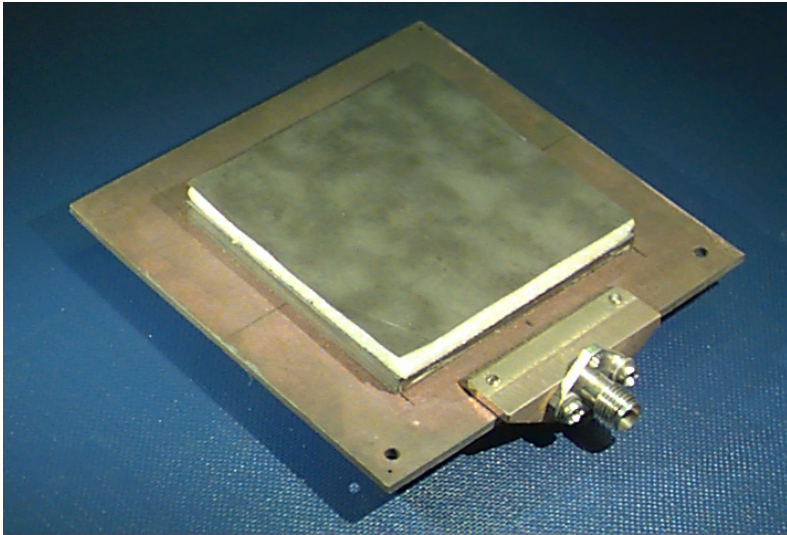


# ECE 6345

Spring 2015

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ECE Dept.



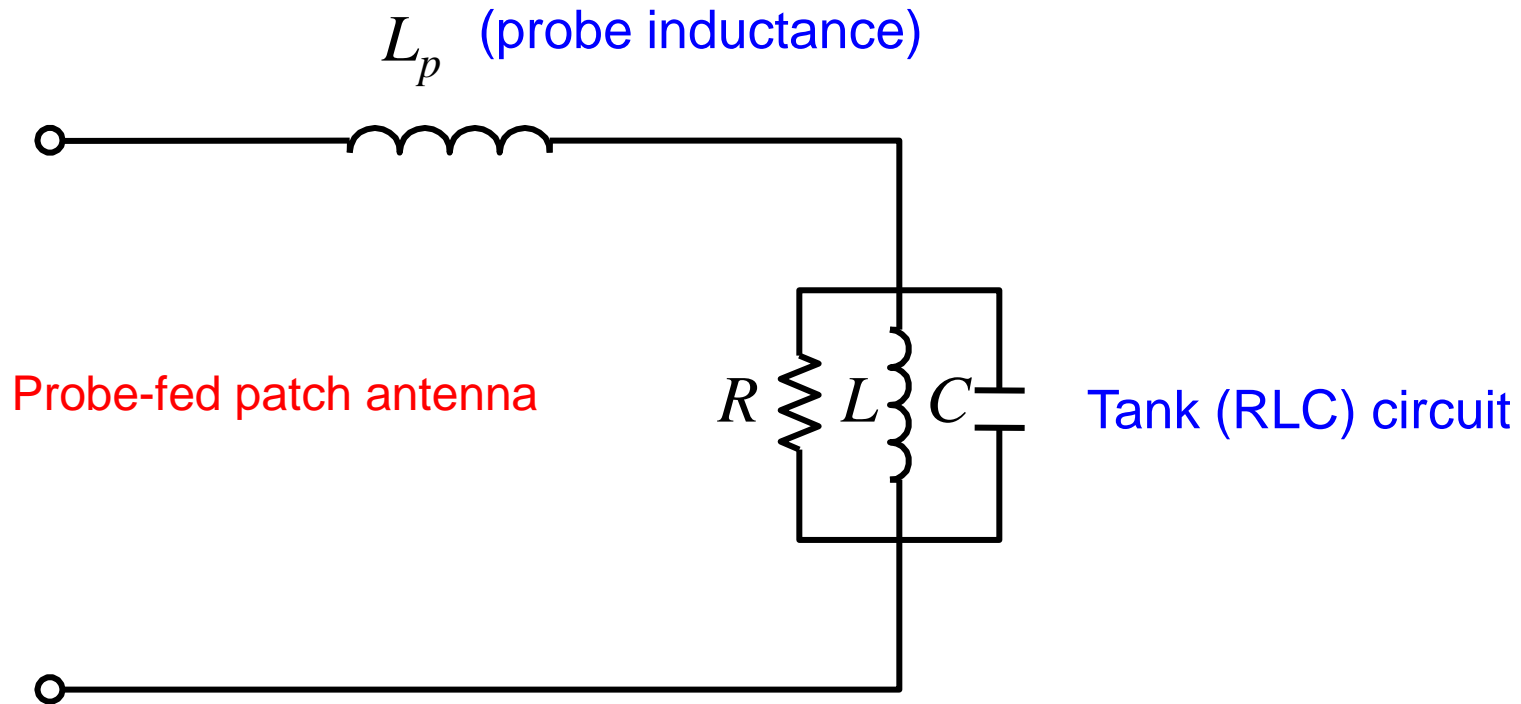
Notes 1

# Overview

In this set of notes we discuss the CAD model of the microstrip antenna.

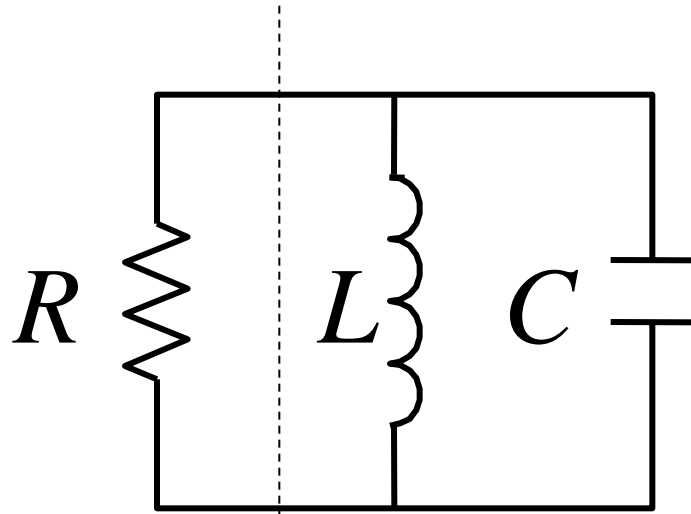
- Discuss complex resonance frequency
- Derive formula for  $Q$
- Derive formula for input impedance
- Derive formula for impedance bandwidth

# CAD Model of Microstrip Antennas



The circuit model is justified from the eigenfunction method in the cavity model, discussed later.

# Tank Circuit: complex resonance frequency



$$G \equiv \frac{1}{R}$$

Transverse Resonance Equation (TRE):

$$\overleftarrow{Y} \quad \leftarrow \quad \rightarrow \quad \overrightarrow{Y}$$

$$\overleftarrow{Y} = -\overrightarrow{Y}$$

The complex resonance frequency is denoted as  $\omega_0$ .

# Resonance Frequency (cont.)

TRE:

$$G = - \left[ j\omega_0 C + \frac{1}{j\omega_0 L} \right]$$

$$j\omega_0 LG = -1 + \omega_0^2 LC$$

$$\omega_0^2 (LC) + \omega_0 (-jLG) + (-1) = 0$$

$$\omega_0 = \frac{jLG \pm \sqrt{-L^2 G^2 + 4LC}}{2LC}$$

$$G \rightarrow 0, \quad \omega_0 \rightarrow \frac{1}{\sqrt{LC}} \quad \text{so choose + sign}$$

# Resonance Frequency (cont.)

$$\omega_0 = j \frac{1}{2} \left( \frac{1}{RC} \right) + \frac{1}{\sqrt{LC}} \sqrt{1 - \frac{1}{4} \frac{L}{R^2 C}}$$

Denote:  $\omega_0 = \omega'_0 + j\omega''_0$

$$\omega'_0 = \frac{1}{\sqrt{LC}} \sqrt{1 - \frac{1}{4} \frac{L}{R^2 C}}$$

$$\omega''_0 = \frac{1}{2} \left( \frac{1}{RC} \right)$$

# Resonance Frequency (cont.)

$$\omega'_0 = \frac{1}{\sqrt{LC}} \sqrt{1 - \frac{1}{4} \frac{L}{R^2 C}} \qquad \omega''_0 = \frac{1}{2} \left( \frac{1}{RC} \right)$$

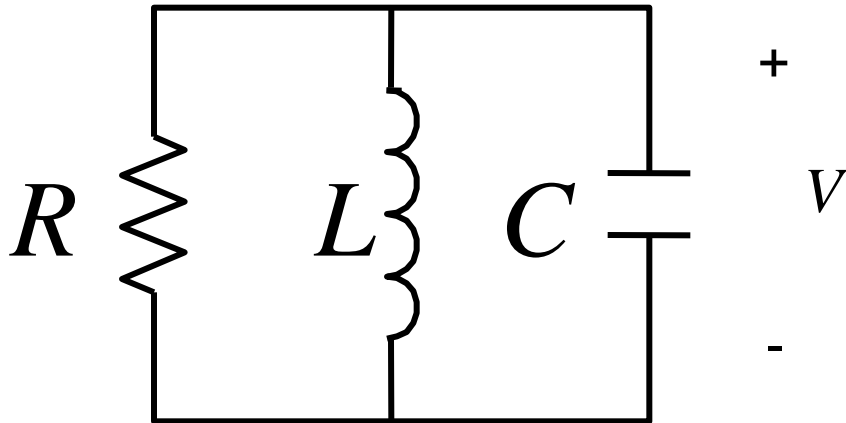
Assume  $R \gg \sqrt{\frac{L}{C}}$  (a good resonator)

We then have:

$$\omega'_0 \approx \frac{1}{\sqrt{LC}}$$

$$\omega''_0 = \frac{1}{2} \left( \frac{1}{RC} \right)$$

# Natural Response (no source)



The complex resonance frequency is  $\omega_0$ .

$$\omega_0 = \omega'_0 + j\omega''_0$$

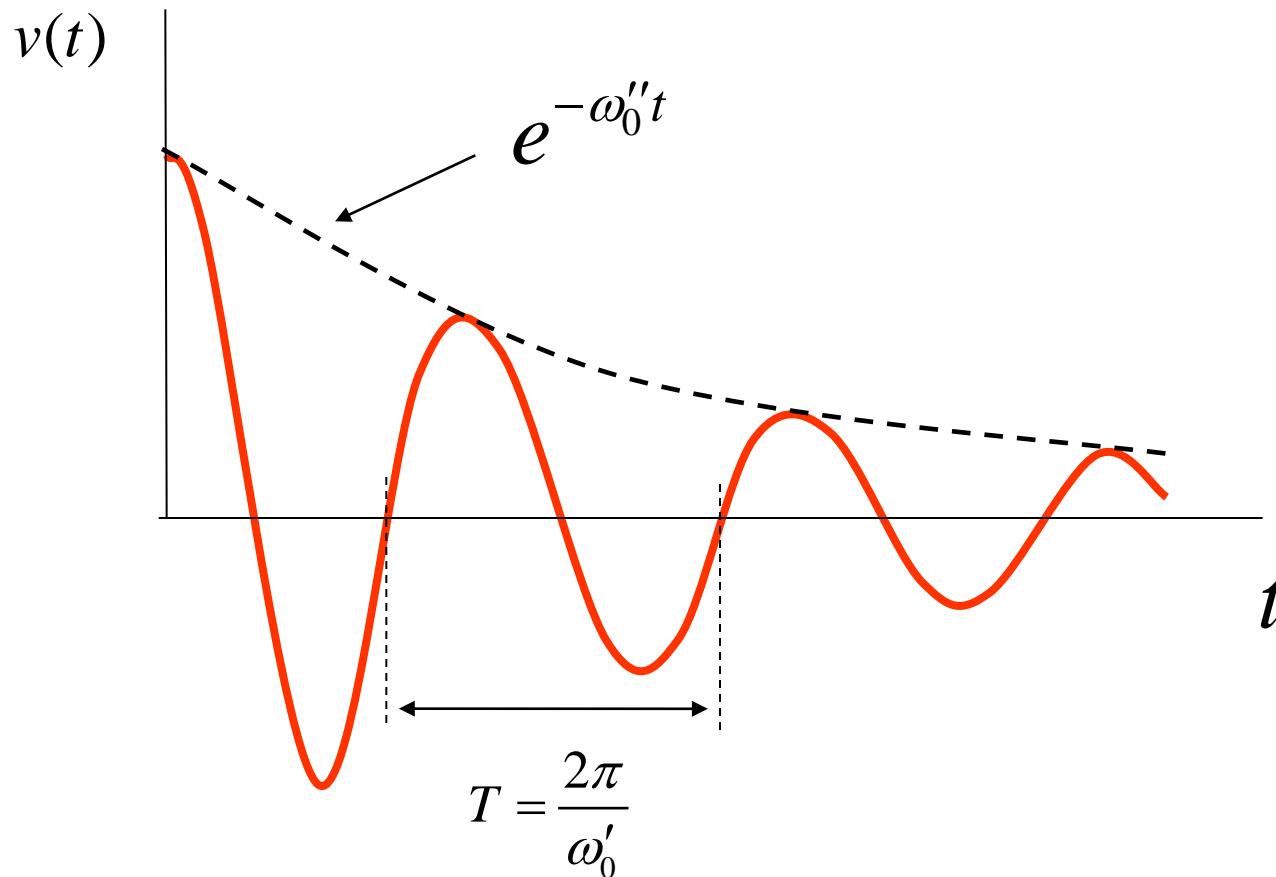
In the time domain:  $v(t) = \text{Re}\left(V e^{j\omega_0 t}\right)$  (Take  $V = 1$ )

so 
$$v(t) = \text{Re}\left(e^{j\omega'_0 t} e^{-\omega''_0 t}\right)$$

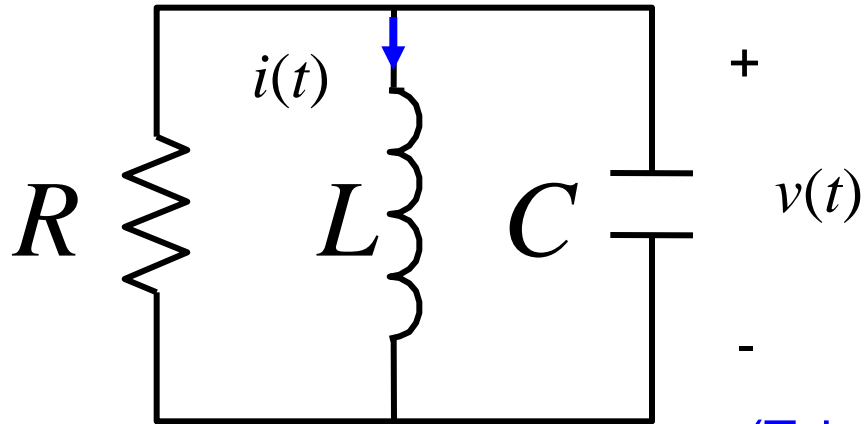


# Natural Response (cont.)

$$v(t) = e^{-\omega_0'' t} \cos(\omega_0' t)$$



# Stored Energy



For the capacitor:

$$U_E(t) = \frac{1}{2} C v^2(t)$$

$$= \frac{1}{2} C e^{-2\omega_0''t} \cos^2(\omega_0't)$$

(Take  $V = 1$ )

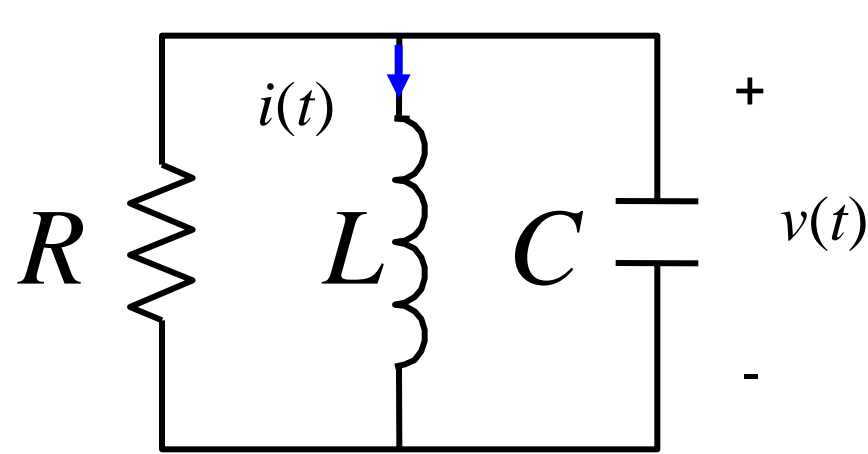
For the inductor:

$$I = \frac{V}{j\omega_0 L} = \frac{1}{j\omega_0 L} \approx \frac{1}{j\omega_0' L}$$

Therefore,

$$i(t) = \text{Re}\left(I e^{j\omega_0 t}\right) \approx \frac{1}{\omega_0' L} \text{Re}\left(\frac{1}{j} e^{-\omega_0'' t} e^{+j\omega_0' t}\right) = \frac{1}{\omega_0' L} e^{-\omega_0'' t} \sin(\omega_0' t)$$

# Stored Energy (cont.)



$$\begin{aligned}U_H(t) &= \frac{1}{2} L i^2(t) \\&= \frac{1}{2} L \left( \frac{1}{\omega'_0 L} \right)^2 e^{-2\omega''_0 t} \sin^2(\omega'_0 t) \\&= \frac{1}{2} C e^{-2\omega''_0 t} \sin^2(\omega'_0 t)\end{aligned}$$

Note:  $\langle U_E(t) \rangle = \langle U_H(t) \rangle = \frac{1}{4} C e^{-2\omega''_0 t}$

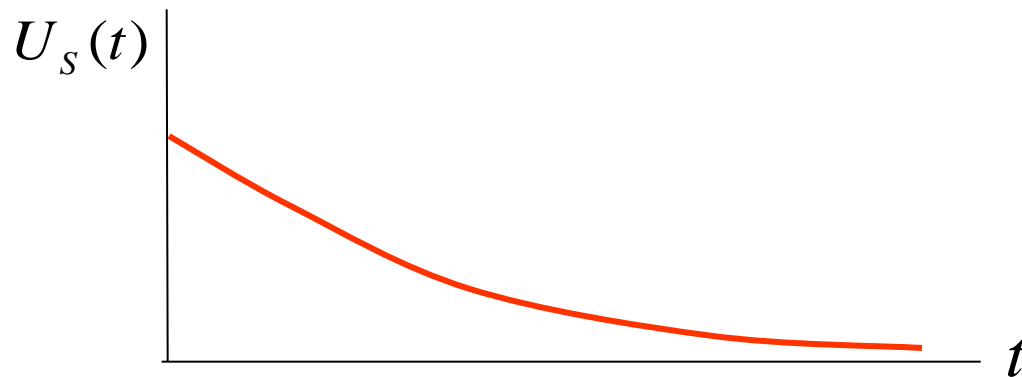
Also, note that  $U_S(t) = U_E(t) + U_H(t) = \frac{1}{2} C e^{-2\omega''_0 t}$

# Stored Energy (cont.)

$$U_S(t) = \langle U_S(t) \rangle = \langle U_E(t) \rangle + \langle U_H(t) \rangle = 2 \langle U_E(t) \rangle$$

Hence

$$U_S(t) = \frac{1}{2} C e^{-2\omega_0'' t} = U_S(0) e^{-2\omega_0'' t}$$



# Q of Cavity

$$Q \equiv 2\pi \left( \frac{U_S}{U_D^T} \right)$$

$U_D^T$  = energy dissipated per cycle (period  $T$ )

$$Q \equiv \frac{2\pi}{T} \left( \frac{U_S}{U_D^T / T} \right) \quad \text{or} \quad Q \equiv \omega'_0 \left( \frac{U_S}{P_D^{AVE}} \right)$$

$P_D^{AVE}$  = average power dissipated

(this includes radiation loss)

# Q of Cavity (cont.)

$$\begin{aligned} Q &= \omega'_0 \left( \frac{\frac{1}{2} C e^{-2\omega''_0 t}}{\langle G v(t)^2 \rangle} \right) = \omega'_0 \left( \frac{\frac{1}{2} C e^{-2\omega''_0 t}}{G \langle e^{-2\omega''_0 t} \cos^2(\omega'_0 t) \rangle} \right) \\ &= \omega'_0 \left( \frac{\frac{1}{2} C e^{-2\omega''_0 t}}{\frac{1}{2} G e^{-2\omega''_0 t}} \right) \\ &= \frac{\omega'_0 C}{G} \\ &= \omega'_0 RC \end{aligned}$$

Note:

$$\begin{aligned} &\langle e^{-2\omega''_0 t} \cos^2(\omega'_0 t) \rangle \\ &\approx e^{-2\omega''_0 t} \langle \cos^2(\omega'_0 t) \rangle \\ &= e^{-2\omega''_0 t} \left( \frac{1}{2} \right) \end{aligned}$$

# $Q$ of Cavity (cont.)

We then have

$$Q = \omega'_0 RC = \frac{1}{\omega'_0} \left( \frac{R}{L} \right) = R \sqrt{\frac{C}{L}}$$

Recall that

$$\omega'_0 \approx \frac{1}{\sqrt{LC}} \quad \omega''_0 = \frac{1}{2} \left( \frac{1}{RC} \right)$$

Hence

$$\omega_0 \approx \omega'_0 \left( 1 + j \frac{1}{2} \left( \frac{1}{RC} \right) \sqrt{LC} \right)$$

# $Q$ of Cavity (cont.)

$$\begin{aligned}\omega &\approx \omega'_0 \left( 1 + j \frac{1}{2} \left( \frac{1}{RC} \right) \sqrt{LC} \right) \\ &= \omega'_0 \left( 1 + j \frac{1}{2} \left( \frac{1}{R} \right) \sqrt{\frac{L}{C}} \right)\end{aligned}$$

Hence

$$\omega_0 \approx \omega'_0 \left( 1 + j \frac{1}{2Q} \right)$$

so

$$Q = \frac{1}{2} \frac{\omega'_0}{\omega''_0}$$



# $Q$ of Cavity (Cont.)

We can thus write

$$v(t) = e^{-\left(\frac{\omega'_0}{2Q}\right)t} \cos(\omega'_0 t)$$

$$U_S(t) = U_S(0) e^{-\left(\frac{\omega'_0}{Q}\right)t}$$

# Input Impedance

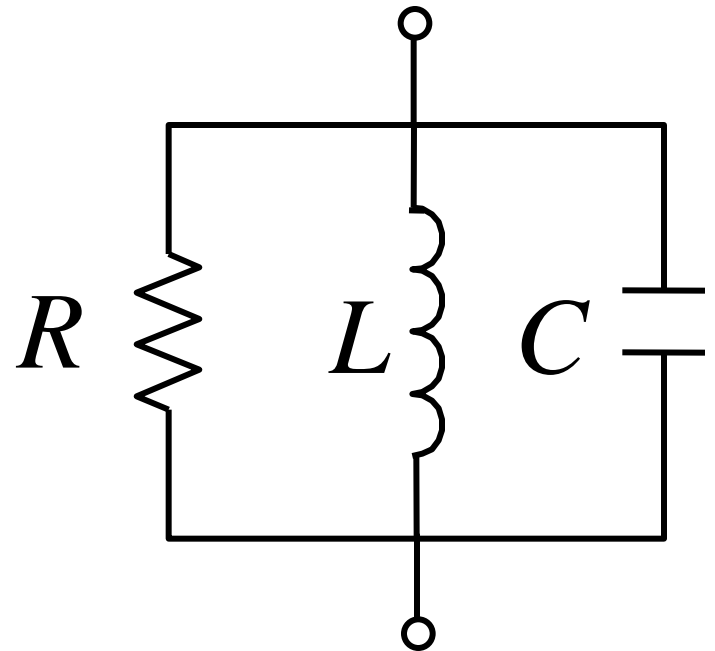
$$Y_{RLC} = G + j\omega C + \frac{1}{j\omega L}$$

$$Z_{RLC} = \frac{1}{G + j\omega C + \frac{1}{j\omega L}}$$

$$= \frac{R}{1 + j\omega RC + \frac{R}{j\omega L}}$$

$$= \frac{R}{1 + j\left(\omega RC - \frac{R}{\omega L}\right)}$$

The probe inductance is neglected here.



# Input Impedance (cont.)

or

$$\begin{aligned} Z_{RLC} &= \frac{R}{1 + j \left( \frac{\omega}{\omega'_0} (\omega'_0 RC) - \frac{R}{\omega'_0 L} \left( \frac{\omega'_0}{\omega} \right) \right)} \\ &= \frac{R}{1 + j \left( \frac{\omega}{\omega'_0} Q - Q \left( \frac{\omega'_0}{\omega} \right) \right)} \end{aligned}$$

Define

$$f_r \equiv \frac{f}{f_0} = \frac{\omega}{\omega'_0}$$

where

$$f_0 \equiv \frac{\omega'_0}{2\pi}$$

(real resonance frequency)

Then we have:

$$Z_{RCL} = \frac{R}{1 + jQ \left( f_r - \frac{1}{f_r} \right)}$$

# Input Impedance (cont.)

Define:

$$F \equiv f_r - \frac{1}{f_r}$$

$$\begin{aligned} F &= \frac{1}{f_r} (f_r^2 - 1) \\ &= \frac{1}{f_r} (f_r - 1)(f_r + 1) \\ &\approx 2(f_r - 1) \text{ for } f_r \approx 1 \end{aligned}$$

Hence, we have

$$Z_{RLC} = \frac{R}{1 + jQF} \approx \frac{R}{1 + j2Q(f_r - 1)}$$

# Input Impedance (cont.)

$$Z_{RLC} = R \left[ \frac{1}{1 + jQF} \right]$$

Define:  $x \equiv QF = Q \left( f_r - \frac{1}{f_r} \right) \approx 2Q(f_r - 1)$

$$\bar{Z}_{RLC} \equiv \frac{Z_{RLC}}{R}$$

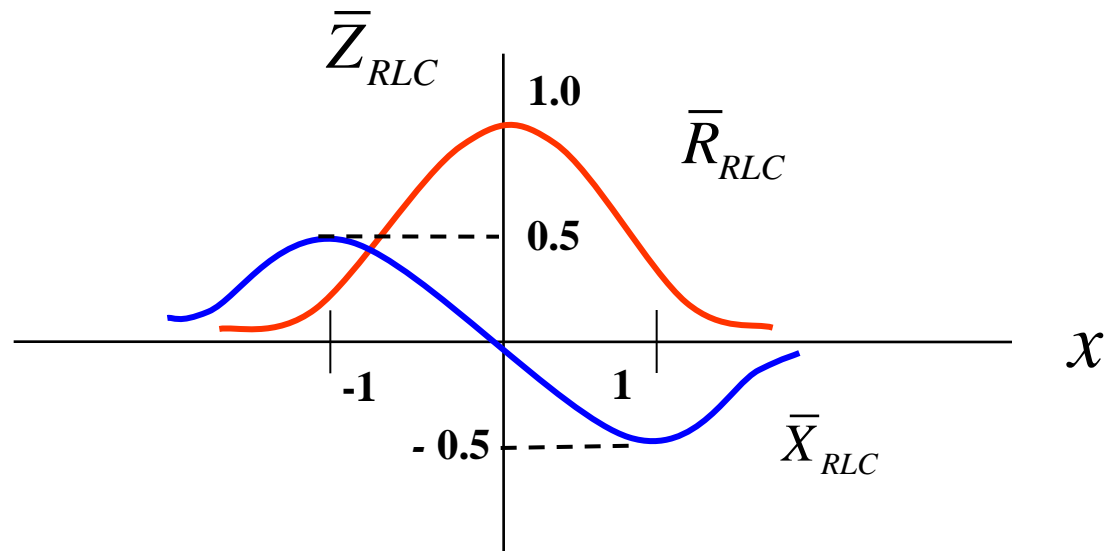
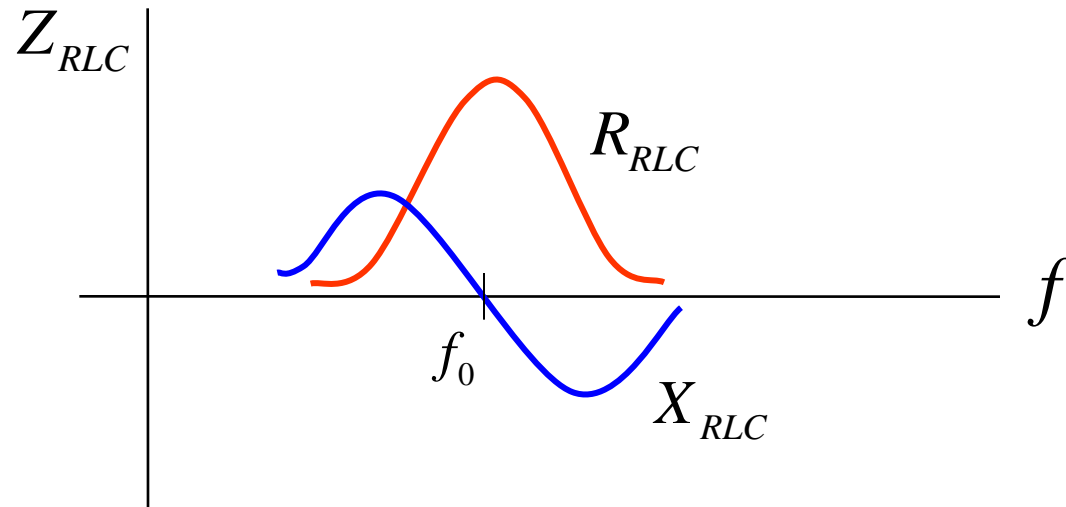
We then have

$$\bar{Z}_{RLC} = \frac{1}{1 + jx}$$

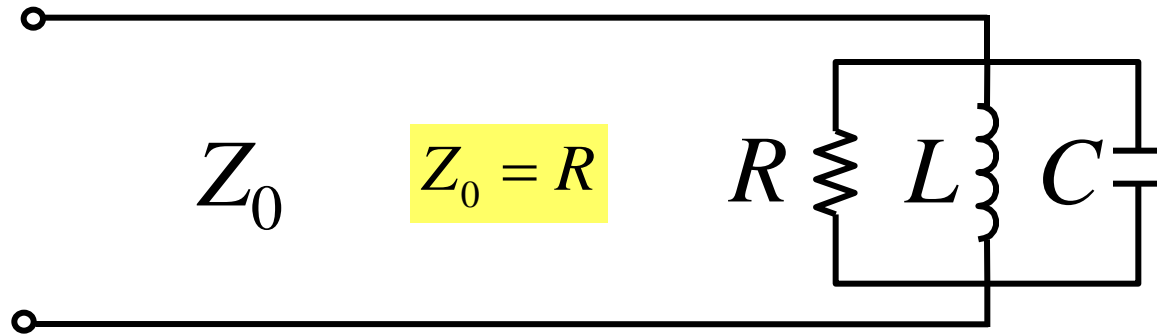
$$\bar{R}_{RLC} = \frac{1}{1 + x^2}$$

$$\bar{X}_{RLC} = \frac{-x}{1 + x^2}$$

# Input Impedance (cont.)



# Reflection Coefficient



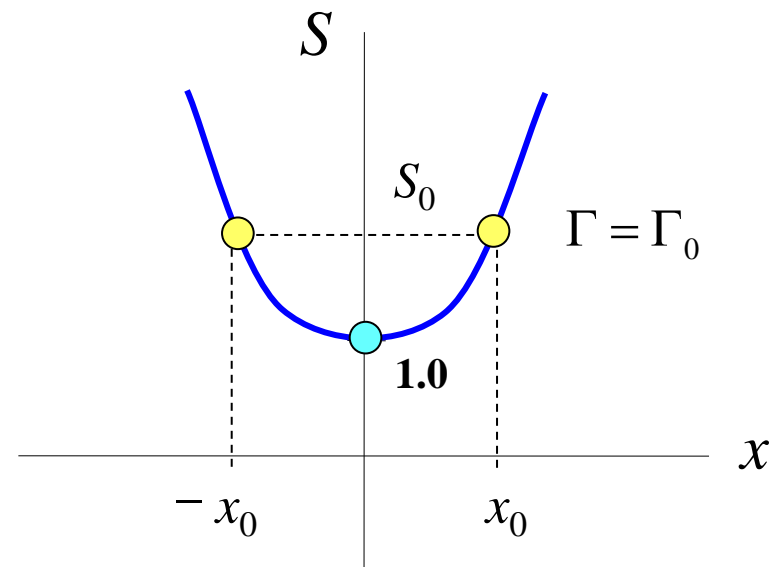
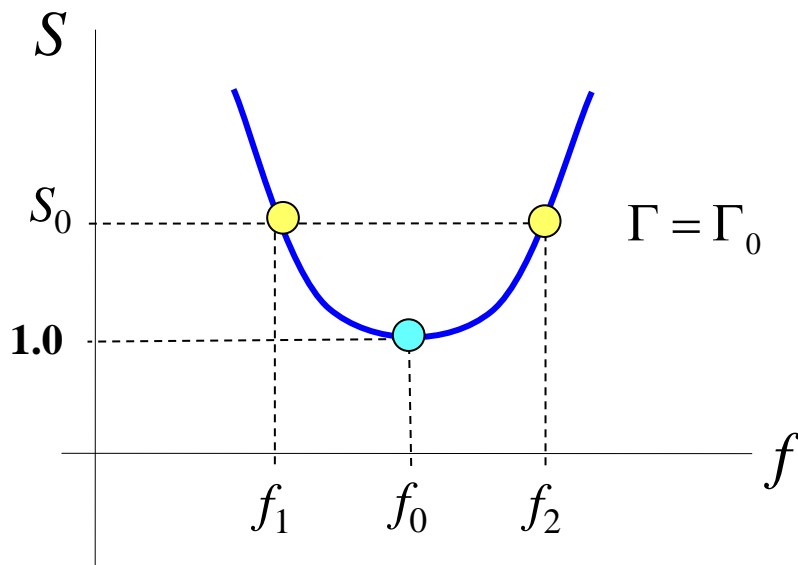
$$\begin{aligned}\Gamma &= \frac{Z_{RLC} - Z_0}{Z_{RLC} + Z_0} = \frac{\bar{Z}_{RLC} - 1}{\bar{Z}_{RLC} + 1} \\ &= \frac{1 - \bar{Y}_{RLC}}{1 + \bar{Y}_{RLC}} = \frac{1 - (1 + jx)}{1 + (1 + jx)} \\ &= \frac{-jx}{2 + jx}\end{aligned}$$

# Bandwidth

$$|\Gamma| = \frac{|x|}{\sqrt{4 + x^2}} \qquad S = SWR = \frac{1 + |\Gamma|}{1 - |\Gamma|}$$

Bandwidth definition is based on  $SWR < S_0$

The value  $S_0$  is often chosen as 2.0.)





# Bandwidth (cont.)

Fractional bandwidth:  $BW = \frac{f_2 - f_1}{f_0} = f_{r2} - f_{r1}$

Recall that  $x \equiv QF = Q \left( f_r - \frac{1}{f_r} \right)$

We can solve for  $f_r$  in terms of  $x$ :

$$f_r - \frac{1}{f_r} = \frac{x}{Q} \quad \Rightarrow \quad f_r^2 - f_r \left( \frac{x}{Q} \right) - 1 = 0$$

so

$$f_r = \frac{\frac{x}{Q} \pm \sqrt{\frac{x^2}{Q^2} + 4}}{2}$$

# Bandwidth (cont.)

To determine correct sign, enforce that  $x \rightarrow 0, f_r \rightarrow 1$

(So choose the plus sign.)

Hence

$$f_r = \frac{\frac{x}{Q} + \sqrt{\frac{x^2}{Q^2} + 4}}{2}$$

Therefore

$$f_{r2} = \frac{\frac{x_0}{Q} + \sqrt{\frac{x_0^2}{Q^2} + 4}}{2}$$

$$f_{r1} = \frac{-\frac{x_0}{Q} + \sqrt{\frac{x_0^2}{Q^2} + 4}}{2}$$

# Bandwidth (cont.)

$$\text{Hence, } BW = \frac{x_0}{Q}$$

Now we need to solve for  $x_0$ :

$$S_0 = \frac{1 + |\Gamma_0|}{1 - |\Gamma_0|} \quad \text{so} \quad |\Gamma_0| = \frac{S_0 - 1}{S_0 + 1}$$

$$\text{Also, } |\Gamma| = \frac{|x|}{\sqrt{4 + x^2}} \Rightarrow |\Gamma_0| = \frac{|x_0|}{\sqrt{4 + x_0^2}} = \frac{x_0}{\sqrt{4 + x_0^2}}$$

# Bandwidth (cont.)

Therefore 
$$\frac{x_0}{\sqrt{4 + x_0^2}} = \frac{S_0 - 1}{S_0 + 1}$$

so

$$\frac{x_0^2}{4 + x_0^2} = \left( \frac{S_0 - 1}{S_0 + 1} \right)^2 \equiv A$$

Thus we have 
$$4A + x_0^2 A = x_0^2$$

or

$$x_0^2 (1 - A) = 4A$$

# Bandwidth (cont.)

The solution is:

$$\begin{aligned}x_0 &= 2\sqrt{\frac{A}{1-A}} \\&= 2\frac{\left(\frac{S_0 - 1}{S_0 + 1}\right)}{\sqrt{1 - \left(\frac{S_0 - 1}{S_0 + 1}\right)^2}} \\&= 2\frac{S_0 - 1}{\sqrt{(S_0 + 1)^2 - (S_0 - 1)^2}} \\&= 2\left(\frac{S_0 - 1}{\sqrt{4S_0}}\right)\end{aligned}$$

# Bandwidth (cont.)

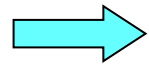
Hence,  $x_0 = \frac{S_0 - 1}{\sqrt{S_0}}$

We then have  $BW = \frac{1}{Q} \left( \frac{S_0 - 1}{\sqrt{S_0}} \right)$

For  $S_0 = 2$  we have:  $BW = \frac{1}{\sqrt{2}Q}$

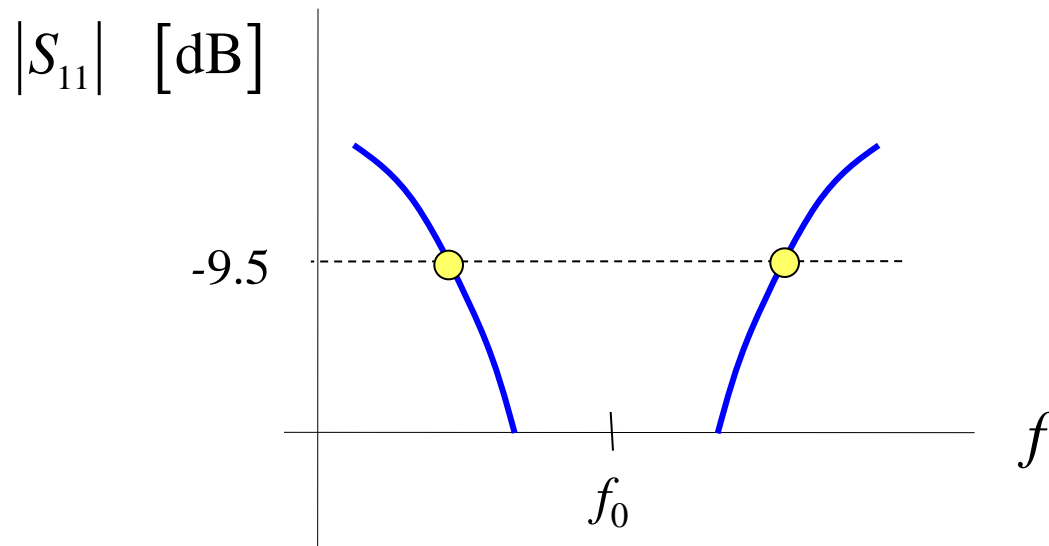
# Bandwidth (cont.)

Note:  $S_0 = SWR = 2.0$

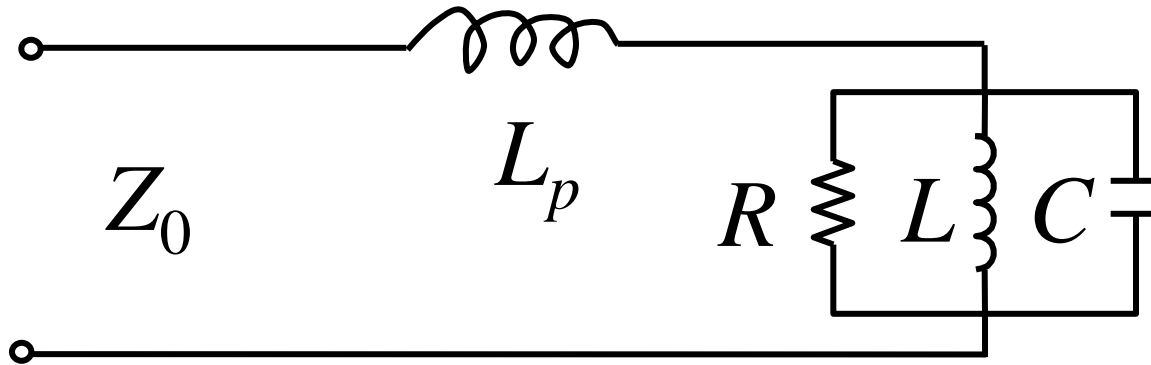


$$|\Gamma| = |S_{11}| = \frac{1}{3}$$

$$20 \log_{10} |S_{11}| = -9.5 \text{ [dB]}$$



# Complete Model

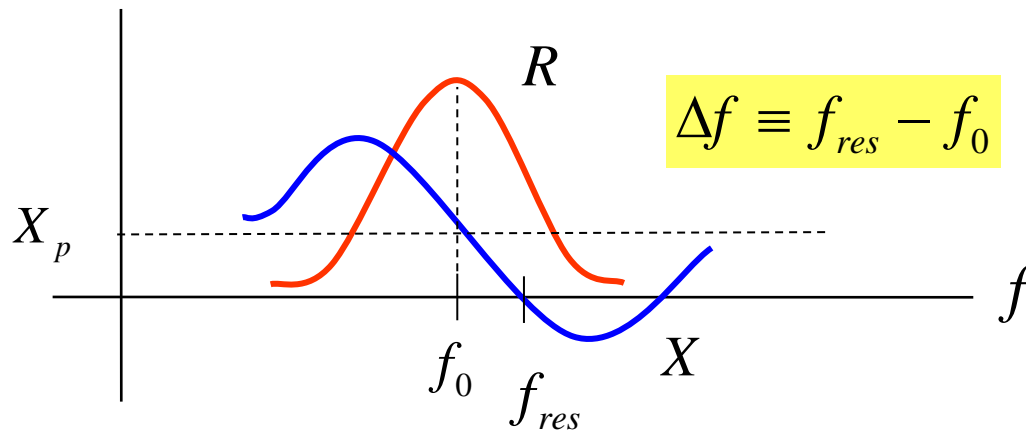


$$X_p = \omega L_p \approx \omega'_0 L_p$$

$$Z_{in} = jX_p + Z_{RLC}$$

$$\frac{\Delta f}{f_0} \approx (BW) \left( \frac{1}{\sqrt{2}} \right) \left( \frac{X_p}{R} \right)$$

(This will be derived in a HW problem.)



$$\Delta f \equiv f_{res} - f_0$$



# Complete Model (cont.)

Define  $\bar{X}_p \equiv X_p / R$

In terms of the normalized variable  $x$ , the resonance frequency  $f_{res}$  where the input impedance is purely real, corresponds to

$$x_{res} = \frac{1 - \sqrt{1 - 4\bar{X}_p^2}}{2\bar{X}_p}$$

(This will be derived in a HW problem.)

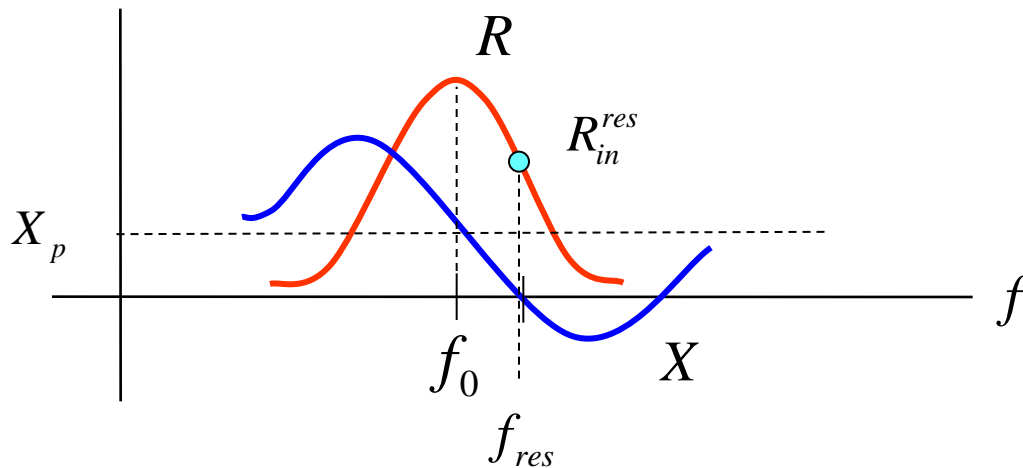
If  $X_p \ll R$  then  $x_{res} \approx \bar{X}_p$

(This follows from a binomial expansion of the square-root term in the numerator.)

# Complete Model (cont.)

At the resonance frequency, the input resistance is then

$$R_{in}^{res} = \frac{R}{1 + x_{res}^2} \quad \Rightarrow \quad R_{in}^{res} \approx \frac{R}{1 + \left(\frac{X_p}{R}\right)^2}$$



# Complete Model (cont.)

$$R_{in}^{res} \approx \frac{R}{1 + \left(\frac{X_p}{R}\right)^2}$$

Note that the probe reactance changes the input resistance at resonance.

Given a specified value of the input resistance at resonance (e.g.,  $R_{in}^{res} = 50 \Omega$ ), we wish to solve for the corresponding value of  $R$ .

Note that the CAD formula for resonant input resistance (in the short-course notes) gives us the value of  $R$  in terms of the feed location.

# Complete Model (cont.)

To solve for  $R$ , use

$$R = R_{in}^{res} + R_{in}^{res} \left( \frac{X_p^2}{R^2} \right)$$

and solve iteratively:

$$R^{(i)} = R_{in}^{res} + R_{in}^{res} \left( \frac{X_p^2}{(R^{(i-1)})^2} \right)$$

Zero iteration:

$$R = R_{in}^{res}$$

First iteration:

$$R = R_{in}^{res} + \left( \frac{X_p^2}{R_{in}^{res}} \right)$$

Second iteration:

$$R = R_{in}^{res} + R_{in}^{res} \left( \frac{X_p^2}{\left( R_{in}^{res} + \frac{X_p^2}{R_{in}^{res}} \right)^2} \right)$$