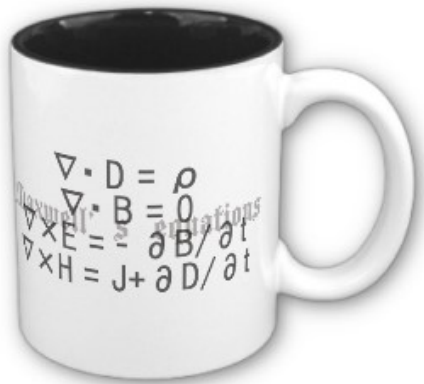


# ECE 6340

## Intermediate EM Waves

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**Notes 10**

# Wave Equation

Assume  $\underline{J}^i = \underline{0}$  ,  $\underline{M}^i = \underline{0}$

The region is source free and homogeneous.

so  $\nabla \times \underline{E} = -j\omega\mu\underline{H}$

$$\nabla \times \underline{H} = j\omega\varepsilon_c \underline{E}$$

Then we have  $\nabla \times (\nabla \times \underline{E}) = -j\omega\mu(\nabla \times \underline{H})$   
 $= -j\omega\mu(j\omega\varepsilon_c \underline{E})$

Use  $k^2 = \omega^2 \mu \varepsilon_c$

Hence

$$\nabla \times (\nabla \times \underline{E}) - k^2 \underline{E} = \underline{0}$$

Vector Wave Equation

# Wave Equation (cont.)

Similarly,  $\nabla \times (\nabla \times \underline{H}) - k^2 \underline{H} = \underline{0}$

Taking the divergence,

$$\nabla \times (\nabla \times \underline{E}) - k^2 \underline{E} = \underline{0}$$



$$\nabla \cdot [\nabla \times (\nabla \times \underline{E})] - k^2 \nabla \cdot \underline{E} = 0$$

Hence,  $\nabla \cdot \underline{E} = 0$

Hence  $\rho_v = 0$  in the time-harmonic steady state for a homogeneous material, provided  $\omega \neq 0$ .

The zero-divergence condition is built in to the vector wave equation.

# Vector Laplacian

Recall  $\nabla^2 \underline{V} \equiv \nabla(\nabla \cdot \underline{V}) - \nabla \times (\nabla \times \underline{V})$

In rectangular coordinates,

$$\nabla^2 \underline{V} = \hat{x} \nabla^2 V_x + \hat{y} \nabla^2 V_y + \hat{z} \nabla^2 V_z$$

The vector wave equation can thus be written as

$$\underbrace{\nabla(\nabla \cdot \underline{E}) - \nabla^2 \underline{E} - k^2 \underline{E}}_{\nabla \times \nabla \times \underline{E}} = 0$$

or

$$\nabla^2 \underline{E} + k^2 \underline{E} = \underline{0}$$

“Vector Helmholtz Equation”

# Vector Laplacian (cont.)

Take the rectangular components:

$$\nabla^2 E_x + k^2 E_x = 0$$

$$\nabla^2 E_y + k^2 E_y = 0$$

$$\nabla^2 E_z + k^2 E_z = 0$$

Note:  $\nabla \cdot \underline{E} = 0$  is **not implied** from the  
vector Helmholtz Equation.

(see the example later)

# Waveguide Fields

From the “guided-wave theorem,” the fields of a guided wave varying as  $\exp(-jk_z z)$  may be written in terms of the longitudinal ( $z$ ) components as

$$E_x = \frac{-j\omega\mu}{k^2 - k_z^2} \frac{\partial H_z}{\partial y} - \frac{jk_z}{k^2 - k_z^2} \frac{\partial E_z}{\partial x}$$

$$E_y = \frac{j\omega\mu}{k^2 - k_z^2} \frac{\partial H_z}{\partial x} - \frac{jk_z}{k^2 - k_z^2} \frac{\partial E_z}{\partial y}$$

$$H_x = \frac{j\omega\epsilon_c}{k^2 - k_z^2} \frac{\partial E_z}{\partial y} - \frac{jk_z}{k^2 - k_z^2} \frac{\partial H_z}{\partial x}$$

$$H_y = \frac{-j\omega\epsilon_c}{k^2 - k_z^2} \frac{\partial E_z}{\partial x} - \frac{jk_z}{k^2 - k_z^2} \frac{\partial H_z}{\partial y}$$

# Waveguide Fields (cont.)

These may be written more compactly as

$$\underline{E}_t = \frac{j\omega\mu}{k^2 - k_z^2} (\hat{z} \times \nabla_t H_z) - \frac{jk_z}{k^2 - k_z^2} (\nabla_t E_z)$$

$$\underline{H}_t = \frac{-j\omega\epsilon_c}{k^2 - k_z^2} (\hat{z} \times \nabla_t E_z) - \frac{jk_z}{k^2 - k_z^2} (\nabla_t H_z)$$

# Waveguide Equations (cont.)

For the special case of  $TE_z$  and  $TM_z$  fields, we can combine these as follows:

$$\begin{array}{l}
 \left. \begin{array}{l}
 \underline{E}_t = \frac{j\omega\mu}{k^2 - k_z^2} (\hat{z} \times \nabla_t H_z) - \frac{jk_z}{k^2 - k_z^2} (\nabla_t E_z) \\
 \underline{H}_t = \frac{-j\omega\epsilon_c}{k^2 - k_z^2} (\hat{z} \times \nabla_t E_z) - \frac{jk_z}{k^2 - k_z^2} (\nabla_t H_z)
 \end{array} \right\} TE_z
 \end{array}$$

$E_z = 0$

Combine

$E_z = 0$



# Waveguide Equations (cont.)

$$\begin{aligned}\underline{E}_t &= \frac{j\omega\mu}{k^2 - k_z^2} \hat{\underline{z}} \times \left( - \left( \frac{k^2 - k_z^2}{jk_z} \right) \underline{H}_t \right) \\ &= - \left( \frac{\omega\mu}{k_z} \right) (\hat{\underline{z}} \times \underline{H}_t)\end{aligned}$$

Define the wave impedance:  $Z^{TE} = \left( \frac{\omega\mu}{k_z} \right)$

Then we have  $\underline{E}_t = -Z^{TE} (\hat{\underline{z}} \times \underline{H}_t)$

# Waveguide Equations (cont.)

Summary for  $TE_z$  and  $TM_z$  cases:

$$\underline{E}_t = -Z^{TE} (\hat{z} \times \underline{H}_t), \quad TE_z \text{ Fields}$$

$$\underline{E}_t = -Z^{TM} (\hat{z} \times \underline{H}_t), \quad TM_z \text{ Fields}$$

$$Z^{TE} = \left( \frac{\omega\mu}{k_z} \right)$$

$$Z^{TM} = \left( \frac{k_z}{\omega\epsilon_c} \right)$$

**Note:** For a wave going in the  $-z$  direction, we replace the minus sign with a plus sign.

# Waveguide Equations (cont.)

Taking the cross product of both sides with  $\hat{\underline{z}}$  gives us

$$\underline{H}_t = +\frac{1}{Z^{TE}}(\hat{\underline{z}} \times \underline{E}_t), \quad \text{TE}_z \text{ Fields}$$

$$\underline{H}_t = +\frac{1}{Z^{TM}}(\hat{\underline{z}} \times \underline{E}_t), \quad \text{TM}_z \text{ Fields}$$

Note: For any transverse vector we have

$$\hat{\underline{z}} \times (\hat{\underline{z}} \times \underline{V}_t) = -\underline{V}_t$$

**Note:** For a wave going in the  $-z$  direction, we replace the plus sign with a minus sign.

# Waveguide Equations (cont.)

For the special case of a TEM wave:

$$k_z = k$$

$$Z^{TE} = \frac{\omega\mu}{k_z} = \frac{\omega\mu}{k} = \frac{\omega\mu}{\omega\sqrt{\mu\epsilon_c}} = \eta$$

$$Z^{TM} = \frac{k_z}{\omega\epsilon_c} = \frac{k}{\omega\epsilon_c} = \frac{\omega\sqrt{\mu\epsilon_c}}{\omega\epsilon_c} = \eta$$

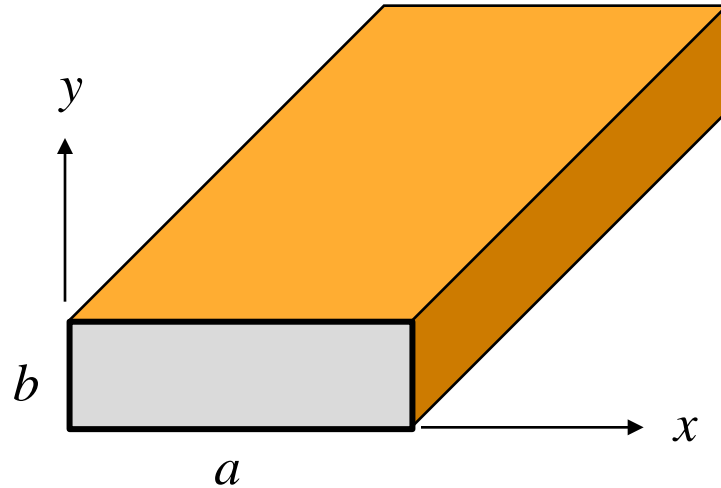
$$Z^{TE} = Z^{TM} = \eta \quad (\text{TEM wave})$$

# Example: Spurious $TE_{00}$ Mode

Rectangular Waveguide

$$\nabla^2 E_z + k^2 E_z = 0$$

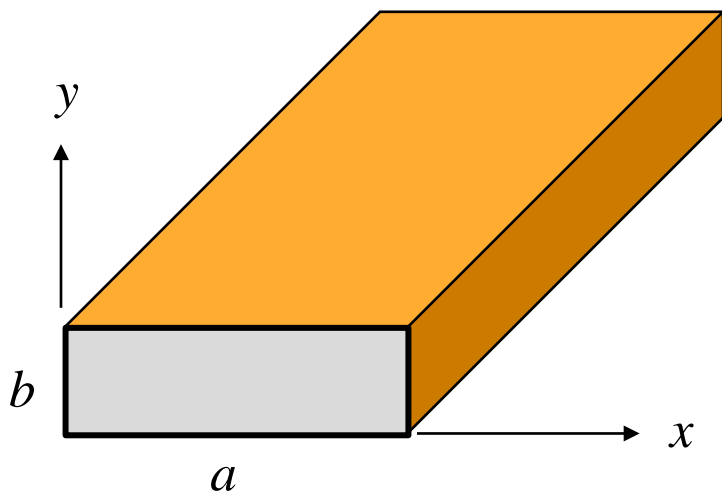
$$\nabla^2 H_z + k^2 H_z = 0$$



In this example we show that the  $TE_{00}$  mode of the rectangular waveguide is non-physical, since it violates the divergence condition.

# Example (cont.)

Solution from separation of variables (given later):



$$k_z = \sqrt{k^2 - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2}$$

$$\text{TM}_{mn}: \quad E_z = E_0 \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right) e^{-jk_z z}$$

$$\text{TE}_{mn}: \quad H_z = H_0 \cos\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right) e^{-jk_z z}$$

# Example (cont.)

**Question:** What is the lowest-order mode (smallest  $m$  and  $n$ ) for a non-trivial field?

$$\text{TM}_{mn}: E_z = E_0 \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right) e^{-jk_z z}$$

$$\text{TE}_{mn}: H_z = H_0 \cos\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right) e^{-jk_z z}$$

$$\text{TM}_{mn}: m = 1, n = 1$$

$$\text{TE}_{mn}: m = 0, n = 0 ?$$

# Example (cont.)

TE<sub>mn</sub> mode:

$$H_z = H_0 \cos\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right) e^{-jk_z z} \quad k_z = \sqrt{k^2 - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2}$$

TE<sub>00</sub> mode:

$$\begin{aligned} H_z &= H_0 e^{-jkz} \\ H_x &= H_y = 0 \end{aligned}$$

The last equation holds since the normal component of the magnetic field is zero on the PEC, and there is no  $(x, y)$  variation inside the waveguide.

$$\underline{H} \cdot \underline{\hat{n}} = 0$$



# Example (cont.)

Check the divergence:

$$\begin{aligned}\nabla \cdot \underline{H} &= \frac{\cancel{\partial H_x}}{\partial x} + \frac{\cancel{\partial H_y}}{\partial y} + \frac{\partial H_z}{\partial z} \\ &= -jk H_0 e^{-jkz}\end{aligned}$$

The TE<sub>00</sub> mode is a non-physical mode.

Hence  $\nabla \cdot \underline{H} \neq 0$  (invalid mode)

The TE<sub>10</sub> mode is actually the mode with the lowest cut-off frequency.

Whenever we solve the vector Helmholtz equation, we should check to make sure that the divergence condition is satisfied!

It turns out that all of the other modes of the rectangular waveguide are physical.