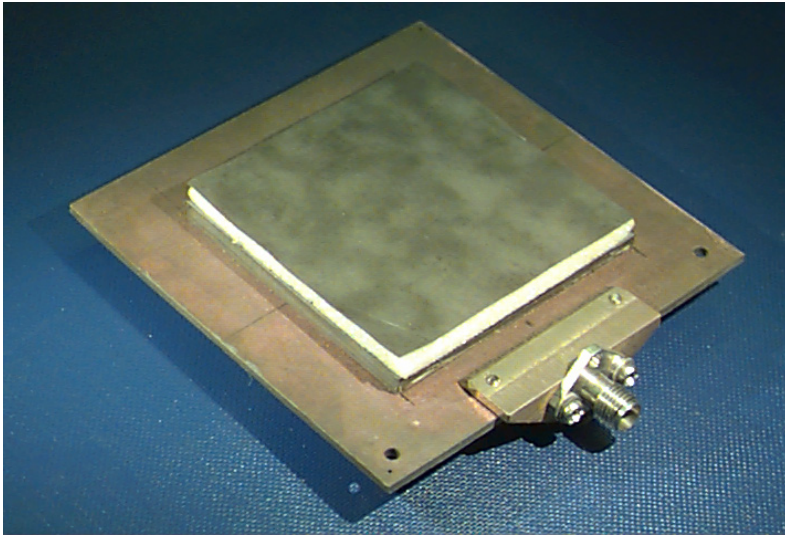


ECE 6345

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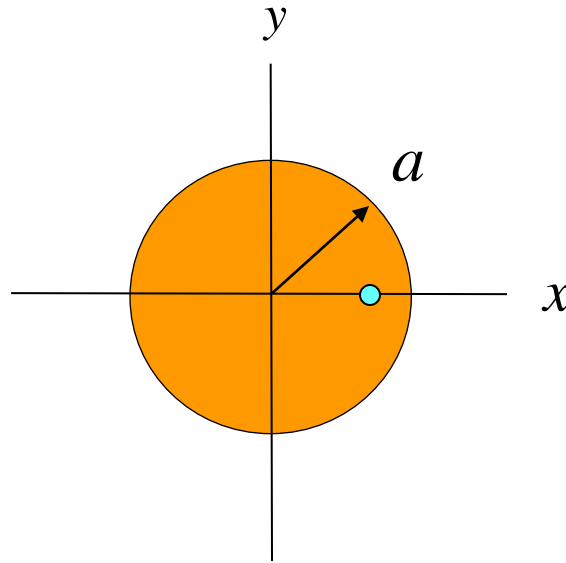
Notes 10

Overview

In this set of notes we derive the far-field pattern of a circular patch operating in the dominant TM_{11} mode.

We use the magnetic current model.

Circular Patch: TM_{11} Mode



$$E_z(\rho, \phi) = A \cos \phi J_1(k\rho)$$

This corresponds to a probe on the x axis.

$$k = \frac{x'_{11}}{a}$$

$$x'_{11} = 1.841$$

Circular Patch (cont.)

Magnetic current model:

$$\begin{aligned}\underline{M}_s &= -\underline{\hat{n}} \times \underline{\hat{E}} \\ &= -\underline{\hat{\rho}} \times \underline{\hat{E}}\end{aligned}$$

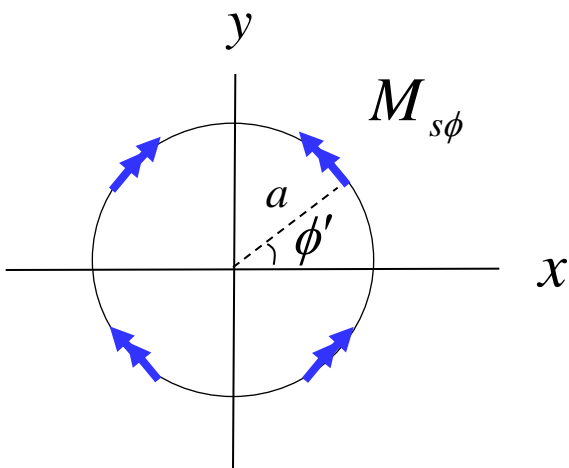
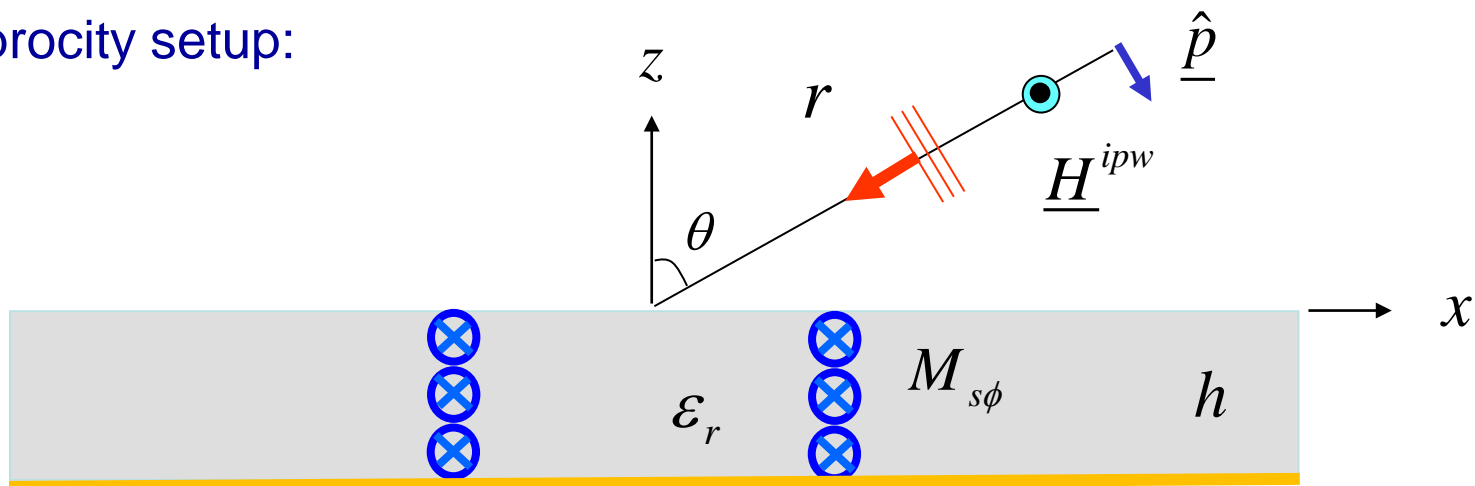
$$M_{s\phi} = E_z(a, \phi) = A \cos \phi J_1(ka)$$

Choose $A = \frac{1}{J_1(ka)}$ $M_{s\phi} = \cos \phi$

($V_0 = -h$ [V] at patch edge on x axis)

Far Field of Circular Patch

Reciprocity setup:



Far Field of Circular Patch (cont.)

Far-field:

$$\begin{aligned} E_p^{FF}(r, \theta, \phi) &= \langle a, b \rangle \\ &= \langle b, a \rangle \\ &= - \int_S \underline{H}^{pw} \cdot \underline{M}_s^a dS \\ &= - \int_S H_{\phi'}^{pw}(\rho', \phi', z') \cos \phi' dS' \\ &= - \int_0^{2\pi} \int_{-h}^0 H_{\phi'}^{pw}(a, \phi', z') \cos \phi' a dz' d\phi' \end{aligned}$$

The primes here denotes source coordinates.

$$E_p^{FF}(r, \theta, \phi) = - \int_0^{2\pi} \int_{-h}^0 H_{\phi'}^{pw}(a, \phi', z') \cos \phi' a dz' d\phi'$$

$$H_{\phi'}^{pw} = H_x^{pw}(-\sin \phi') + H_y^{pw}(\cos \phi')$$

Far Field of Circular Patch (cont.)

Inside the substrate we have (see Notes 9):

$$H_{x,y}^{pw}(x', y', z') = H_{x,y}^{pw}(0, 0, 0) e^{j(k_x x' + k_y y')} \left[\sec(k_{z1} h) \cos(k_{z1} (z' + h)) \right]$$

$$(k_{z1} = k_0 N_1(\theta))$$

The exponent term may be put in cylindrical coordinates as follows:

$$\begin{aligned} k_x x' + k_y y' &= (k_0 \sin \theta \cos \phi)(a \cos \phi') + (k_0 \sin \theta \sin \phi)(a \sin \phi') \\ &= k_0 a \sin \theta (\cos \phi \cos \phi' + \sin \phi \sin \phi') \\ &= k_0 \sin \theta \cos(\phi' - \phi) \end{aligned}$$

Far Field of Circular Patch (cont.)

Hence

$$H_{\phi'}^{pw} = \sec(k_{z1}h) \cos k_{z1}(z' + h) e^{j(k_0a) \sin \theta \cos(\phi' - \phi)} \cdot \left[-\sin \phi' H_x^{pw}(0,0,0) + \cos \phi' H_y^{pw}(0,0,0) \right]$$

Since the horizontal magnetic field components are modeled as current in the TEN, we have

$$H_{x,y}^{pw}(0,0,0) = H_{x,y}^{ipw}(0,0,0)(1 - \Gamma(\theta))$$

$$p = \theta : \text{TM}, \quad p = \phi : \text{TE}$$

Far Field of Circular Patch (cont.)

TM_z ($\hat{\underline{p}} = \hat{\underline{\theta}}$)

$$H_x^{ipw}(0,0,0) = \left(-\frac{E_0}{\eta_0} \hat{\underline{\phi}} \right) \cdot \hat{\underline{x}} = \frac{E_0}{\eta_0} (\sin \phi)$$

$$H_y^{ipw}(0,0,0) = \left(-\frac{E_0}{\eta_0} \hat{\underline{\phi}} \right) \cdot \hat{\underline{y}} = \frac{E_0}{\eta_0} (-\cos \phi)$$

TE_z ($\hat{\underline{p}} = \hat{\underline{\phi}}$)

$$H_x^{ipw}(0,0,0) = \left(\frac{E_0}{\eta_0} \hat{\underline{\theta}} \right) \cdot \hat{\underline{x}} = \frac{E_0}{\eta_0} (\cos \theta \cos \phi)$$

$$H_y^{ipw}(0,0,0) = \left(\frac{E_0}{\eta_0} \hat{\underline{\theta}} \right) \cdot \hat{\underline{y}} = \frac{E_0}{\eta_0} (\cos \theta \sin \phi)$$

Far Field E_θ

$$\text{TM}_z \quad (\hat{p} = \hat{\theta})$$

Substituting for H_x and H_y , we have

$$H_{\phi'}^{pw} = \sec(k_{z1}h) \cos k_{z1}(z' + h) e^{j(k_0a) \sin \theta \cos(\phi' - \phi)} \cdot \left(\frac{E_0}{\eta_0} \right) [-\sin \phi' \sin \phi - \cos \phi' \cos \phi] (1 - \Gamma^{TM}(\theta))$$

Note: $[-\sin \phi' \sin \phi - \cos \phi' \cos \phi] = -\cos(\phi' - \phi)$

Hence, we have

$$E_\theta^{FF}(r, \theta, \phi) = \int_0^{2\pi} \int_{-h}^0 \frac{E_0}{\eta_0} \sec(k_{z1}h) \cos k_{z1}(z' + h) e^{j(k_0a) \sin \theta \cos(\phi' - \phi)} \cdot (1 - \Gamma^{TM}(\theta)) (\cos(\phi' - \phi)) \cos \phi' a dz' d\phi'$$

Far Field E_θ (cont.)

For the z' integral we have

$$\sec(k_{z1}h) \int_{-h}^0 \cos k_{z1}(z' + h) dz' = h \operatorname{tanc}(k_{z1}h)$$

so that

$$E_\theta^{FF}(r, \theta, \phi) = a \left(\frac{E_0}{\eta_0} \right) h \operatorname{tanc}(k_{z1}h) (1 - \Gamma^{TM}(\theta)) I_{TM}$$

where

$$I_{TM} \equiv \int_0^{2\pi} e^{jq \cos(\phi' - \phi)} \cos(\phi' - \phi) \cos \phi' d\phi'$$

$$q \equiv (k_0 a) \sin \theta$$

Let $\phi'' = \phi' - \phi$

$$I_{TM} = \int_{-\phi}^{2\pi - \phi} e^{jq \cos(\phi'')} \cos(\phi'') \cos(\phi'' + \phi) d\phi''$$

Far Field E_θ (cont.)

We have that

$$\cos(\phi'' + \phi) = \cos \phi'' \cos \phi - \sin \phi'' \sin \phi$$

$$\text{and } \int_{-\phi}^{2\pi-\phi} () d\phi'' = \int_0^{2\pi} () d\phi''$$

so that

$$\begin{aligned} I_{TM} &\equiv \int_0^{2\pi} e^{jq \cos(\phi'')} \cos(\phi'') \cos(\phi'' + \phi) d\phi'' \\ &= \cos \phi \int_0^{2\pi} e^{jq \cos \phi''} \cos^2 \phi'' d\phi'' \\ &\quad - \sin \phi \int_0^{2\pi} e^{jq \cos \phi''} \sin \phi'' \cos \phi'' d\phi'' \end{aligned}$$

Integrates to zero
(odd function)

Now use $\cos^2 \phi'' = 1 - \sin^2 \phi''$

Far Field E_θ (cont.)

$$I_{TM} = \cos \phi \int_0^{2\pi} e^{jq \cos \phi''} d\phi'' - \cos \phi \int_0^{2\pi} e^{jq \cos \phi''} \sin^2 \phi'' d\phi''$$

Now we use the following identity:

$$\int_0^{2\pi} e^{jq \cos \phi''} \sin^{2n} \phi'' d\phi'' = J_n(q) \left[\frac{2^{n+1} \sqrt{\pi} \Gamma\left(n + \frac{1}{2}\right)}{q^n} \right]$$

$$n = 0, 1, 2, \dots$$

where

$$\Gamma\left(\frac{1}{2}\right) = \sqrt{\pi}$$

$$\Gamma\left(\frac{3}{2}\right) = \frac{1}{2} \sqrt{\pi}$$

Far Field E_θ (cont.)

Hence

$$\int_0^{2\pi} e^{jq \cos \phi''} d\phi'' = 2\pi J_0(q) \quad (n=0)$$

and

$$\int_0^{2\pi} e^{jq \cos \phi''} \sin^2 \phi'' d\phi'' = 2\pi \left(\frac{J_1(q)}{q} \right) \quad (n=1)$$

and thus

$$I_{TM} = \cos \phi (2\pi) \left[J_0(q) - \frac{J_1(q)}{q} \right]$$

Next, use

$$J'_n(x) = J_{n-1}(x) - \frac{n}{x} J_n(x)$$

so that

$$J'_1(x) = J_0(x) - \frac{J_1(x)}{x}$$

Far Field E_θ (cont.)

Hence

$$I_{TM} = 2\pi \cos \phi J_1'(q)$$

The far field is then

$$E_\theta^{FF}(r, \theta, \phi) = \frac{E_0}{\eta_0} (ah) \operatorname{tanc}(k_{z1}h) Q(\theta) \cdot 2\pi \cos \phi J_1'(k_0 a \sin \theta)$$

where

$$Q(\theta) = 1 - \Gamma^{TM}(\theta)$$

Far Field E_ϕ

$$\text{TE}_z \left(\underline{\hat{p}} = \underline{\hat{\phi}} \right)$$

Performing similar steps, we have

$$H_{\phi'}^{pw} = \sec(k_{z1}h) \cos k_{z1}(z' - h) (1 - \Gamma^{TE}(\theta)) e^{j(k_0 a) \sin \theta \cos(\phi' - \phi)} \\ \cdot \left(\frac{E_0}{\eta_0} \right) [-\sin \phi' \cos \theta \cos \phi + \cos \phi' \cos \theta \sin \phi]$$

Using reciprocity and performing the integration in z'' , we have

$$E_\phi^{FF}(r, \theta, \phi) = a \left(\frac{E_0}{\eta_0} \right) (h) \text{tanc}(k_{z1}h) (1 - \Gamma^{TE}(\theta)) \\ \cdot \cos \theta \int_0^{2\pi} e^{jq \cos(\phi' - \phi)} \sin(\phi' - \phi) \cos \phi' d\phi'$$

Far Field E_ϕ (cont.)

Evaluating the integral, we have

$$\begin{aligned} I_{TE} &\equiv \int_0^{2\pi} e^{jq \cos(\phi' - \phi)} \sin(\phi' - \phi) \cos \phi' d\phi' \\ &= \int_0^{2\pi} e^{jq \cos \phi''} \sin \phi'' \cos(\phi'' + \phi) d\phi'' \\ &= \int_0^{2\pi} e^{jq \cos \phi''} \sin \phi'' [\cos \phi'' \cos \phi - \sin \phi'' \sin \phi] d\phi'' \\ &= \cos \phi \int_0^{2\pi} e^{jq \cos \phi''} \sin \phi'' \cos \phi'' d\phi'' - \sin \phi \int_0^{2\pi} e^{jq \cos \phi''} \sin^2 \phi'' d\phi'' \end{aligned}$$

integrates to zero
(odd function)

$$= -\sin \phi 2\pi \left(\frac{J_1(q)}{q} \right)$$

Far Field E_ϕ (cont.)

Hence

$$E_\phi^{FF}(r, \theta, \phi) = -\frac{E_0}{\eta_0} (ah) \operatorname{tanc}(k_{z1}h) \sin \phi 2\pi \left(\frac{J_1(k_0 a \sin \theta)}{k_0 a \sin \theta} \right) P(\theta)$$

where

$$P(\theta) \equiv \cos \theta (1 - \Gamma^{TE}(\theta))$$