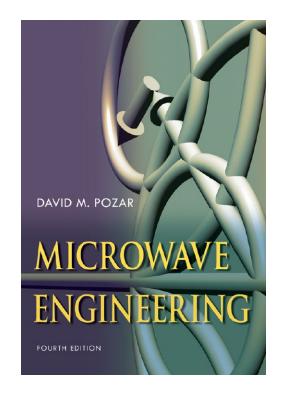
Adapted from notes by Prof. Jeffery T. Williams

ECE 5317-6351 Microwave Engineering

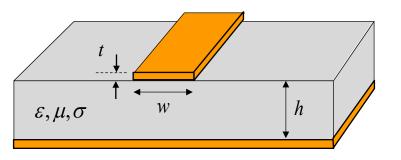
Fall 2019

Prof. David R. Jackson Dept. of ECE

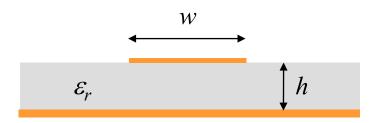


Notes 11

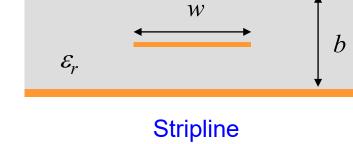
Waveguiding Structures Part 6: Planar Transmission Lines

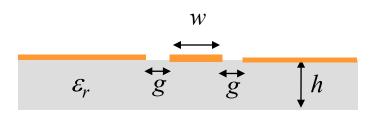


Planar Transmission Lines

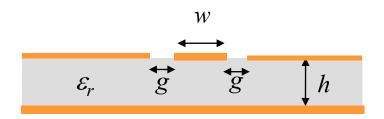


Microstrip

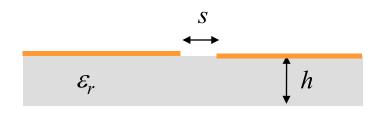




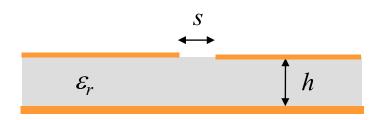
Coplanar Waveguide (CPW)



Conductor-backed CPW

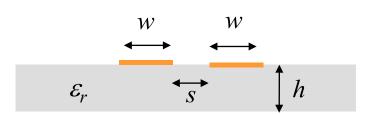


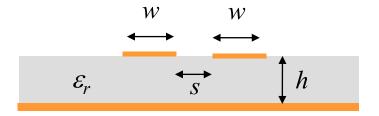
Slotline



Conductor-backed Slotline

Planar Transmission Lines (cont.)





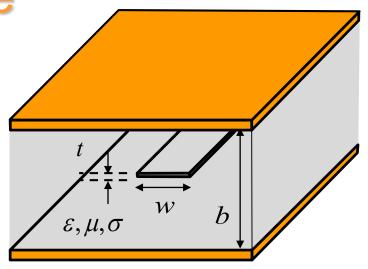
Coplanar Strips (CPS)

Conductor-backed CPS

- Stripline is a planar version of coax.
- Coplanar strips (CPS) is a planar version of twin lead.

Stripline

- Common on circuit boards
- Fabricated with two circuit boards
- Homogenous dielectric (perfect TEM mode*)

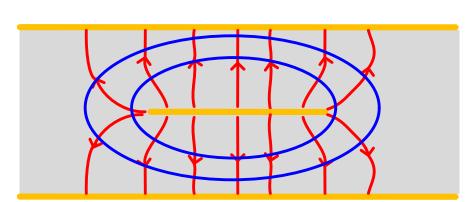


TEM mode

(also TE & TM Modes at high frequency)

Field structure for TEM mode:

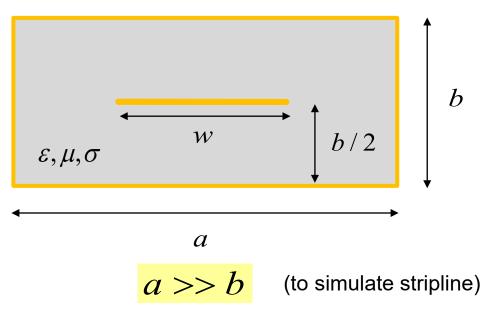
Electric Field — Magnetic Field —



^{*} The mode is a <u>perfect</u> TEM mode if there is no conductor loss.

- Analysis of stripline is not simple.
- TEM mode fields can be obtained from an electrostatic analysis (e.g., conformal mapping).

A <u>closed</u> stripline structure is analyzed in the Pozar book by using an approximate numerical method:



Conformal Mapping Solution (R. H. T. Bates)

Exact solution (for t = 0):

$$K =$$
complete elliptic integral of the first kind

$$K(k) \equiv \int_{0}^{\pi/2} \frac{1}{\sqrt{1 - k^2 \sin^2 \theta}} d\theta$$

$$k = \operatorname{sech}\left(\frac{\pi w}{2b}\right)$$

$$k' = \tanh\left(\frac{\pi w}{2b}\right)$$

R. H. T. Bates, "The characteristic impedance of the shielded slab line," *IEEE Trans. Microwave Theory* and Techniques, vol. 4, pp. 28-33, Jan. 1956.

Curve fitting this exact solution:

$$Z_0 = \left(\frac{\eta_0}{4\sqrt{\varepsilon_r}}\right) \frac{b}{w_e + \frac{\ln(4)}{\pi}b}$$
Note: $\frac{\ln(4)}{\pi} = 0.441$

Effective width

Fringing term

$$\frac{w_e}{b} = \frac{w}{b} - \begin{cases} 0 & \text{; for } \frac{w}{b} \ge 0.35\\ \left(0.35 - \frac{w}{b}\right)^2 & \text{; for } 0.1 \le \frac{w}{b} \le 0.35 \end{cases}$$

Note:
$$Z_0^{\text{ideal}} = \frac{1}{2} \eta \left(\frac{b/2}{w} \right) = \frac{\eta b}{4w} = \frac{\eta_0 b}{4w\sqrt{\varepsilon_r}}$$
 The factor of 1/2 in front is from the parallel combination of two ideal PPWs.

Inverting this solution to find w for given Z_0 :

$$\frac{w}{b} = \begin{cases} X; & \text{for } \sqrt{\varepsilon_r} Z_0 \le 120 \ \left[\Omega\right] \\ \\ 0.85 - \sqrt{0.6 - X} & ; & \text{for } \sqrt{\varepsilon_r} Z_0 \ge 120 \ \left[\Omega\right] \end{cases}$$

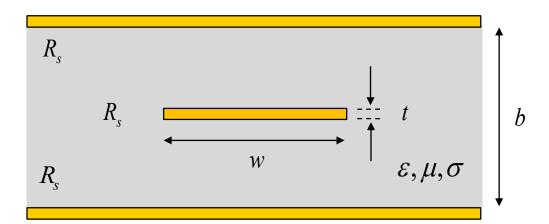
$$X \equiv \frac{\eta_0}{4\sqrt{\varepsilon_r} Z_0} - \frac{\ln(4)}{\pi}$$

Attenuation

Dielectric Loss:

$$\alpha_d = k'' \approx \frac{k'}{2} \tan \delta_d \approx \frac{k_0 \sqrt{\varepsilon_r}}{2} \tan \delta_d$$
 (TEM formula)

$$k = k' - jk'' = \omega \sqrt{\mu_0 \varepsilon_c}$$
 $\varepsilon_c = \varepsilon - j\frac{\sigma}{\omega}$ $\tan \delta_d = \frac{\varepsilon_c''}{\varepsilon_c'}$



Conductor Loss:

$$\alpha_c \approx \begin{cases} (2.7 \times 10^{-3}) \frac{4R_s \varepsilon_r Z_0}{\eta_0(b-t)} A \; ; \; \; \text{for} \, \sqrt{\varepsilon_r} Z_0 \leq 120 \left[\Omega\right] \; \; \left(\text{wider strips}\right) \\ 0.16 \bigg(\frac{R_s}{Z_0 b}\bigg) B \; ; \; \; \text{for} \, \sqrt{\varepsilon_r} Z_0 \geq 120 \left[\Omega\right] \; \; \left(\text{narrower strips}\right) \end{cases} \\ \text{where} \\ A = 1 + 2 \frac{w}{(b-t)} + \frac{1}{\pi} \bigg(\frac{b+t}{b-t}\bigg) \ln \bigg(\frac{2b-t}{t}\bigg) \\ B = 1 + \frac{b}{\bigg(\frac{w}{2} + 0.7t\bigg)} + \bigg(\frac{1}{2} + 0.414 \frac{t}{w} + \frac{1}{2\pi} \ln \bigg(4\pi \frac{w}{t}\bigg)\bigg) \end{cases}$$

Note: We cannot let $t \to 0$ when we calculate the conductor loss.

Note about conductor attenuation:

It is necessary to assume a nonzero conductor thickness in order to accurately calculate the conductor attenuation.

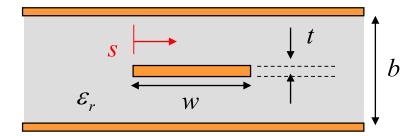
The perturbational method predicts an infinite attenuation if a zero thickness is assumed.

$$t=0$$
: $J_{sz} \propto \frac{1}{\sqrt{S}}$ as $s \to 0$

$$\alpha_c = \frac{P_l(0)}{2P_0}$$

$$P_l(0) = \frac{R_s}{2} \int_{C_1 + C_2} |\underline{J}_s|^2 d\ell \rightarrow \infty$$





Practical note:

A standard metal thickness for PCBs is 0.7 [mils] (17.5 $[\mu m]$), called "half-ounce copper".

$$1 \text{ mil} = 0.001 \text{ inch}$$

Microstrip

 ε, μ, σ $\downarrow h$

Inhomogeneous dielectric

 \Rightarrow No TEM mode

Note: Pozar uses (W, d)

TEM mode would require $k_z = k$ in each region, but k_z must be unique!

- Requires advanced analysis techniques
- Exact fields are <u>hybrid</u> modes (E_z and H_z)

For $h/\lambda_0 << 1$, the dominant mode is <u>quasi-TEM</u>.

Part of the field lines are in air, and part of the field lines are inside the substrate.

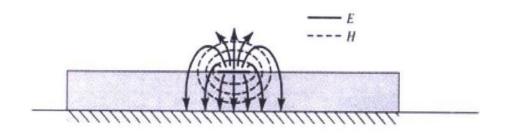
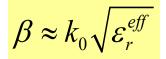


Figure from Pozar book

Note:

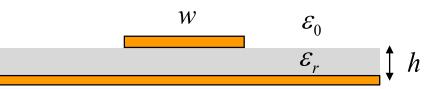
The flux lines get more concentrated in the substrate region as the frequency increases.

Equivalent TEM problem:

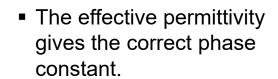


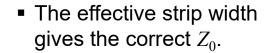
$$\Rightarrow \varepsilon_r^{eff} = \left(\frac{\beta}{k_0}\right)^2$$

Actual problem



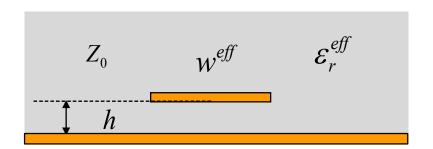






$$Z_0 = Z_0^{air} / \sqrt{\mathcal{E}_r^{\it eff}}$$

$$Z_0^{\it air}: \ \mathcal{E}_r^{\it eff} o 1$$
 $\left({
m since} \ Z_0 = \sqrt{L/C}
ight)$



Equivalent TEM problem

Effective permittivity:

$$\varepsilon_r^{eff} = \frac{\varepsilon_r + 1}{2} + \frac{\varepsilon_r - 1}{2} \left(\frac{1}{\sqrt{1 + 12\frac{h}{w}}} \right)$$

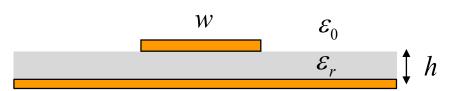
Note:

This formula ignores "dispersion", i.e., the fact that the effective permittivity is actually a function of frequency.

Limiting cases:

$$w/h \to 0: \quad \varepsilon_r^{eff} \to \frac{\varepsilon_r + 1}{2}$$
 (narrow strip)

$$w/h \to \infty$$
: $\varepsilon_r^{eff} \to \varepsilon_r$ (wide strip)



Characteristic Impedance:

$$Z_0 = \begin{cases} \frac{60}{\sqrt{\varepsilon_r^{\textit{eff}}}} \, \ln\!\left(\frac{8h}{w} + \frac{w}{4h}\right); & \text{for } \frac{w}{h} \leq 1 \\ \frac{\eta_0}{\sqrt{\varepsilon_r^{\textit{eff}}} \left(\frac{w}{h} + 1.393 + 0.667 \ln\!\left(\frac{w}{h} + 1.444\right)\right)} & ; & \text{for } \frac{w}{h} \geq 1 \end{cases}$$

Note:

This formula ignores the fact that the characteristic impedance is actually a function of frequency.

Inverting this solution to find w for a given Z_0 :

$$\frac{w}{h} = \begin{cases} \frac{8e^A}{e^{2A} - 2} ; & \text{for } \frac{w}{h} \le 2\\ \frac{2}{\pi} \left[B - 1 - \ln(2B - 1) + \frac{\varepsilon_r - 1}{2\varepsilon_r} \left(\ln(B - 1) + 0.39 - \frac{0.61}{\varepsilon_r} \right) \right]; & \text{for } \frac{w}{h} \ge 2 \end{cases}$$

where

$$A = \frac{Z_0}{60} \sqrt{\frac{\varepsilon_r + 1}{2}} + \frac{\varepsilon_r - 1}{\varepsilon_r + 1} \left(0.33 + \frac{0.11}{\varepsilon_r} \right)$$

$$B = \frac{\eta_0 \pi}{2Z_0 \sqrt{\varepsilon_r}}$$

More accurate formulas for characteristic impedance that account for dispersion (frequency variation) and conductor thickness:

$$Z_{0}(f) = Z_{0}(0) \left(\frac{\varepsilon_{r}^{eff}(f) - 1}{\varepsilon_{r}^{eff}(0) - 1} \right) \sqrt{\frac{\varepsilon_{r}^{eff}(0)}{\varepsilon_{r}^{eff}(f)}}$$

$$Z_{0}(0) = \frac{\eta_{0}}{\sqrt{\varepsilon_{r}^{eff}(0)} \left[(w'/h) + 1.393 + 0.667 \ln((w'/h) + 1.444) \right]} \quad (w/h \ge 1)$$

$$w' = w + \frac{t}{\pi} \left(1 + \ln \left(\frac{2h}{t} \right) \right)$$

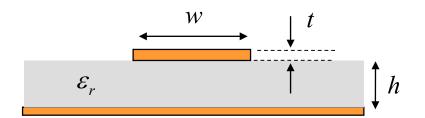
$$\mathcal{E}_{r}$$

where

$$\varepsilon_r^{eff}(f) = \left(\sqrt{\varepsilon_r^{eff}(0)} + \frac{\sqrt{\varepsilon_r} - \sqrt{\varepsilon_r^{eff}(0)}}{1 + 4F^{-1.5}}\right)^2 \qquad (w/h \ge 1)$$

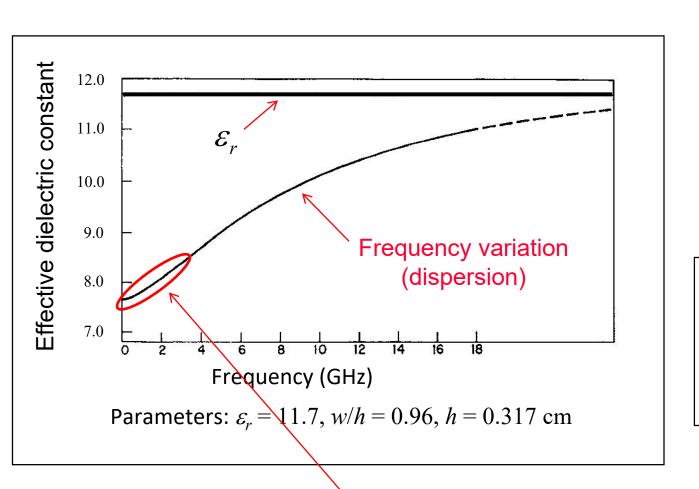
$$\varepsilon_r^{eff}\left(0\right) = \frac{\varepsilon_r + 1}{2} + \left(\frac{\varepsilon_r - 1}{2}\right) \left(\frac{1}{\sqrt{1 + 12(h/w)}}\right) - \left(\frac{\varepsilon_r - 1}{4.6}\right) \left(\frac{t/h}{\sqrt{w/h}}\right)$$

$$F = 4\left(\frac{h}{\lambda_0}\right)\sqrt{\varepsilon_r - 1}\left(0.5 + \left(1 + 0.868\ln\left(1 + \frac{w}{h}\right)\right)^2\right)$$



As
$$f \to 0$$
:
$$\varepsilon_r^{eff}(f) \to \varepsilon_r^{eff}(0)$$

As
$$f \to \infty$$
:
$$\varepsilon_r^{eff}(f) \to \varepsilon_r$$



Note:

The phase velocity is a function of frequency, which causes pulse distortion.

Quasi-TEM region

$$\varepsilon_r^{eff} = \left(\frac{\beta}{k_0}\right)^2$$

$$v_p = \frac{c}{\sqrt{\varepsilon_r^{eff}}}$$

"A frequency-dependent solution for microstrip transmission lines," E. J. Denlinger, *IEEE Trans. Microwave Theory and Techniques*, Vol. 19, pp. 30-39, Jan. 1971.

Note:

The flux lines get more concentrated in the substrate region as the frequency increases.

Attenuation

Dielectric loss:

"filling factor"

$$\alpha_d \approx \frac{k_0 \sqrt{\varepsilon_r}}{2} \tan \delta_d \left[\sqrt{\frac{\varepsilon_r}{\varepsilon_r^{eff}}} \frac{\left(\varepsilon_r^{eff} - 1\right)}{\left(\varepsilon_r - 1\right)} \right]$$

$$\varepsilon_r^{eff} \to 1: \quad \alpha_d \to 0$$

Conductor loss:

$$\alpha_c \approx \frac{R_s}{Z_0 w} \approx \frac{R_s}{\eta h}$$

$$\left(Z_0 \approx \frac{\eta h}{w}\right)$$

$$\varepsilon_r^{eff} \to \varepsilon_r: \quad \alpha_d \to \frac{k_0 \sqrt{\varepsilon_r}}{2} \tan \delta_d$$

very crude ("parallel-plate") approximation (More accurate formulas are given on next slide.)

More accurate formulas for conductor attenuation:

$$\frac{1}{2\pi} < \frac{w}{h} \le 2$$

$$\frac{1}{2\pi} < \frac{w}{h} \le 2 \qquad \alpha_c = \left(\frac{R_s}{hZ_0}\right) \left(\frac{1}{2\pi}\right) \left[1 - \left(\frac{w'}{4h}\right)^2\right] \left[1 + \frac{h}{w'} + \frac{h}{\pi w'} \left(\ln\left(\frac{2h}{t}\right) - \frac{t}{h}\right)\right]$$

$$\frac{w}{h} \ge 2$$

$$\frac{w}{h} \ge 2 \qquad \alpha_c = \left(\frac{R_s}{hZ_0}\right) \left[\frac{w'}{h} + \frac{2}{\pi} \ln\left(2\pi e\left(\frac{w'}{2h} + 0.94\right)\right)\right]^{-2} \left[\frac{w'}{h} + \frac{w'/(\pi h)}{\frac{w'}{2h} + 0.94}\right] \left[1 + \frac{h}{w'} + \frac{h}{\pi w'}\left(\ln\left(\frac{2h}{t}\right) - \frac{t}{h}\right)\right]$$

This is the number e = 2.71828 multiplying the term in parenthesis.

$$w' = w + \frac{t}{\pi} \left(1 + \ln \left(\frac{2h}{t} \right) \right)$$

$$\mathcal{E}_r$$

$$w' = w + \frac{t}{\pi} \left(1 + \ln \left(\frac{2h}{t} \right) \right)$$

REFERENCES

- L. G. Maloratsky, Passive RF and Microwave Integrated Circuits, Elsevier, 2004.
- I. Bahl and P. Bhartia, Microwave Solid State Circuit Design, Wiley, 2003.
- R. A. Pucel, D. J. Masse, and C. P. Hartwig, "Losses in Microstrip," *IEEE Trans. Microwave Theory and Techniques*, pp. 342-350, June 1968.
- R. A. Pucel, D. J. Masse, and C. P. Hartwig, "Corrections to 'Losses in Microstrip'," *IEEE Trans. Microwave Theory and Techniques*, Dec. 1968, p. 1064.



This is a public-domain software for calculating the properties of some common planar transmission lines.

https://www.awr.com/software/options/tx-line





