Introduction to Microstrip Antennas

David R. Jackson
Dept. of ECE
University of Houston
Contact Information

David R. Jackson
Dept. of ECE
N308 Engineering Building 1
University of Houston
Houston, TX 77204-4005

Phone: 713-743-4426
Fax: 713-743-4444
Email: djackson@uh.edu
Purpose of Short Course

- Provide an introduction to microstrip antennas.
- Provide a physical and mathematical basis for understanding how microstrip antennas work.
- Provide a physical understanding of the basic physical properties of microstrip antennas.
- Provide an overview of some of the recent advances and trends in the area (but not an exhaustive survey – directed towards understanding the fundamental principles).
Additional Resources

- Some basic references are provided at the end of these viewgraphs.
- You are welcome to visit a website that goes along with a course at the University of Houston on microstrip antennas (PowerPoint viewgraphs from the course may be found there, along with the viewgraphs from this short course).

ECE 6345: Microstrip Antennas

http://courses.egr.uh.edu/ECE/ECE6345/

Note:
You are welcome to use anything that you find on this website, as long as you please acknowledge the source.
Outline

- Overview of microstrip antennas
- Feeding methods
- Basic principles of operation
- General characteristics
- CAD Formulas
- Radiation pattern
- Input Impedance
- Circular polarization
- Circular patch
- Improving bandwidth
- Miniaturization
- Reducing surface waves and lateral radiation
Notation

\( c = \text{speed of light in free space} \)

\( \lambda_0 = \text{wavelength of free space} \)

\( k_0 = \text{wavenumber of free space} \)

\( k_1 = \text{wavenumber of substrate} \)

\( \eta_0 = \text{intrinsic impedance of free space} \)

\( \eta_1 = \text{intrinsic impedance of substrate} \)

\( \varepsilon_r = \text{relative permittivity (dielectric constant) of substrate} \)

\( \varepsilon_{r_{\text{eff}}} = \text{effective relative permittivity} \)  
(accurting for fringing of flux lines at edges)

\( \varepsilon_{r_{\text{eff}}}^{\text{complex}} = \text{complex effective relative permittivity} \)  
(used in the cavity model to account for all losses)

\[ c = 2.99792458 \times 10^8 \text{ [m/s]} \]

\[ \lambda_0 = \frac{c}{f} \]

\[ k_0 = \omega \sqrt{\mu_0 \varepsilon_0} = \frac{2\pi}{\lambda_0} \]

\[ k_1 = k_0 \sqrt{\varepsilon_r} \]

\[ \eta_0 = \sqrt{\frac{\mu_0}{\varepsilon_0}} \approx 376.7303 \text{ [\Omega]} \]

\[ \eta_1 = \frac{\eta_0}{\sqrt{\varepsilon_r}} \]

\[ c = \frac{1}{\sqrt{\varepsilon_0 \mu_0}} \]

\[ \mu_0 = 4\pi \times 10^7 \text{ [H/m]} \]

\[ \varepsilon_0 = \frac{1}{\mu_0 c^2} \approx 8.854188 \times 10^{12} \text{ [F/m]} \]
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Overview of Microstrip Antennas

Also called “patch antennas”

- One of the most useful antennas at microwave frequencies ($f > 1 \text{ GHz}$).
- It usually consists of a metal “patch” on top of a grounded dielectric substrate.
- The patch may be in a variety of shapes, but rectangular and circular are the most common.

Microstrip line feed

Coax feed
Overview of Microstrip Antennas

Common Shapes

- Rectangular
- Square
- Circular
- Annular ring
- Elliptical
- Triangular
Overview of Microstrip Antennas

History

- Invented by Bob Munson in 1972 (but earlier work by Dechamps goes back to 1953).
- Became popular starting in the 1970s.


Advantages of Microstrip Antennas

- Low profile (can even be “conformal,” i.e. flexible to conform to a surface).
- Easy to fabricate (use etching and photolithography).
- Easy to feed (coaxial cable, microstrip line, etc.).
- Easy to incorporate with other microstrip circuit elements and integrate into systems.
- Patterns are somewhat hemispherical, with a moderate directivity (about 6-8 dB is typical).
- Easy to use in an array to increase the directivity.
Overview of Microstrip Antennas

Disadvantages of Microstrip Antennas

- Low bandwidth (but can be improved by a variety of techniques). Bandwidths of a few percent are typical. Bandwidth is roughly proportional to the substrate thickness and inversely proportional to the substrate permittivity.

- Efficiency may be lower than with other antennas. Efficiency is limited by conductor and dielectric losses*, and by surface-wave loss**.

- Only used at microwave frequencies and above (the substrate becomes too large at lower frequencies).

- Cannot handle extremely large amounts of power (dielectric breakdown).

* Conductor and dielectric losses become more severe for thinner substrates.

** Surface-wave losses become more severe for thicker substrates (unless air or foam is used).
Overview of Microstrip Antennas

Applications

Applications include:

- Satellite communications
- Microwave communications
- Cell phone antennas
- GPS antennas
Overview of Microstrip Antennas

Microstrip Antenna Integrated into a System: HIC Antenna Base-Station for 28-43 GHz

(Photo courtesy of Dr. Rodney B. Waterhouse)
Overview of Microstrip Antennas

Arrays

Linear array (1-D corporate feed)

2×2 array

2-D 8X8 corporate-fed array

4 × 8 corporate-fed / series-fed array
Overview of Microstrip Antennas

Wraparound Array (conformal)

The substrate is so thin that it can be bent to “conform” to the surface.

(Photo courtesy of Dr. Rodney B. Waterhouse)
Overview of Microstrip Antennas

Rectangular patch

Note:
The fields and current are approximately independent of $y$ for the dominant (1,0) mode.

Note:
$L$ is the resonant dimension (direction of current flow). The width $W$ is usually chosen to be larger than $L$ (to get higher bandwidth). However, usually $W < 2L$ (to avoid problems with the (0,2) mode).

$W = 1.5L$ is typical.
Overview of Microstrip Antennas

Circular Patch

The location of the feed determines the direction of current flow and hence the polarization of the radiated field.
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Feeding Methods

Some of the more common methods for feeding microstrip antennas are shown.

The feeding methods are illustrated for a rectangular patch, but the principles apply for circular and other shapes as well.
Coaxial Feed

A feed along the centerline at $y = W/2$ is the most common (this minimizes higher-order modes and cross-pol).

Note:

Surface current

Feed at $(x_0, y_0)$

$\varepsilon_r$

$W$

$L$

$h$

$x$

$y$

$z$
Advantages:

- Simple
- Directly compatible with coaxial cables
- Easy to obtain input match by adjusting feed position

Disadvantages:

- Significant probe (feed) radiation for thicker substrates
- Significant probe inductance for thicker substrates (limits bandwidth)
- Not easily compatible with arrays

\[ R = R_{edge} \cos^2 \left( \frac{\pi x_0}{L} \right) \]

(The resistance varies as the square of the modal field shape.)
Advantages:
- Simple
- Allows for planar feeding
- Easy to use with arrays
- Easy to obtain input match

Disadvantages:
- Significant line radiation for thicker substrates
- For deep notches, patch current and radiation pattern may show distortion
Feeding Methods

Inset Feed

An investigation has shown that the resonant input resistance varies as:

\[ R_{in} = A \cos^2 \left( \frac{\pi}{2} \left( \frac{2x_0}{L} - B \right) \right) \]

Less accurate approximation:

\[ R \approx R_{edge} \cos^2 \left( \frac{\pi x_0}{L} \right) \]

The coefficients \( A \) and \( B \) depend on the notch width \( S \) but (to a good approximation) not on the line width \( W_f \).

Feeding Methods

Proximity-coupled Feed
(Electromagnetically-coupled Feed)

Advantages:
- Allows for planar feeding
- Less line radiation compared to microstrip feed (the line is closer to the ground plane)
- Can allow for higher bandwidth (no probe inductance, so substrate can be thicker)

Disadvantages:
- Requires multilayer fabrication
- Alignment is important for input match
Feeding Methods

Gap-coupled Feed

Advantages:
- Allows for planar feeding
- Can allow for a match even with high edge impedances, where a notch might be too large (e.g., when using a high permittivity substrate)

Disadvantages:
- Requires accurate gap fabrication
- Requires full-wave design
Advantages:

- Allows for planar feeding
- Feed-line radiation is isolated from patch radiation
- Higher bandwidth is possible since probe inductance is eliminated (allowing for a thick substrate), and also a double-resonance can be created
- Allows for use of different substrates to optimize antenna and feed-circuit performance

Disadvantages:

- Requires multilayer fabrication
- Alignment is important for input match
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Basic Principles of Operation

- The basic principles are illustrated here for a rectangular patch, but the principles apply similarly for other patch shapes.

- We use the cavity model to explain the operation of the patch antenna.

Basic Principles of Operation

Main Ideas:

- The patch acts approximately as a resonant cavity (with short-circuit (PEC) walls on top and bottom, open-circuit (PMC) walls on the edges).
- Radiation is accounted for by using an effective loss tangent for the substrate.
- In a cavity, only certain modes are allowed to exist, at different resonance frequencies.
- If the antenna is excited at a resonance frequency, a strong field is set up inside the cavity, and a strong current on the (bottom) surface of the patch. This produces significant radiation (a good antenna).
Basic Principles of Operation

A microstrip antenna can radiate well, even with a thin substrate, because of resonance.

- As the substrate gets thinner the patch current radiates less, due to image cancellation (current and image are separated by $2h$).
- However, the $Q$ of the resonant cavity mode also increases, making the patch currents stronger at resonance.
- These two effects cancel, allowing the patch to radiate well even for thin substrates (though the bandwidth decreases).

\[
|J_s| \propto Q \propto \frac{1}{h}
\]

\[\mathcal{E}_r\]

\[J_s\]
Basic Principles of Operation

Thin Substrate Approximation

On patch and ground plane: \[ E_t = 0 \quad \Rightarrow \quad \vec{E} = \hat{z} E_z \]

Inside the patch cavity, because of the thin substrate, the electric field vector is approximately independent of \( z \).

Hence \[ E(x, y, z) \approx \hat{z} E_z(x, y) \]
Basic Principles of Operation

Thin Substrate Approximation

Magnetic field inside patch cavity:

\[
H = -\frac{1}{j\omega \mu} \nabla \times E
\]

\[
= -\frac{1}{j\omega \mu} \nabla \times (\hat{z} E_z (x, y))
\]

\[
= -\frac{1}{j\omega \mu} (-\hat{z} \times \nabla E_z (x, y))
\]
Thin Substrate Approximation

$$H(x, y) = \frac{1}{j \omega \mu} \left( \hat{z} \times \nabla E_z(x, y) \right)$$

**Note:**
The magnetic field is purely horizontal. (The mode is $TM_z$.)
On the edges of the patch:

\[ \mathbf{J}_s \cdot \hat{n} = 0 \]

(\(\mathbf{J}_s\) is the sum of the top and bottom surface currents.)

On the bottom surface of the patch conductor, at the edge of the patch, we have:

\[ \mathbf{J}^{bot}_s \cdot \hat{n} \approx 0 \quad \text{(assuming } |\mathbf{J}^{bot}_s| \gg |\mathbf{J}^{top}_s|) \]

Also,

\[ \mathbf{J}^{bot}_s = (-\hat{z} \times \mathbf{H}) \]

\[ \mathbf{H}^{bot}_l \approx 0 \]
Since the magnetic field is approximately independent of $z$, we have an approximate PMC condition on the entire vertical edge.

\[ H_t = 0 \ (\text{PMC}) \]

or

\[ \mathbf{n} \times H(x, y) = 0 \]

Actual patch

PMC Model
Basic Principles of Operation

Magnetic-wall Approximation

\[
\hat{n} \times H(x, y) = 0
\]

\[
H(x, y) = \frac{1}{j \omega \mu} (\hat{z} \times \nabla E_z(x, y))
\]

Hence,

\[
\hat{n} \times (\hat{z} \times \nabla E_z(x, y)) = 0
\]

\[
\hat{z} \left( \hat{n} \cdot \nabla E_z(x, y) \right) = 0
\]

\[
\frac{\partial E_z}{\partial n} = 0 \quad \text{(Neumann B.C.)}
\]
Basic Principles of Operation

Resonance Frequencies

\[ \nabla^2 E_z + k^2 E_z = 0 \]

\[ k = k_1 = k_0 \sqrt{\varepsilon_r} \]

From separation of variables:

\[ E_z = \cos\left(\frac{m\pi x}{L}\right) \cos\left(\frac{n\pi y}{W}\right) \]

(TM\textsubscript{mn} mode)

We then have

\[ \left[-\left(\frac{m\pi}{L}\right)^2 - \left(\frac{n\pi}{W}\right)^2 + k_1^2\right] E_z = 0 \]

Hence

\[ \left[-\left(\frac{m\pi}{L}\right)^2 - \left(\frac{n\pi}{W}\right)^2 + k_1^2\right] = 0 \]

Note:
We ignore the loss tangent of the substrate for the calculation of the resonance frequencies.
We thus have

\[ k_1^2 = \left( \frac{m\pi}{L} \right)^2 + \left( \frac{n\pi}{W} \right)^2 \]

Recall that

\[ k_1 = k_0 \sqrt{\varepsilon_r} = \omega \sqrt{\mu_0 \varepsilon_0 \varepsilon_r} \]

\[ \omega = 2\pi f \]

Hence

\[ f = \frac{c}{2\pi \sqrt{\varepsilon_r}} \sqrt{\left( \frac{m\pi}{L} \right)^2 + \left( \frac{n\pi}{W} \right)^2} \]

\[ c = \frac{1}{\sqrt{\mu_0 \varepsilon_0}} \]
Hence \( f = f_{mn} \) (resonance frequency of \((m,n)\) mode)

where

\[
f_{mn} = \frac{c}{2\pi\sqrt{\varepsilon_r}} \sqrt{\left(\frac{m\pi}{L}\right)^2 + \left(\frac{n\pi}{W}\right)^2}
\]
Basic Principles of Operation

Dominant (1,0) mode

This structure operates as a “fat planar dipole.”

This mode is usually used because the radiation pattern has a broadside beam.

\[ f_{10} = \frac{c}{2\sqrt{\varepsilon_r}} \left( \frac{1}{L} \right) \]

\[ E_z = \cos \left( \frac{\pi x}{L} \right) \]

\[ H(x, y) = -\hat{y} \left( \frac{1}{j\omega\mu} \right) \left( \frac{\pi}{L} \right) \sin \left( \frac{\pi x}{L} \right) \]

\[ J_s = \hat{x} \left( \frac{-1}{j\omega\mu_0} \right) \left( \frac{\pi}{L} \right) \sin \left( \frac{\pi x}{L} \right) \]

The resonant length \( L \) is about 0.5 guided wavelengths in the \( x \) direction (see next slide).
The resonance frequency is mainly controlled by the patch length \( L \) and the substrate permittivity.

Approximately, (assuming PMC walls)

\[
\begin{align*}
k_1^2 &= \left( \frac{m\pi}{L} \right)^2 + \left( \frac{n\pi}{W} \right)^2 \\
(1,0) \text{ mode: } k_1L &= \pi \\
k_1 &= \frac{2\pi}{\lambda_d}
\end{align*}
\]

This is equivalent to saying that the length \( L \) is one-half of a wavelength in the dielectric.

\[
L = \frac{\lambda_d}{2} = \frac{\lambda_0}{2} \sqrt{\varepsilon_r}
\]

Comment:
A higher substrate permittivity allows for a smaller antenna (miniaturization), but with a lower bandwidth.
The resonance frequency calculation can be improved by adding a “fringing length extension” $\Delta L$ to each edge of the patch to get an “effective length” $L_e$.

$$L_e = L + 2\Delta L$$

Note: Some authors use *effective permittivity* in this equation. (This would change the value of $L_e$.)

$$f_{10} = \frac{c}{2\sqrt{\varepsilon_r}} \left( \frac{1}{L_e} \right)$$
Basic Principles of Operation

Resonance Frequency of Dominant Mode

Hammerstad formula:

$$\Delta L / h = 0.412 \left[ \frac{(\varepsilon_r^{\text{eff}} + 0.3) \left( \frac{W}{h} + 0.264 \right)}{(\varepsilon_r^{\text{eff}} - 0.258) \left( \frac{W}{h} + 0.8 \right)} \right]$$

$$\varepsilon_r^{\text{eff}} = \frac{\varepsilon_r + 1}{2} + \left( \frac{\varepsilon_r - 1}{2} \right) \left[ 1 + 12 \left( \frac{h}{W} \right) \right]^{-1/2}$$

Note:
Even though the Hammerstad formula involves an effective permittivity, we still use the actual substrate permittivity in the resonance frequency formula.

$$f_{10} = \frac{c}{2\sqrt{\varepsilon_r}} \left( \frac{1}{L + 2\Delta L} \right)$$
Basic Principles of Operation

Resonance Frequency of Dominant Mode

Note: $\Delta L \approx 0.5 \ h$

This is a good “rule of thumb” to give a quick estimate.
The resonance frequency has been normalized by the zero-order value (without fringing).
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The bandwidth is directly proportional to substrate thickness \( h \).

However, if \( h \) is greater than about 0.05 \( \lambda_0 \), the probe inductance (for a coaxial feed) becomes large enough so that matching is difficult – the bandwidth will decrease.

The bandwidth is inversely proportional to \( \varepsilon_r \) (a foam substrate gives a high bandwidth).

The bandwidth of a rectangular patch is proportional to the patch width \( W \) (but we need to keep \( W < 2L \); see the next slide).
General Characteristics

Width Restriction for a Rectangular Patch

\[ f_{mn} = \frac{c}{2\pi \sqrt{\varepsilon_r}} \sqrt{\left( \frac{m\pi}{L} \right)^2 + \left( \frac{n\pi}{W} \right)^2} \]

\[ W < 2L \]

\[ f_{01} = \frac{c}{2\sqrt{\varepsilon_r}} \left( \frac{1}{W} \right) \]

\[ f_{10} = \frac{c}{2\sqrt{\varepsilon_r}} \left( \frac{1}{L} \right) \]

\[ f_{02} = \frac{c}{2\sqrt{\varepsilon_r}} \left( \frac{2}{W} \right) \]

\[ f_{02} - f_{01} = \frac{c}{\sqrt{\varepsilon_r}} \left( \frac{1}{W} - \frac{1}{2L} \right) \]

\[ W = 1.5 \ L \text{ is typical.} \]
General Characteristics

Some Bandwidth Observations

- For a typical substrate thickness \((h / \lambda_0 = 0.02)\), and a typical substrate permittivity \((\varepsilon_r = 2.2)\) the bandwidth is about 3%.

- By using a thick foam substrate, bandwidth of about 10% can be achieved.

- By using special feeding techniques (aperture coupling) and stacked patches, bandwidths of 100% have been achieved.
General Characteristics

Results: Bandwidth

The discrete data points are measured values. The solid curves are from a CAD formula (given later).

\[ \varepsilon_r = 2.2 \text{ or } 10.8 \quad \frac{W}{L} = 1.5 \]
The resonant input resistance is fairly independent of the substrate thickness $h$ unless $h$ gets small (the variation is then mainly due to dielectric and conductor loss).

The resonant input resistance is proportional to $\varepsilon_r$.

The resonant input resistance is directly controlled by the location of the feed point (maximum at edges $x = 0$ or $x = L$, zero at center of patch).
Note:
The patch is usually fed along the centerline \( y_0 = W / 2 \) to maintain symmetry and thus minimize excitation of undesirable modes (which cause cross-pol).

Desired mode: (1,0)
For a given mode, it can be shown that the resonant input resistance is proportional to the square of the cavity-mode field at the feed point.

This is seen from the cavity-model eigenfunction analysis (please see the reference).

\[ R_{in} \propto E_z^2 (x_0, y_0) \]

For \((1,0)\) mode:

\[ R_{in} \propto \cos^2 \left( \frac{\pi x_0}{L} \right) \]

Hence, for $(1,0)$ mode:

\[ R_{in} = R_{edge} \cos^2 \left( \frac{\pi x_0}{L} \right) \]

The value of $R_{edge}$ depends strongly on the substrate permittivity (it is proportional to the permittivity).

For a typical patch, it is often in the range of 100-200 Ohms.
Results: Resonant Input Resistance

The solid curves are from a CAD formula (given later.)

\[ \varepsilon_r = 10.8 \]

Region where loss is important

\[ \varepsilon_r = 2.2 \text{ or } 10.8 \quad W/L = 1.5 \]
Radiation Efficiency

- Radiation efficiency is the ratio of power radiated into space, to the total input power.

\[ e_r = \frac{P_r}{P_{tot}} \]

- The radiation efficiency is less than 100% due to
  - Conductor loss
  - Dielectric loss
  - Surface-wave excitation
General Characteristics

Radiation Efficiency (cont.)

- Surface wave 
- TM_0 
- Cos (\phi) pattern 
- J_s 

y

x
General Characteristics

Radiation Efficiency (cont.)

Hence,

\[
e_r = \frac{P_r}{P_{\text{tot}}} = \frac{P_r}{P_r + (P_c + P_d + P_{\text{sw}})}
\]

\[P_r = \text{radiated power}\]

\[P_{\text{tot}} = \text{total input power}\]

\[P_c = \text{power dissipated by conductors}\]

\[P_d = \text{power dissipated by dielectric}\]

\[P_{\text{sw}} = \text{power launched into surface wave}\]
Conductor and dielectric loss is more important for thinner substrates (the $Q$ of the cavity is higher, and thus the resonance is more seriously affected by loss).

Conductor loss increases with frequency (proportional to $f^{1/2}$) due to the skin effect. It can be very serious at millimeter-wave frequencies.

Conductor loss is usually more important than dielectric loss for typical substrate thicknesses and loss tangents.

$$R_s = \frac{1}{\sigma \delta} \quad \delta = \sqrt{\frac{2}{\omega \mu \sigma}}$$

$R_s$ is the surface resistance of the metal. The skin depth of the metal is $\delta$.

$$R_s = \sqrt{\frac{\omega \mu_0}{2\sigma}} \propto \sqrt{f}$$
Surface-wave power is more important for thicker substrates or for higher-substrate permittivities. (The surface-wave power can be minimized by using a thin substrate or a foam substrate.)

- For a foam substrate, a high radiation efficiency is obtained by making the substrate thicker (minimizing the conductor and dielectric losses). There is no surface-wave power to worry about.

- For a typical substrate such as $\varepsilon_r = 2.2$, the radiation efficiency is maximum for $h / \lambda_0 \approx 0.02$. 
General Characteristics

Results: Efficiency (Conductor and dielectric losses are neglected.)

\[ \varepsilon_r = 2.2 \quad \text{or} \quad 10.8 \]

Note: CAD plot uses the Pozar formula (given later).

\[ \frac{W}{L} = 1.5 \quad \text{Note: CAD plot uses the Pozar formula (given later).} \]
Results: Efficiency (All losses are accounted for.)

**General Characteristics**

Note: CAD plot uses the Pozar formula (given later).
General Characteristics

Radiation Pattern

E-plane: co-pol is $E_{\theta}$

H-plane: co-pol is $E_{\phi}$

Note:
For radiation patterns, it is usually more convenient to place the origin at the middle of the patch (this keeps the formulas as simple as possible).
Comments on radiation patterns:

- The E-plane pattern is typically broader than the H-plane pattern.

- The truncation of the ground plane will cause edge diffraction, which tends to degrade the pattern by introducing:
  - Rippling in the forward direction
  - Back-radiation

- Pattern distortion is more severe in the E-plane, due to the angle dependence of the vertical polarization $E_\theta$ on the ground plane. (It varies as $\cos (\phi)$).
General Characteristics

Radiation Patterns

Edge diffraction is the most serious in the E plane.

\[ E_\theta \text{ varies as } \cos \phi \]
General Characteristics

Radiation Patterns

**E-plane pattern**

Red: infinite substrate and ground plane

Blue: 1 meter ground plane

**Note:**
The E-plane pattern “tucks in” and tends to zero at the horizon due to the presence of the infinite substrate.
General Characteristics

Radiation Patterns

**H-plane pattern**

*Red:* infinite substrate and ground plane

*Blue:* 1 meter ground plane

![Diagram showing radiation patterns with labels for H-plane pattern, red for infinite substrate and ground plane, and blue for 1 meter ground plane.](image)
General Characteristics

Directivity

- The directivity is fairly insensitive to the substrate thickness.

- The directivity is higher for lower permittivity, because the patch is larger.
General Characteristics

Results: Directivity (relative to isotropic)

Directivity (dB)

$\varepsilon_r = 2.2$

$\varepsilon_r = 2.2$ or $10.8$, $W/L = 1.5$
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CAD Formulas

CAD formulas for the important properties of the *rectangular microstrip antenna* will be shown.

- Radiation efficiency
- Bandwidth ($Q$)
- Resonant input resistance
- Directivity


CAD Formulas

Radiation Efficiency

\[ e_r = \frac{e_r^{hed}}{1 + e_r^{hed} \left[ \ell_d + \left( \frac{R_s^{ave}}{\pi \eta_0} \right) \left( \frac{1}{h / \lambda_0} \right) \right] \left[ \left( \frac{3}{16} \right) \left( \frac{\varepsilon_r}{\rho c_1} \right) \left( \frac{L}{W} \right) \left( \frac{1}{h / \lambda_0} \right) \right]} \]

Comment:
The efficiency becomes small as the substrate gets thin, if there is dielectric or conductor loss.

where

\[ \ell_d = \tan \delta = \text{loss tangent of substrate} \]

\[ R_s = \text{surface resistance of metal} = \frac{1}{\sigma \delta} = \sqrt{\frac{\omega \mu}{2\sigma}} \]

\[ R_s^{ave} = \left( R_s^{patch} + R_s^{ground} \right) / 2 \]

Note: “hed” refers to a unit-amplitude horizontal electric dipole.
CAD Formulas

Radiation Efficiency (cont.)

\[ e_r^{hed} = \frac{P_{sp}^{hed}}{P_{sp}^{hed} + P_{sw}^{hed}} = \frac{1}{1 + \frac{P_{sw}^{hed}}{P_{sp}^{hed}}} \]

where

\[ P_{sp}^{hed} = \frac{1}{\lambda_0^2} (k_0 h)^2 \left( 80\pi^2 c_1 \right) \]

\[ P_{sw}^{hed} = \frac{1}{\lambda_0^2} (k_0 h)^3 \left[ 60\pi^3 \left( 1 - \frac{1}{\varepsilon_r} \right)^3 \right] \]

**Note:** “hed” refers to a unit-amplitude horizontal electric dipole.

**Note:** When we say “unit amplitude” here, we assume peak (not RMS) values.
Hence, we have

\[ e_r^{hed} = \frac{1}{1 + \frac{3}{4} \pi (k_0 h) \left( \frac{1}{c_1} \right)^3 \left( 1 - \frac{1}{\varepsilon_r} \right)^3} \]

Physically, this term is the radiation efficiency of a horizontal electric dipole (hed) on top of the substrate.
The constants are defined as follows:

\[
c_1 = 1 - \frac{1}{\varepsilon_r} + \frac{2}{5\varepsilon_r^2}
\]

\[
p = 1 + \frac{a_2}{10}(k_0 W)^2 + \left(a_2^2 + 2a_4\right)\left(\frac{3}{560}\right)(k_0 W)^4 + c_2\left(\frac{1}{5}\right)(k_0 L)^2
\]

\[
+ a_2 c_2 \left(\frac{1}{70}\right)(k_0 W)^2 (k_0 L)^2
\]

\[
c_2 = -0.0914153
\]

\[
a_2 = -0.16605
\]

\[
a_4 = 0.00761
\]
Improved formula for HED surface-wave power (due to Pozar)

\[
P_{\text{sw}}^{\text{hed}} = \frac{\eta_0 k_0^2}{8} \frac{\varepsilon_r \left(x_0^2 - 1\right)^{3/2}}{\varepsilon_r \left(1 + x_1\right) + (k_0 h) \sqrt{x_0^2 - 1} \left(1 + \varepsilon_r^2 x_1\right)}
\]

**Note:** \(x_0\) in this formula is not the feed location!

\[
x_1 = \frac{x_0^2 - 1}{\varepsilon_r - x_0^2}
\]

\[
x_0 = 1 + \frac{-\varepsilon_r^2 + \alpha_0 \alpha_1 + \varepsilon_r \sqrt{\varepsilon_r^2 - 2 \alpha_0 \alpha_1 + \alpha_0^2}}{\varepsilon_r^2 - \alpha_1^2}
\]

\[
\alpha_0 = s \tan \left[(k_0 h)s\right]
\]

\[
\alpha_1 = -\frac{1}{s} \left[\tan \left[(k_0 h)s\right] + \frac{(k_0 h)s}{\cos^2 \left[(k_0 h)s\right]}\right]
\]

\[
s = \sqrt{\varepsilon_r - 1}
\]


**Note:** The above formula for the surface-wave power is different from that given in Pozar’s paper by a factor of 2, since Pozar used RMS instead of peak values.
CAD Formulas

Bandwidth

\[
BW = \frac{1}{\sqrt{2}} \left[ \ell_d + \left( \frac{R_s^{ave}}{\pi \eta_0} \right) \left( \frac{1}{h / \lambda_0} \right) + \left( \frac{16}{3} \right) \left( \frac{p c_1}{\varepsilon_r} \right) \left( \frac{h}{\lambda_0} \right) \left( \frac{W}{L} \right) \left( \frac{1}{e_r^{hed}} \right) \right]
\]

\[
Q = \frac{1}{\sqrt{2} \ BW}
\]

Comments:
For a lossless patch, the bandwidth is approximately proportional to the patch width and to the substrate thickness. It is inversely proportional to the substrate permittivity. For very thin substrates the bandwidth will increase for a lossy patch, but at the expense of efficiency.

\( BW \) is defined from the frequency limits \( f_1 \) and \( f_2 \) at which \( SWR = 2.0 \).

\[
BW = \frac{f_2 - f_1}{f_0}
\]
(multiply by 100 if you want to get %)
CAD Formulas

Quality Factor $Q$

$$Q \equiv \omega_0 \frac{U_s}{P}$$

$U_s = \text{energy stored in patch cavity}$

$P = \text{power that is radiated and dissipated by patch}$

$$\frac{1}{Q} = \frac{P}{\omega_0 U_s}$$

$$P = P_d + P_c + P_{sp} + P_{sw}$$

$$\frac{1}{Q} = \frac{1}{Q_d} + \frac{1}{Q_c} + \frac{1}{Q_{sp}} + \frac{1}{Q_{sw}}$$
**Q Components**

\[ Q_d = \frac{1}{\tan \delta} \]

\[ Q_c = \left( \frac{\eta_0}{2} \right) \left[ \frac{(k_0 h)}{R_s^{\text{ave}}} \right] \]

\[ R_s^{\text{ave}} = \left( R_s^{\text{patch}} + R_s^{\text{ground}} \right) / 2 \]

\[ Q_{sp} \approx \frac{3}{16} \left( \frac{\varepsilon_r}{pc_1} \right) \left( \frac{L}{W} \right) \left( \frac{1}{h / \lambda_0} \right) \]

\[ Q_{sw} = Q_{sp} \left( \frac{e_r^{\text{hed}}}{1 - e_r^{\text{hed}}} \right) \]

The constants \( p \) and \( c_1 \) were defined previously.

\[ e_r^{\text{hed}} = \frac{1}{1 + \frac{3}{4} \pi (k_0 h) \left( \frac{1}{c_1} \right) \left( 1 - \frac{1}{\varepsilon_r} \right)^3} \]
Resonant Input Resistance

**Probe-feed Patch**

\[ R = R_{\text{in}}^{\text{max}} = R_{\text{edge}} \cos^2 \left( \frac{\pi x_0}{L} \right) \]

\[
R_{\text{edge}} = \frac{\left( \frac{4\eta_0}{\pi} \right) \left( \frac{L}{W} \right) \left( \frac{h}{\lambda_0} \right)}{\ell_d + \left( \frac{R_s}{\pi \eta_0} \right) \left( \frac{1}{h/\lambda_0} \right) + \left( \frac{16}{3} \right) \left( \frac{pc_1}{\varepsilon_r} \right) \left( \frac{W}{L} \right) \left( \frac{h}{\lambda_0} \right) \left( \frac{1}{\varepsilon_r^{hed}} \right)}
\]

**Comments:**
For a lossless patch, the resonant resistance is approximately independent of the substrate thickness. For a lossy patch it tends to zero as the substrate gets very thin. For a lossless patch it is inversely proportional to the square of the patch width and it is proportional to the substrate permittivity.
Approximate CAD formula for probe (feed) reactance (in Ohms)

\[ X_p = \frac{\eta_0}{2\pi} (k_0 h) \left[ -\gamma + \ln \left( \frac{2}{\sqrt{\varepsilon_r (k_0 a)}} \right) \right] \]

This is based on an infinite parallel-plate model.

\[ X_p = \omega L_p \]

\[ \gamma \doteq 0.577216 \] (Euler’s constant)

\[ \eta_0 = \sqrt{\mu_0 / \varepsilon_0} = 376.7303 \ \Omega \]
CAD Formulas

Observations:

- Feed (probe) reactance increases proportionally with substrate thickness $h$.

- Feed reactance increases for smaller probe radius.

$$X_p = \frac{\eta_0}{2\pi} (k_0 h) \left[ -\gamma + \ln \left( \frac{2}{\sqrt{\varepsilon_r (k_0 a)}} \right) \right]$$

Important point:

If the substrate gets too thick, the probe reactance will make it difficult to get an input match, and the bandwidth will suffer.

(Compensating techniques will be discussed later.)
Results: Probe Reactance ($X_f = X_p = \omega L_p$)

The normalized feed location ratio $x_r$ is zero at the center of the patch ($x = L/2$), and is 1.0 at the patch edge ($x = L$).

$x_r = 2 \left( \frac{x_0}{L} \right) - 1$

CAD Formulas

- $\varepsilon_r = 2.2$
- $W/L = 1.5$
- $h = 0.0254 \lambda_0$
- $a = 0.5 \text{ mm}$
CAD Formulas

Directivity

\[ D = \left( \frac{3}{pc_1} \right) \left[ \frac{\varepsilon_r}{\varepsilon_r + \tan^2(k_1h)} \right] \left( \tan c^2(k_1h) \right) \]

\[ k_1 = k_0 \sqrt{\varepsilon_r} \]

where

\[ \tan c(x) \equiv \frac{\tan(x)}{x} \]

The constants \( p \) and \( c_1 \) were defined previously.
For thin substrates:

\[ D \approx \frac{3}{p c_1} \]

(The directivity is essentially independent of the substrate thickness.)
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There are two models often used for calculating the radiation pattern:

- Electric current model
- Magnetic current model

**Note:**

The origin is placed at the center of the patch, at the top of the substrate, for the pattern calculations.
Electric current model:
We keep the physical currents flowing on the patch (and feed).

\[ \mathbf{J}_{s}^{\text{patch}} = \mathbf{J}_{s}^{\text{top}} + \mathbf{J}_{s}^{\text{bot}} \]
Magnetic current model:
We apply the *equivalence principle* and invoke the (approximate) PMC condition at the edges.

\[
\begin{align*}
J_s^e &= \hat{n} \times H \\
M_s^e &= -\hat{n} \times E
\end{align*}
\]

The equivalent surface current is *approximately* zero on the top surface (weak fields) and the sides (PMC). We can ignore it on the ground plane (it does not radiate).
Theorem

The electric and magnetic models yield identical patterns at the resonance frequency of the cavity mode.

Assumption:

The electric and magnetic current models are based on the fields of a single cavity mode, corresponding to an ideal cavity with PMC walls.

Comments on the Substrate Effects

- The substrate can be neglected to simplify the far-field calculation.
- When considering the substrate, it is most convenient to assume an infinite substrate (in order to obtain a closed-form solution).
- Reciprocity can be used to calculate the far-field pattern of electric or magnetic current sources inside of an infinite layered structure.
- When an infinite substrate is assumed, the far-field pattern always goes to zero at the horizon.

Comments on the Two Models

- For the rectangular patch, the electric current model is the simplest since there is only one electric surface current (as opposed to four edges).
- For the rectangular patch, the magnetic current model allows us to classify the “radiating” and “nonradiating” edges.

\[ J_s = \hat{x} A_0 \cos \left( \frac{\pi x}{L} \right) \]

Note:
On the nonradiating edges, the magnetic currents are in opposite directions across the centerline \((x = 0)\).

\[ M_s^e = -\hat{n} \times E \]

\[ E_z = -\sin \left( \frac{\pi x}{L} \right) \]
Radiation Pattern

Rectangular Patch Pattern Formula

(The formula is based on the electric current model.)

\[ J_s = \hat{x} \cos \left( \frac{\pi x}{L} \right) \]

The origin is at the center of the patch.

(1,0) mode

The probe is on the x axis.
The far-field pattern can be determined by reciprocity.

\[ E_i (r, \theta, \phi) = E_i^{\text{hex}} (r, \theta, \phi) \left( \frac{\pi WL}{2} \right) \left[ \sin \left( \frac{k_y W}{2} \right) \right] \left[ \sin \left( \frac{k_y W}{2} \right) \right] \left[ \cos \left( \frac{k_x L}{2} \right) \right] \left( \frac{\pi}{2} - \left( \frac{k_x L}{2} \right)^2 \right) \]

\[ i = \theta \text{ or } \phi \]

\[ k_x = k_0 \sin \theta \cos \phi \]

\[ k_y = k_0 \sin \theta \sin \phi \]

The “hex” pattern is for a horizontal electric dipole in the \( x \) direction, sitting on top of the substrate.

\[ E^{\text{hex}}_{\phi}(r, \theta, \phi) = -E_0 \sin \phi \ F(\theta) \]
\[ E^{\text{hex}}_{\theta}(r, \theta, \phi) = E_0 \cos \phi \ G(\theta) \]

where

\[ E_0 = \left( \frac{-j\omega \mu_0}{4\pi r} \right) e^{-jk_0r} \]

\[ F(\theta) = 1 + \Gamma^{TE}(\theta) = \frac{2 \tan \left(k_0h N(\theta)\right)}{\tan \left(k_0h N(\theta)\right) - jN(\theta) \sec \theta} \]

\[ G(\theta) = \cos \theta \left(1 + \Gamma^{TM}(\theta)\right) = \frac{2 \tan \left(k_0h N(\theta)\right) \cos \theta}{\tan \left(k_0h N(\theta)\right) - j \frac{\varepsilon_r}{N(\theta)} \cos \theta} \]

\[ N(\theta) = \sqrt{\varepsilon_r - \sin^2(\theta)} \]

**Note:** To account for lossy substrate, use
\[ \varepsilon_r \rightarrow \varepsilon_{rc} = \varepsilon_r \left(1 - j \tan \delta\right) \]
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Various models have been proposed over the years for calculating the input impedance of a microstrip patch antenna.

- **Transmission line model**
  - The first model introduced
  - Very simple

- **Cavity model (eigenfunction expansion)**
  - Simple yet accurate for thin substrates
  - Gives physical insight into operation

- **CAD circuit model**
  - Extremely simple and almost as accurate as the cavity model

- **Spectral-domain method**
  - More challenging to implement
  - Accounts rigorously for both radiation and surface-wave excitation

- **Commercial software**
  - Very accurate
  - Can be time consuming
Results for a typical patch show that the first three methods agree very well, provided the correct $Q$ is used and the probe inductance is accounted for.
Input Impedance

Comparison of CAD with Full-Wave

\[ \varepsilon_r = 2.2 \]
\[ \tan \delta = 0.001 \]
\[ h = 1.524 \text{ mm} \]
\[ L = 6.255 \text{ cm} \]
\[ W / L = 1.5 \]
\[ \sigma = 3.0 \times 10^7 \text{ S/m} \]
\[ x_0 = 6.255 \text{ cm} \]
\[ y_0 = 0 \]
\[ a = 0.635 \text{ mm} \]

Results from full-wave analysis agree well with the simple CAD circuit model, except for a shift in resonance frequency.
The circuit model discussed assumes a probe feed. Other circuit models exist for other types of feeds.

Note:
The mathematical justification of the CAD circuit model comes from a cavity-model eigenfunction analysis.

Near the resonance frequency, the patch cavity can be approximately modeled as a resonant $RLC$ circuit.

- The resistance $R$ accounts for radiation and losses.
- A probe inductance $L_p$ is added in series, to account for the “probe inductance” of a probe feed.

Input Impedance

Probe-fed Patch

![Diagram of a probe-fed patch antenna with a resonant RLC circuit](image)
Input Impedance

\[
Z_{in} \approx j\omega L_p + \frac{R}{1 + jQ \left( \frac{f - f_0}{f_0} \right)}
\]

\[
Q = \frac{R}{\omega_0 L} \quad BW = \frac{1}{\sqrt{2}Q}
\]

**BW** is defined here by \(SWR < 2.0\) when the \(RLC\) circuit is fed by a matched line \((Z_0 = R)\).

\[
\omega_0 = 2\pi f_0 = \frac{1}{\sqrt{LC}}
\]

\[Z_{in} = R_{in} + jX_{in}\]
Input Impedance

\[
R_{in} = \frac{R}{1 + Q \left( \frac{f}{f_0} - \frac{f_0}{f} \right)}^2
\]

\[R_{in}^{\text{max}} = R_{in} \mid f = f_0 = R\]

\(R\) is the input resistance at the resonance of the patch cavity (the frequency that maximizes \(R_{in}\)).

\(f = f_0\) (resonance of \(RLC\) circuit)
Input Impedance

The input resistance is determined once we know four parameters:

- \( f_0 \): the resonance frequency of the patch cavity
- \( R \): the input resistance at the cavity resonance frequency \( f_0 \)
- \( Q \): the quality factor of the patch cavity
- \( L_p \): the probe inductance

CAD formulas for all of these four parameters have been given earlier.

\[
Z_{in} \approx j\omega L_p + \frac{R}{1 + jQ \left( \frac{f}{f_0} - \frac{f_0}{f} \right)}
\]

\[ (R, f_0, Q) \]
Input Impedance

Typical plot of input impedance

Without probe inductance

With probe inductance
Input Impedance

Results: Input Resistance vs. Frequency

Frequency where the input resistance is maximum ($f_0$): $R_{in} = R$

$\varepsilon_r = 2.2$  \quad W/L = 1.5  \quad L = 3.0$ cm
Results: Input Reactance vs. Frequency

Frequency where the input resistance is maximum ($f_0$)

Shift due to probe reactance

Frequency where the input impedance is real

$\varepsilon_r = 2.2 \quad W/L = 1.5 \quad L = 3.0 \text{ cm}$

Note: “exact” means the cavity model will all infinite modes.
Optimization to get exactly 50 $\Omega$ at the desired resonance frequency:

- Vary the length $L$ first until you find the value that gives an input reactance of zero at the desired frequency.

- Then adjust the feed position $x_0$ to make the real part of the input impedance 50 $\Omega$ at this frequency.
Design a probe-fed rectangular patch antenna on a substrate having a relative permittivity of 2.33 and a thickness of 62 mils (0.1575 cm). (This is Rogers RT Duroid 5870.) Choose an aspect ratio of $W / L = 1.5$. The patch should resonate at the operating frequency of 1.575 GHz (the GPS L1 frequency). Ignore the probe inductance in your design, but account for fringing at the patch edges when you determine the dimensions. At the operating frequency the input impedance should be $50 \, \Omega$ (ignoring the probe inductance). Assume an SMA connector is used to feed the patch along the centerline (at $y = W / 2$), and that the inner conductor of the SMA connector has a radius of 0.635 mm. The copper patch and ground plane have a conductivity of $\sigma = 3.0 \times 10^7 \, \text{S/m}$ and the dielectric substrate has a loss tangent of $\tan \delta = 0.001$.

1) Calculate the following:

- The final patch dimensions $L$ and $W$ (in cm)
- The feed location $x_0$ (distance of the feed from the closest patch edge, in cm)
- The bandwidth of the antenna (SWR $< 2$ definition, expressed in percent)
- The radiation efficiency of the antenna (accounting for conductor, dielectric, and surface-wave loss, and expressed in percent)
- The probe reactance $X_p$ at the operating frequency (in $\Omega$)
- The expected complex input impedance (in $\Omega$) at the operating frequency, accounting for the probe inductance
- Directivity
- Gain
Design Example

Design a probe-fed rectangular patch antenna on a substrate having a relative permittivity of 2.33 and a thickness of 62 mils (0.1575 cm). (This is Rogers RT Duroid 5870.) Choose an aspect ratio of $W/L = 1.5$. The patch should resonate at the operating frequency of 1.575 GHz (the GPS L1 frequency). Ignore the probe inductance in your design, but account for fringing at the patch edges when you determine the dimensions. At the operating frequency the input impedance should be 50 Ω (ignoring the probe inductance). Assume an SMA connector is used to feed the patch along the centerline (at $y = W/2$), and that the inner conductor of the SMA connector has a radius of 0.635 mm. The copper patch and ground plane have a conductivity of $\sigma = 3.0 \times 10^7$ S/m and the dielectric substrate has a loss tangent of $\tan\delta = 0.001$.

2) Find $(f_0, R, X_p, \text{ and } Q)$ and plot the input impedance vs. frequency using the CAD circuit model.

3) Keep $W/L = 1.5$, but now vary the length $L$ of the patch and the feed position until you find the value that makes the input impedance exactly $50 + j(0)$ Ω at 1.575 GHz.
Part 1

Results from the CAD formulas:

1) \( L = 6.071 \text{ cm}, \ W = 9.106 \text{ cm} \)
2) \( x_0 = 1.832 \text{ cm} \)
3) \( BW = 1.23\% \)
4) \( e_r = 82.9\% \)
5) \( X_p = 11.1 \Omega \)
6) \( Z_{in} = 50.0 + j(11.1) \Omega \)
7) \( D = 5.85 \) (7.67 dB)
8) \( G = (D)(e_r) = 4.85 \) (6.86 dB)
Design Example

Part 2

Results from the CAD formulas:

\[ f_0 = 1.575 \times 10^9 \text{ Hz} \]
\[ R = 50 \ \Omega \]
\[ Q = 57.5 \]
\[ X_p = 11.1 \ \Omega \]

\[ Z_{in} \approx jX_p + \frac{R}{1 + jQ \left( \frac{f}{f_0} - \frac{f_0}{f} \right)} \]
Part 3

After optimization:

\[ L = 6.083 \text{ cm} \]
\[ x_0 = 1.800 \text{ cm} \]

\[ Z_{in} = 50 + j(0) \Omega \]
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Three main techniques:

1) **Single feed** with “nearly degenerate” eigenmodes (compact but small CP bandwidth).

2) **Dual feed** with delay line or 90° hybrid phase shifter (broader CP bandwidth but uses more space).

3) **Synchronous subarray technique** (produces high-quality CP due to cancellation effect, but requires even more space).

The techniques will be illustrated with a rectangular patch.
Circular Polarization

Single Feed Method

The feed is on the diagonal. The patch is nearly (but not exactly) square.

$L \approx W$

Basic principle: The two dominant modes (1,0) and (0,1) are excited with equal amplitude, but with a $\pm 45^\circ$ phase.
Circular Polarization

Design equations:

\[ f_{CP} = \frac{f_x + f_y}{2} \]

The optimum CP frequency is the average of the \( x \) and \( y \) resonance frequencies.

\[
\begin{align*}
  f_x &= f_{CP} \left( 1 + \frac{1}{2Q} \right) \\
  f_y &= f_{CP} \left( 1 \pm \frac{1}{2Q} \right)
\end{align*}
\]

Top sign for LHCP, bottom sign for RHCP.

The frequency \( f_{CP} \) is also the resonance frequency: \( Z_{in} = R_{in} = R_x = R_y \)

The resonant input resistance of the CP patch at \( f_{CP} \) is the same as what a linearly-polarized patch fed at the same position would be.
Circular Polarization

Other Variations

**Note:** Diagonal modes are used as degenerate modes

Patch with slot

Patch with truncated corners
Here we compare bandwidths (impedance and axial-ratio):

**Linearly-polarized (LP) patch:**

\[
BW_{SWR}^{LP} = \frac{1}{\sqrt{2Q}} \quad (SWR < 2)
\]

**Circularly-polarized (CP) single-feed patch:**

\[
BW_{SWR}^{CP} = \frac{\sqrt{2}}{Q} \quad (SWR < 2)
\]
\[
BW_{AR}^{CP} = \frac{0.348}{Q} \quad (AR < \sqrt{2} \quad (3\text{dB}))
\]

The axial-ratio bandwidth is **small** when using the single-feed method.

Circular Polarization

Dual-Feed Method

Phase shift realized with delay line:

\[ P + \lambda_g/4 \]
Circular Polarization

Phase shift realized with 90° quadrature hybrid (branchline coupler)

This gives us a higher bandwidth than the simple power divider, but requires a load resistor.
Circular Polarization

Synchronous Rotation

Multiple elements are rotated in space and fed with phase shifts.

Because of symmetry, radiation from higher-order modes (or probes) tends to be reduced, resulting in good cross-pol.
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Circular Patch
Circular Patch

Resonance Frequency

From separation of variables:

\[ E_z = \cos(m\phi) J_m(k_1\rho) \]

\[ k_1 = k_0 \sqrt{\varepsilon_r} \]

\[ J_m = \text{Bessel function of first kind, order } m. \]

\[ \left. \frac{\partial E_z}{\partial \rho} \right|_{\rho=a} = 0 \]

\[ J_m'(k_1a) = 0 \]
Circular Patch

Resonance Frequency

\[ J'_m (k_1 a) = 0 \]

This gives us

\[ k_1 a = x'_{mn} \]

\((n^{th} \text{ root of } J'_m \text{ Bessel function})\)

\[ f_{mn} = \frac{c}{2\pi \sqrt{\varepsilon_r}} x'_{mn} \]

\[ c = \frac{1}{\sqrt{\mu_0 \varepsilon_0}} \]
Circular Patch

Resonance Frequency

Table of values for $x'_{mn}$

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<th>$n/m$</th>
<th>0</th>
<th>1</th>
<th>2</th>
<th>3</th>
<th>4</th>
<th>5</th>
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<td>1</td>
<td>3.832</td>
<td>1.841</td>
<td>3.054</td>
<td>4.201</td>
<td>5.317</td>
<td>5.416</td>
</tr>
</tbody>
</table>

Dominant mode: TM$_{11}$

$$f_{11} = \frac{c}{2\pi a \sqrt{\varepsilon_r}} \ x'_{11}$$

$x'_{11} \approx 1.841$
Circular Patch

Dominant mode: $\text{TM}_{11}$

The circular patch is somewhat similar to a square patch.

Circular patch

Square patch

$W = L$
Circular Patch

Fringing extension

\[ a_e = a + \Delta a \]

\[ f_{11} = \frac{c}{2\pi a_e \sqrt{\varepsilon_r}} \]

"Long/Shen Formula":

\[ a_e = a \sqrt{1 + \frac{2h}{\pi a \varepsilon_r} \left[ \ln \left( \frac{\pi a}{2h} \right) + 1.7726 \right]} \]

or

\[ \Delta a \approx \frac{h}{\pi \varepsilon_r} \left[ \ln \left( \frac{\pi a}{2h} \right) + 1.7726 \right] \]

Circular Patch

Patterns

(The patterns are based on the magnetic current model.)

\[ E_z(\rho, \phi) = \cos\phi \left( \frac{J_1(k_1\rho)}{J_1(k_1a)} \right) \left( \frac{1}{h} \right) \]

\[ k_1 = k_0 \sqrt{\varepsilon_r} \]

(The edge voltage has a maximum of one volt.)
Circular Patch

Patterns

\[ E^R_\theta (r, \theta, \phi) = 2\pi a \frac{E_0}{\eta_0} \tanh(k_z h) \cos \phi \ J'_1(k_0 a \sin \theta) Q(\theta) \]

\[ E^R_\phi (r, \theta, \phi) = -2\pi a \frac{E_0}{\eta_0} \tanh(k_z h) \sin \phi \left( \frac{J_1(k_0 a \sin \theta)}{k_0 a \sin \theta} \right) P(\theta) \]

where

\[ \tanh(x) = \frac{\tan(x)}{x} \]

\[ P(\theta) = \cos \theta \left( 1 - \Gamma^{TE} (\theta) \right) = \cos \theta \left[ \frac{-2jN(\theta)}{\tan(k_0 h N(\theta)) - jN(\theta) \sec \theta} \right] \]

\[ Q(\theta) = 1 - \Gamma^{TM} (\theta) = \frac{-2j \left( \frac{\varepsilon_r}{N(\theta)} \right) \cos \theta}{\tan(k_0 h N(\theta)) - j \frac{\varepsilon_r}{N(\theta)} \cos \theta} \]

\[ E_0 = \left( -j \omega \mu_0 \right) e^{-jk_0 r} \]

\[ N(\theta) = \sqrt{\varepsilon_r - \sin^2(\theta)} \]

Note: To account for lossy substrate, use \( \varepsilon_r \rightarrow \varepsilon_{rc} = \varepsilon_r (1 - j \tan \delta) \)
Circular Patch

Input Resistance

\[ R_{in} \approx R_{edge} \left[ \frac{J_1^2(k_1\rho_0)}{J_1^2(k_1a)} \right] \]

\[ k_1 = k_0 \sqrt{\varepsilon_r} \]
Circular Patch

Input Resistance (cont.)

\[
R_{\text{edge}} = \left[ \frac{1}{2P_{sp}} \right] e_r
\]

where

\[
P_{sp} = \frac{\pi}{8\eta_0} (k_0 a)^2 \int_0^{\pi/2} \tan^2 (k_0 hN(\theta)) \sin^2 (k_0 a \sin \theta) + |P(\theta)|^2 J_{inc}^2 (k_0 a \sin \theta) \sin \theta \, d\theta
\]

\[J_{inc}(x) = J_1(x) / x\]

\[
P_{sp} = \text{power radiated into space by circular patch with maximum edge voltage of one volt.}
\]

\[
e_r = \text{radiation efficiency}
\]
Circular Patch

Input Resistance (cont.)

CAD Formula:

\[ P_{sp} = \frac{\pi}{8\eta_0} (k_0 a)^2 I_c \]

\[ I_c = \frac{4}{3} p_c \]

\[ p_c = \sum_{k=0}^{6} (k_0 a)^{2k} e^{2k} \]

\[ e_0 = 1 \]
\[ e_2 = -0.400000 \]
\[ e_4 = 0.0785710 \]
\[ e_6 = -7.27509 \times 10^{-3} \]
\[ e_8 = 3.81786 \times 10^{-4} \]
\[ e_{10} = -1.09839 \times 10^{-5} \]
\[ e_{12} = 1.47731 \times 10^{-7} \]
Outline

- Overview of microstrip antennas
- Feeding methods
- Basic principles of operation
- General characteristics
- CAD Formulas
- Radiation pattern
- Input Impedance
- Circular polarization
- Circular patch
- Improving bandwidth
- Miniaturization
- Reducing surface waves and lateral radiation
Improving Bandwidth

Some of the techniques that have been successfully developed are illustrated here.

The literature may be consulted for additional designs and variations.
Improving Bandwidth

Probe Compensation

L-shaped probe:

As the substrate thickness increases the probe inductance limits the bandwidth – so we compensate for it.

Capacitive “top hat” on probe:

Top view
Improving Bandwidth

SSFIP: Strip Slot Foam Inverted Patch (a version of the ACP).

- Bandwidths greater than 25% have been achieved.
- Increased bandwidth is due to the thick foam substrate and also a dual-tuned resonance (patch+slot).

**Note:** There is no probe inductance to worry about here.

Improving Bandwidth

Stacked Patches

- Bandwidth increase is due to thick low-permittivity antenna substrates and a dual or triple-tuned resonance.
- Bandwidths of 25% have been achieved using a probe feed.
- Bandwidths of 100% have been achieved using an ACP feed.
Improving Bandwidth

Stacked Patches

Bandwidth \( (S_{11} = -10 \text{ dB}) \) is about 100%

(Stacked patch with ACP feed)

(Photo courtesy of Dr. Rodney B. Waterhouse)
Improving Bandwidth

Stacked Patches

Stacked patch with ACP feed

Two extra loops are observed on the Smith chart.

(Photo courtesy of Dr. Rodney B. Waterhouse)
Radiating Edges Gap Coupled Microstrip Antennas (REGCOMA).

Non-Radiating Edges Gap Coupled Microstrip Antennas (NEGCOMA)

Four-Edges Gap Coupled Microstrip Antennas (FEGCOMA)

Bandwidth improvement factor: REGCOMA: 3.0, NEGCOMA: 3.0, FEGCOMA: 5.0?

Mush of this work was pioneered by K. C. Gupta.
Improving Bandwidth

Direct-Coupled Patches

- Radiating Edges Direct Coupled Microstrip Antennas (REDCOMA).
- Non-Radiating Edges Direct Coupled Microstrip Antennas (NEDCOMA)
- Four-Edges Direct Coupled Microstrip Antennas (FEDCOMA)

Bandwidth improvement factor:
REDCOMA: 5.0, NEDCOMA: 5.0, FEDCOMA: 7.0
The introduction of a U-shaped slot can give a significant bandwidth (10%-40%).

(This is due to a double resonance effect, with two different modes.)

Improving Bandwidth

Double U-Slot

A 44% bandwidth was achieved.

A modification of the U-slot patch.

A bandwidth of 34% was achieved (40% using a capacitive “washer” to compensate for the probe inductance).

Multi-Band Antennas

A multi-band antenna is sometimes more desirable than a broadband antenna, if multiple narrow-band channels are to be covered.

General Principle:

Introduce multiple resonance paths into the antenna.
Multi-Band Antennas

Dual-band E patch

Dual-band patch with parasitic strip
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Miniaturization

- High Permittivity
- Quarter-Wave Patch
- PIFA
- Capacitive Loading
- Slots
- Meandering

**Note:** Miniaturization usually comes at a price of reduced bandwidth!

Usually, bandwidth is proportional to the volume of the patch cavity, as we will see in the examples.
The smaller patch has about one-fourth the bandwidth of the original patch.

(Bandwidth is inversely proportional to the permittivity.)
The new patch has about one-half the bandwidth of the original patch.

Neglecting losses:

\[ Q = \omega_0 \frac{U_s}{P_r} \quad U'_s = \frac{U_s}{2} \quad P'_r = \frac{P_r}{4} \]

\[ Q' = 2Q \]

Note: 1/2 of the radiating magnetic current
Miniaturationization

Smaller Quarter-Wave patch

A quarter-wave patch with the same aspect ratio $W/L$ as the original patch

The new patch has about one-half the bandwidth of the original quarter-wave patch, and hence one-fourth the bandwidth of the regular patch.

(Bandwidth is proportional to the patch width.)
Fewer vias actually gives more miniaturization!

(The edge has a larger inductive impedance: explained on the next slide.)
The Smith chart provides a simple explanation for the length reduction.

Miniaturization
Quarter-Wave Patch with Fewer Vias
A single shorting strip or via is used.

This antenna can be viewed as a limiting case of the via-loaded patch, or as an $LC$ resonator.
The capacitive loading allows for the length of the PIFA to be reduced.

\[ \omega_0 = \frac{1}{\sqrt{LC}} \]
The patch has a monopole-like pattern

The patch operates in the \((0,0)\) mode, as an \(LC\) resonator

(Hao Xu Ph.D. dissertation, University of Houston, 2006)
Miniaturization

Circular Patch Loaded with Vias

Example: Circular Patch Loaded with 2 Vias

Unloaded: resonance frequency = 5.32 GHz.

(Miniaturization factor = 4.8)
The slot forces the current to flow through a longer path, increasing the effective dimensions of the patch.
Meandering forces the current to flow through a longer path, increasing the effective dimensions of the patch.

Meandering also increases the capacitance of the PIFA line.
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Reducing surface waves and lateral radiation
Reducing Surface and Lateral Waves

Reduced Surface Wave (RSW) Antenna

Reducing Surface and Lateral Waves

Reducing surface-wave excitation and lateral radiation reduces edge diffraction and mutual coupling.

- Edge diffraction degrades the radiation pattern on a finite ground plane.
- Mutual coupling causes an desirable coupling between antennas.
Reducing surface-wave excitation and lateral radiation reduces edge diffraction.

- **Space-wave radiation** (desired)
- **Lateral radiation** (undesired)
- **Diffracted field at edge**
- **Surface waves** (undesired)
Reducing the Surface Wave Excitation

**TM$_{11}$ mode:**

\[
E_z(\rho, \phi) = V_0 \left( \frac{-1}{hJ_1(k_1a)} \right) \cos \phi J_1(k_1\rho)
\]

At edge:

\[
E_z = - \frac{V_0}{h} \cos \phi
\]

\[
M_s = -\hat{n} \times \overline{E} = -\hat{\rho} \times (\hat{z} E_z)
\]

\[
M_{s\phi}(\phi) = E_z(a, \phi)
\]

\[
M_{s\phi} = -\frac{V_0}{h} \cos \phi
\]
Reduction Surface and Lateral Waves

Surface-Wave Excitation:

\[ M_{s\phi} = -\frac{V_0}{h} \cos \phi \]

\[ E_z^{TM_0} = A_{TM_0} \cos \phi H_1^{(2)}(\beta_{TM_0} \rho) e^{-jk_z0z} \]

\[ (z > h) \]

\[ A_{TM_0} = AJ'_1(\beta_{TM_0} a) \]

Set

\[ J'_1(\beta_{TM_0} a) = 0 \]
Reducing Surface and Lateral Waves

For TM$_{11}$ mode: $x'_{11} \approx 1.841$

Hence

Note: $\beta_{TM_0} < k_1$ → The RSW patch is too big to be resonant.
The radius $a$ is chosen to make the patch resonant:

$$\beta_{TM_0} b = 1.841$$

$$\frac{J_1(k_1a)}{Y_1(k_1a)} = \frac{J_1'(\frac{k_1 x'_1}{\beta_{TM_0}})}{Y_1'(\frac{k_1 x'_1}{\beta_{TM_0}})}$$
Reducing Surface and Lateral Waves

Reducing the Lateral Radiation

Assume no substrate outside of patch (or very thin substrate):

**Space-Wave Field:**

\[ E_{z}^{SP} = A_{SP} \cos \phi \left( \frac{1}{\rho} \right) e^{-jk_{0}\rho} \quad (z = h) \]

\[ A_{SP} = CJ'_{1}(k_{0}a) \]

Set \( J'_{1}(k_{0}a) = 0 \) \( \Rightarrow \) \( k_{0}a = 1.841 \)
For a thin substrate:

\[ \beta_{TM_0} \approx k_0 \]

The same design reduces both surface-wave fields and lateral-radiation fields.

\[ k_0 a = 1.841 \]

\[ \frac{2a}{\lambda_0} = 0.586 \]

**Note:** The diameter of the RSW antenna is found from

**Note:** The size is approximately independent of the permittivity (the patch cannot be miniaturized by choosing a higher permittivity!).

Reducing Surface and Lateral Waves
Reducing Surface and Lateral Waves
E-plane Radiation Patterns

Measurements were taken on a 1 m diameter circular ground plane at 1.575 GHz.
Reducing surface-wave excitation and lateral radiation reduces mutual coupling.
Reducing Surface and Lateral Waves

Reducing surface-wave excitation and lateral radiation reduces mutual coupling.

References

General references about microstrip antennas:


General references about microstrip antennas (cont.):


More information about the CAD formulas presented here for the rectangular patch may be found in:


References devoted to broadband microstrip antennas:


The End