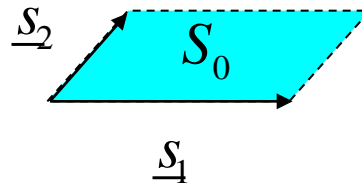
# Skewed Lattice (cont.)

Recall that  $\underline{k}_{tmn}^p = m(\alpha_1 \underline{p}_2) + n(\alpha_2 \underline{p}_1)$

$$\alpha_1 = -\frac{2\pi}{A} \quad \alpha_2 = \frac{2\pi}{A}$$

The wavenumbers for the periodic function are then

$$\underline{k}_{tmn}^p = \left(\frac{2\pi m}{A}\right)(-\underline{p}_2) + \left(\frac{2\pi n}{A}\right)(\underline{p}_1)$$



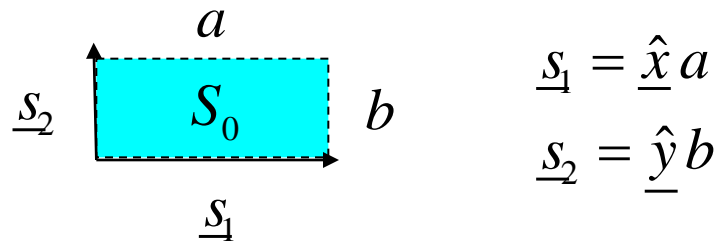
$$\underline{p}_1 = \hat{\underline{z}} \times \underline{s}_1$$

$$\underline{p}_2 = \hat{\underline{z}} \times \underline{s}_2$$

# Skewed Lattice (cont.)

Rectangular lattice:

$$\underline{k}_{tmn}^p = \left( \frac{2\pi m}{A} \right) (-\underline{p}_2) + \left( \frac{2\pi n}{A} \right) (\underline{p}_1)$$



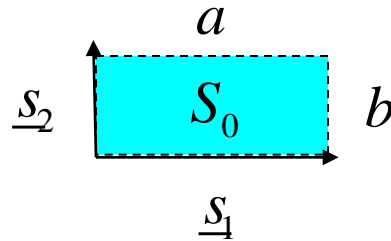
$$\underline{p}_1 = \hat{z} \times (\hat{x} a) = \hat{y} a$$

$$\underline{p}_2 = \hat{z} \times (\hat{y} b) = -\hat{x} b$$

$$\underline{k}_{tmn}^p = \left( \frac{2\pi m}{ab} \right) (\hat{x} b) + \left( \frac{2\pi n}{ab} \right) (\hat{y} a)$$

# Skewed Lattice (cont.)

Rectangular lattice (cont.):



$$\underline{k}_{tmn}^p = \left( \frac{2\pi m}{ab} \right) (\underline{\hat{x}}b) + \left( \frac{2\pi n}{ab} \right) (\underline{\hat{y}}a)$$

$$\underline{k}_{tmn}^p = \underline{\hat{x}} \left( \frac{2\pi m}{a} \right) + \underline{\hat{y}} \left( \frac{2\pi n}{b} \right)$$

$$e^{-j\underline{k}_{tmn}^p \cdot \underline{r}} = e^{-j \left( \frac{2\pi m}{a} \right) x} e^{-j \left( \frac{2\pi n}{b} \right) y}$$

# Skewed Lattice (cont.)

Orthogonality:

$$\int_{S_0} \psi_{mn}(x, y) \psi_{m'n'}^*(x, y) dx dy = \begin{cases} 0, & (m, n) \neq (m', n') \\ A, & (m, n) = (m', n') \end{cases}$$

(proof omitted)