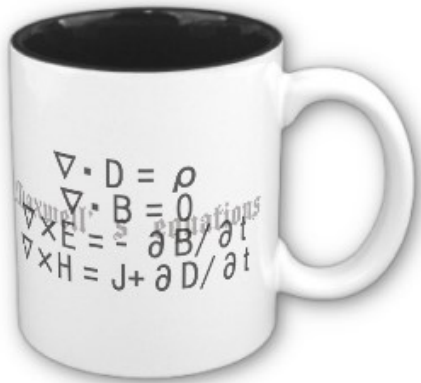


ECE 6340

Intermediate EM Waves

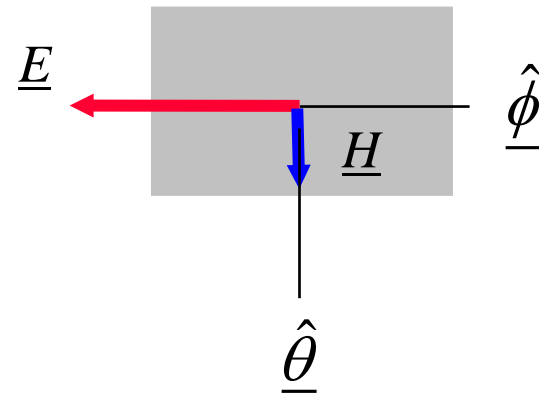
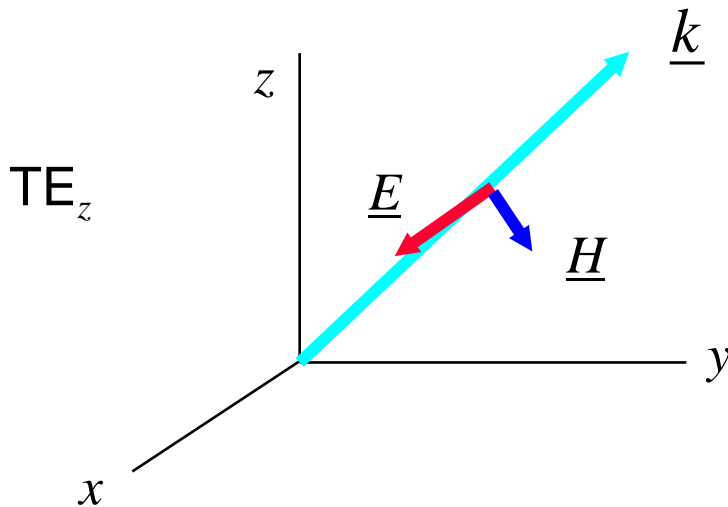
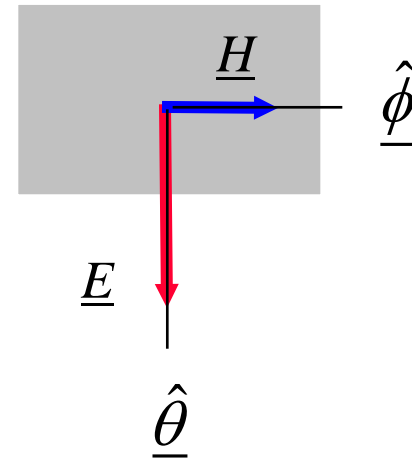
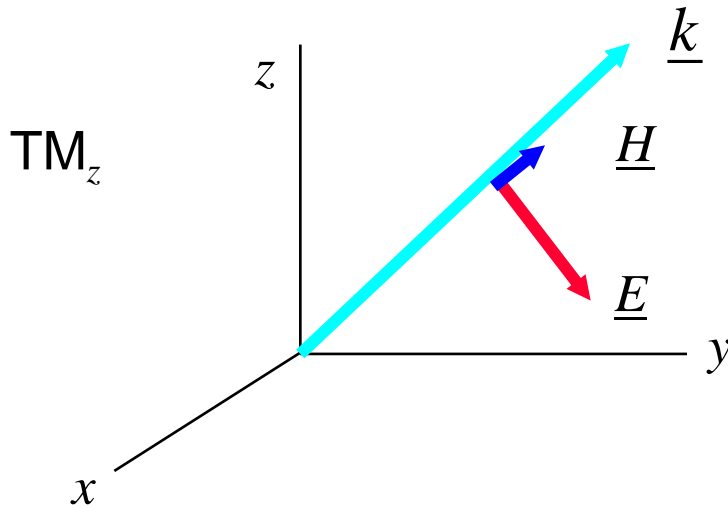
Fall 2016

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Notes 18

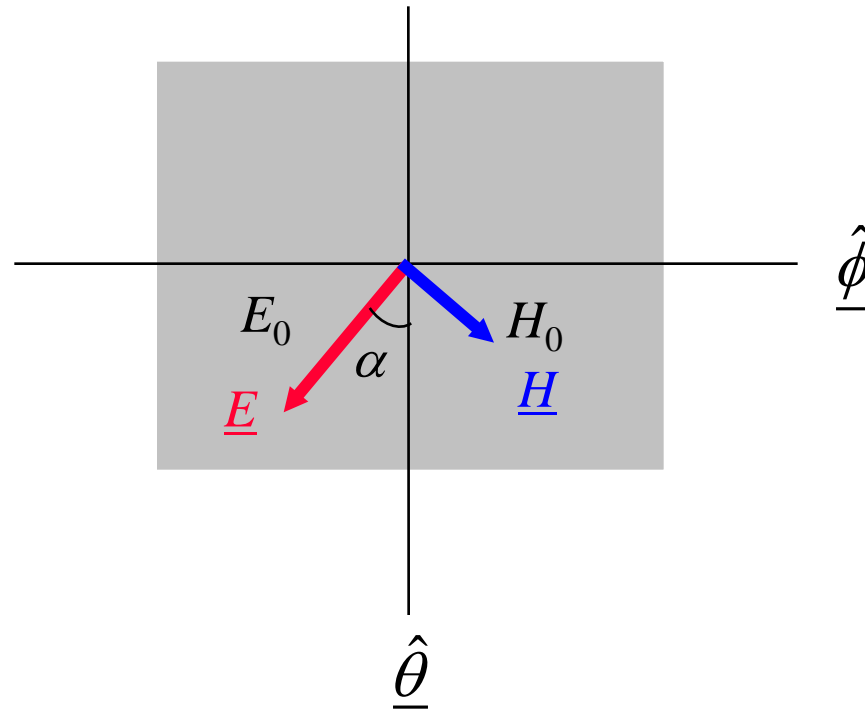
TE_z - TM_z Plane Waves



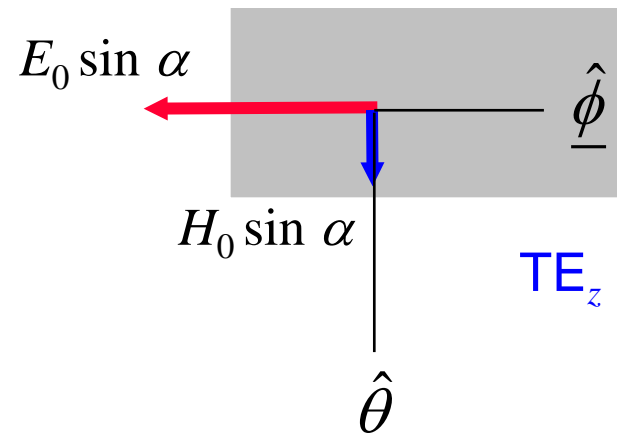
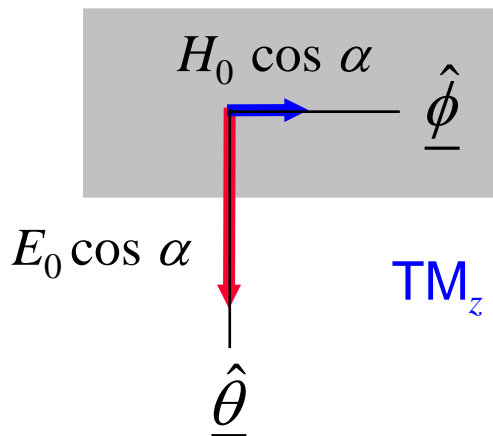
A homogeneous plane wave is shown for simplicity (but the principle is general).

Arbitrary Polarization: Decomposition

Assume that \underline{E} and \underline{H} are real vectors in the figure for simplicity.



TM_z and TE_z parts:



TM_z Plane Wave

$$\nabla \times \underline{H} = j\omega\varepsilon_c \underline{E}$$

or

$$-j\underline{k} \times \underline{H} = j\omega\varepsilon_c \underline{E}$$

or

$$-\underline{k} \times \underline{H} = \omega\varepsilon_c \underline{E}$$

Take x component:

$$\begin{aligned} \omega\varepsilon_c E_x &= \hat{x} \cdot \left[-(\hat{x}k_x + \hat{y}k_y + \hat{z}k_z) \times (\hat{x}H_x + \hat{y}H_y) \right] \\ &= k_z H_y \end{aligned}$$

so

$$\frac{E_x}{H_y} = \frac{k_z}{\omega\varepsilon_c}$$

Assume that the wave is propagating upward (+ z direction) so that k_z is positive (or has a positive real part for a lossy medium).

TM_z Plane Wave (cont.)

Take y component:

$$\begin{aligned}\omega \varepsilon_c E_y &= \underline{\hat{y}} \cdot \left[-(\underline{\hat{x}} k_x + \underline{\hat{y}} k_y + \underline{\hat{z}} k_z) \times (\underline{\hat{x}} H_x + \underline{\hat{y}} H_y) \right] \\ &= -k_z H_x\end{aligned}$$

$$\frac{E_y}{-H_x} = \frac{k_z}{\omega \varepsilon_c}$$

Define:

$$Z^{TM} = \frac{k_z}{\omega \varepsilon_c}$$

TM_z Plane Wave (cont.)

Both results are summarized in a vector equation:

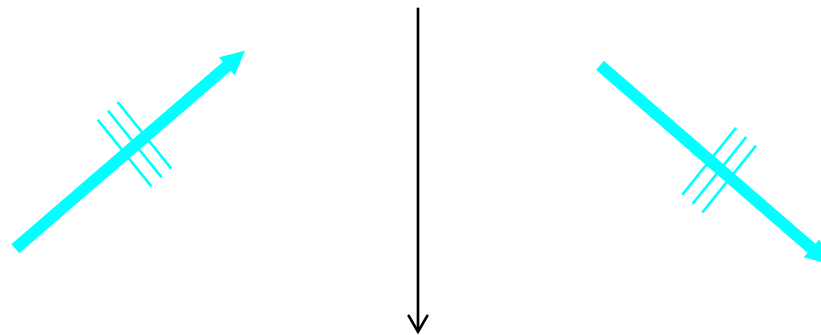
$$\underline{E}_t = -Z^{TM} (\hat{z} \times \underline{H}_t)$$

or

$$\underline{H}_t = \frac{1}{Z^{TM}} (\hat{z} \times \underline{E}_t)$$

Note:
“t” stands for
“transverse,”
meaning the x and y
components.

Consider replacing $k_z \rightarrow -k_z$



$$Z^{TM} \rightarrow -Z^{TM}$$

Recall: $Z^{TM} = \frac{k_z}{\omega \epsilon_c}$

TM_z Plane Wave (cont.)

so

$$\underline{H}_t = -\frac{1}{Z^{TM}} (\hat{z} \times \underline{E}_t)$$

(Here Z^{TM} is the same as before, defined for a $+z$ wave.)

Summary for both cases:

$$\underline{H}_t = \pm \frac{1}{Z^{TM}} (\hat{z} \times \underline{E}_t)$$

$$\underline{E}_t = \mp Z^{TM} (\hat{z} \times \underline{H}_t)$$

TE_z Plane Wave

$$\nabla \times \underline{E} = -j\omega\mu \underline{H}$$

or

$$-j \underline{k} \times \underline{E} = -j\omega\mu \underline{H}$$

or

$$\underline{k} \times \underline{E} = \omega\mu \underline{H}$$

Take x component:

$$\begin{aligned} \omega \mu H_x &= \hat{x} \cdot \left[(\hat{x} k_x + \hat{y} k_y + \hat{z} k_z) \times (\hat{x} E_x + \hat{y} E_y) \right] \\ &= -k_z E_y \end{aligned}$$

so

$$\frac{E_y}{-H_x} = \frac{\omega\mu}{k_z}$$

TE_z Plane Wave (cont.)

Take y component:

$$\begin{aligned}\omega \mu H_y &= \underline{\hat{y}} \cdot \left[(\underline{\hat{x}} k_x + \underline{\hat{y}} k_y + \underline{\hat{z}} k_z) \times (\underline{\hat{x}} E_x + \underline{\hat{y}} E_y) \right] \\ &= k_z E_x\end{aligned}$$

so

$$\frac{E_x}{H_y} = \frac{\omega \mu}{k_z} = \eta \frac{k}{k_z}$$

Define:

$$Z^{TE} = \frac{\omega \mu}{k_z}$$

TE_z Plane Wave (cont.)

Then

$$\underline{E}_t = -Z^{TE} (\hat{z} \times \underline{H}_t)$$

$$\underline{H}_t = \frac{1}{Z^{TE}} (\hat{z} \times \underline{E}_t)$$

Allowing for both directions, +z and -z, we have:

$$\underline{H}_t = \pm \frac{1}{Z^{TE}} (\hat{z} \times \underline{E}_t)$$

$$\underline{E}_t = \mp Z^{TE} (\hat{z} \times \underline{H}_t)$$

Transverse Equivalent Network

Assume a TM_z plane wave going upward
(the z propagation is in the $+z$ direction).

Denote

$$\underline{E}_t(x, y, z) = \underline{\hat{e}}_{TM} V_{TM}(z) \psi_t(x, y)$$

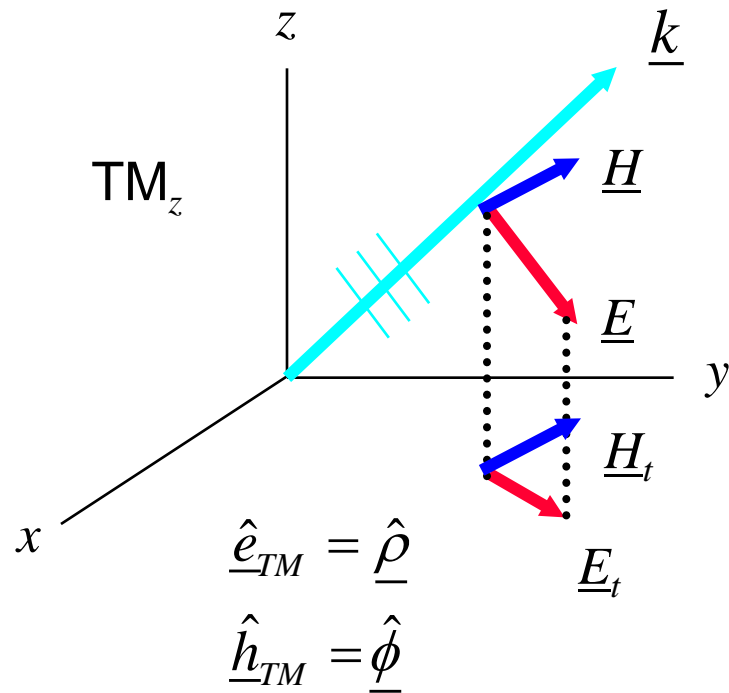
$$\underline{H}_t(x, y, z) = \underline{\hat{h}}_{TM} I_{TM}(z) \psi_t(x, y)$$

where

$$\psi_t(x, y) = e^{-j(k_x x + k_y y)}$$

As we will see,

$V_{TM}(z)$ and $I_{TM}(z)$ behave as voltage and current on a TL



TEN (cont.)

Assume a +z wave:

$$V_{TM}(z) = Ae^{-jk_z z}$$

$$I_{TM}(z) = Be^{-jk_z z}$$

The form is the same as the waves on a TL.

Also
$$\underline{H}_t = \frac{1}{Z^{TM}} (\hat{z} \times \underline{E}_t)$$

Therefore

$$\psi_t(x, y) I_{TM}(z) \hat{h}_{TM} = \frac{1}{Z^{TM}} \psi_t(x, y) V_{TM}(z) (\hat{z} \times \hat{e}_{TM})$$

We choose:

$$\underline{\hat{h}}_{TM} = \hat{z} \times \underline{\hat{e}}_{TM}$$

TEN (cont.)

We then have

$$I_{TM}(z) = \frac{1}{Z_{TM}} V_{TM}(z)$$

If we assume a $-z$ (downward) wave:

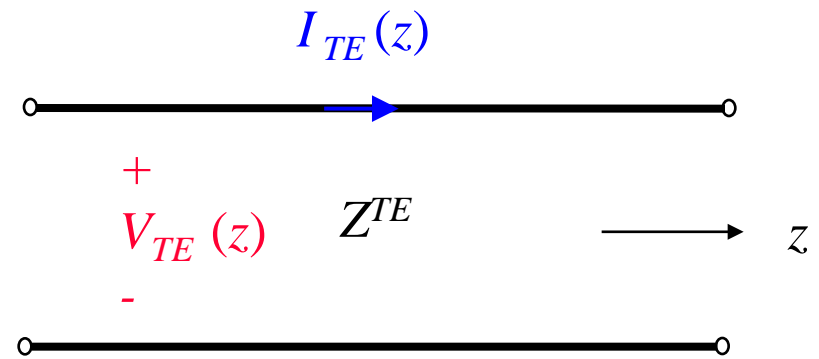
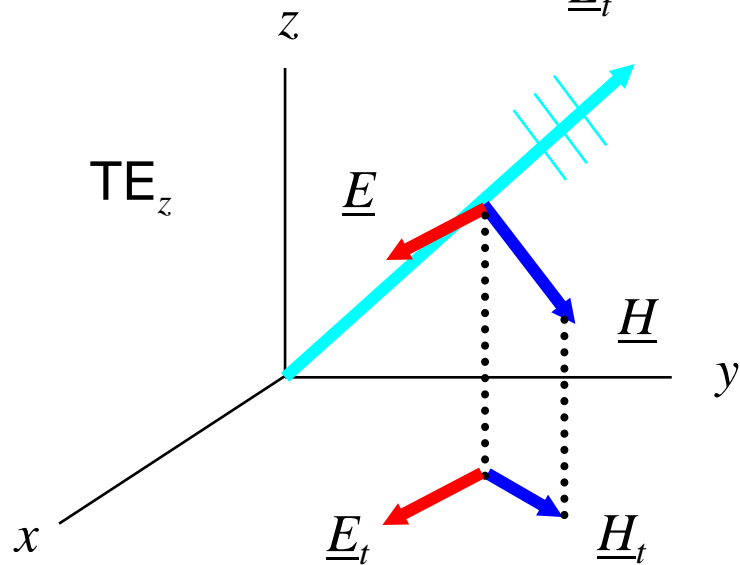
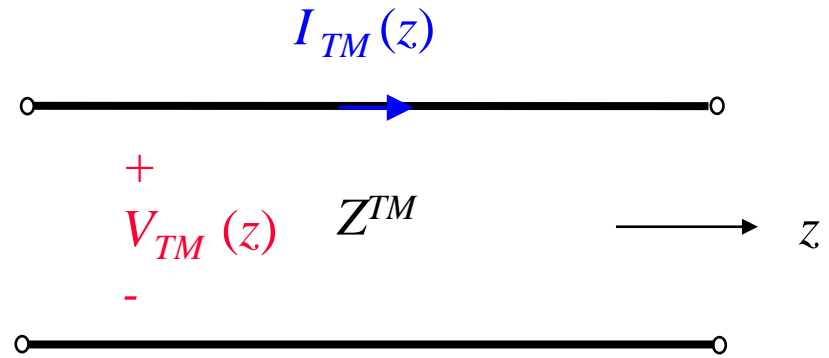
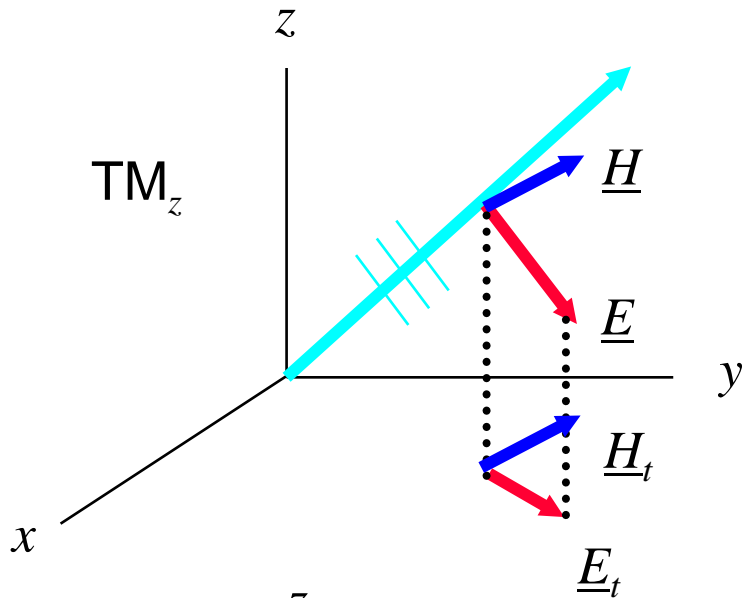
$$I_{TM}(z) = -\frac{1}{Z_{TM}} V_{TM}(z)$$

Hence, in summary:

$$I_{TM}(z) = \pm \frac{1}{Z_{TM}} V_{TM}(z)$$

This proves that the transverse fields behave as voltage and current on a TL.

TEN (cont.)



TEN (cont.)

Note: $V(z)$ and $I(z)$ model only the transverse fields, but we can obtain the z component of the fields from these.

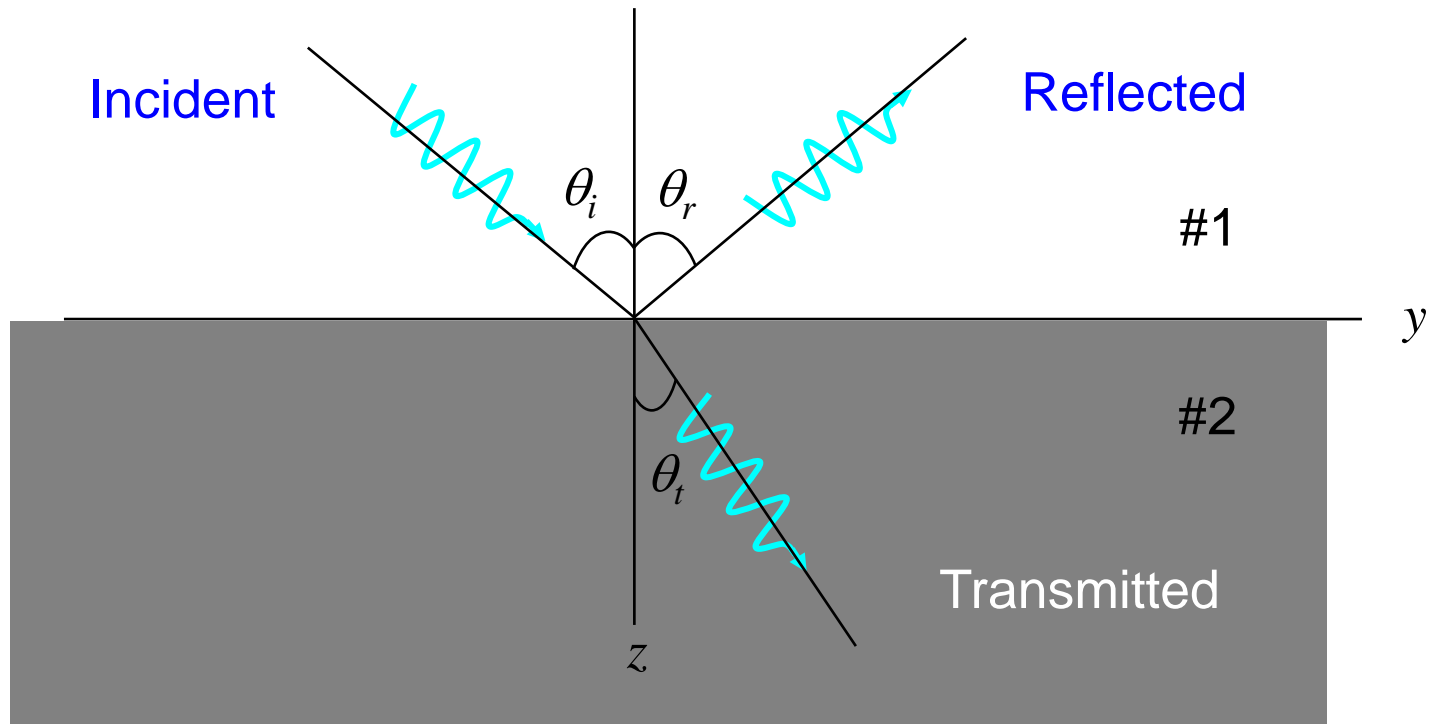
Example: Find E_z for a TM_z plane wave:

$$\nabla \times \underline{H} = j\omega \epsilon_c \underline{E}$$

$$E_z = \frac{1}{j\omega \epsilon_c} \hat{z} \cdot (\nabla \times \underline{H})$$

$$\begin{aligned} E_z &= \frac{1}{j\omega \epsilon_c} \left[\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right] \\ &= \frac{1}{j\omega \epsilon_c} \left[(-jk_x)H_y + (jk_y)H_x \right] \end{aligned}$$

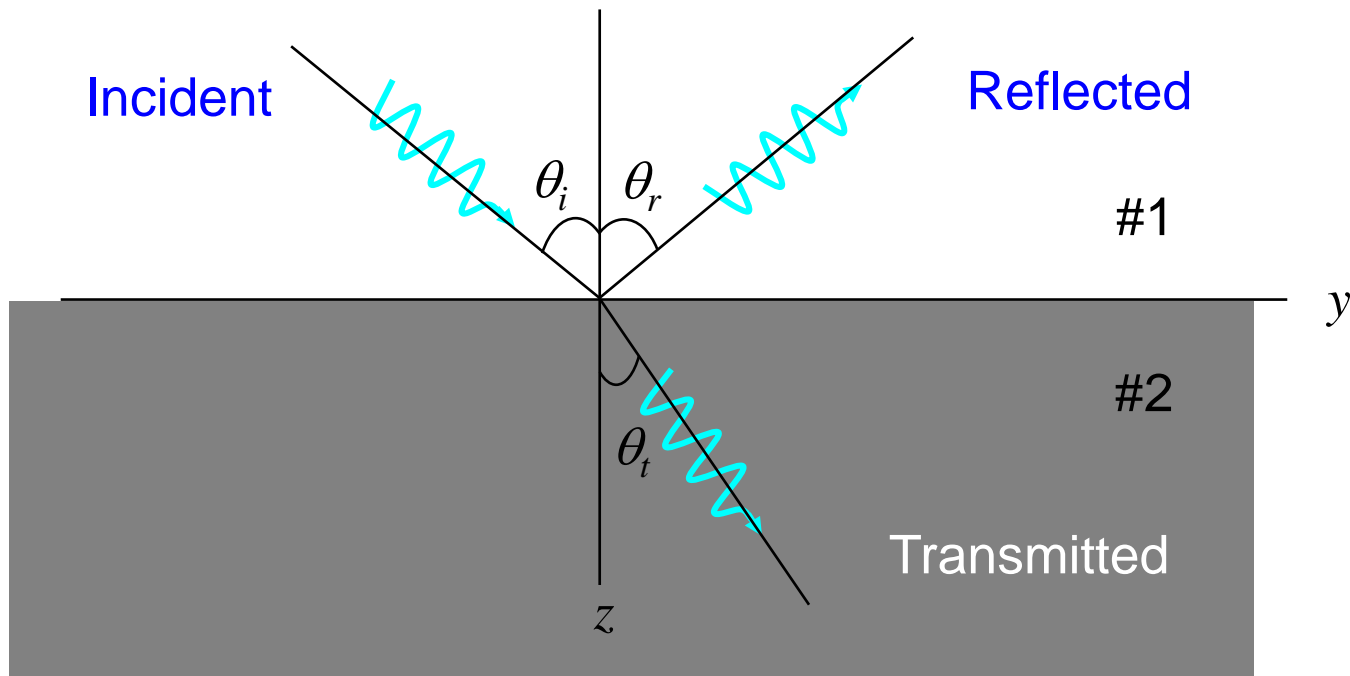
Reflection From Interface



We want to find the directions of the reflected and transmitted waves, and the reflection and transmission coefficients.

Note: The yz plane is the “plane of incidence” (the plane containing the incident wave vector and the unit normal to the boundary).

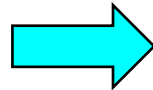
Reflection From Interface



$$k_x = k \sin \theta \cos \phi$$

$$k_y = k \sin \theta \sin \phi$$

$$k_z = k \cos \theta$$



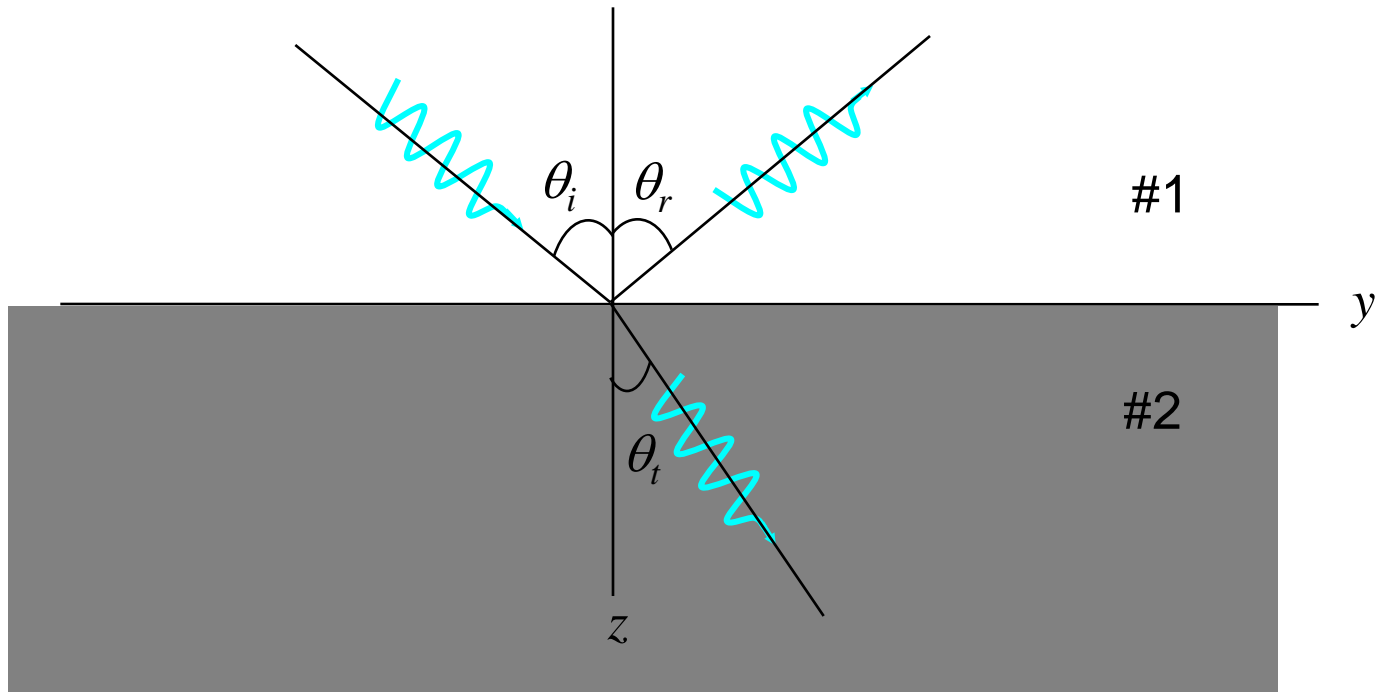
$$\phi = \pi / 2$$

$$k_x = 0$$

$$k_y = k \sin \theta$$

$$k_z = k \cos \theta$$

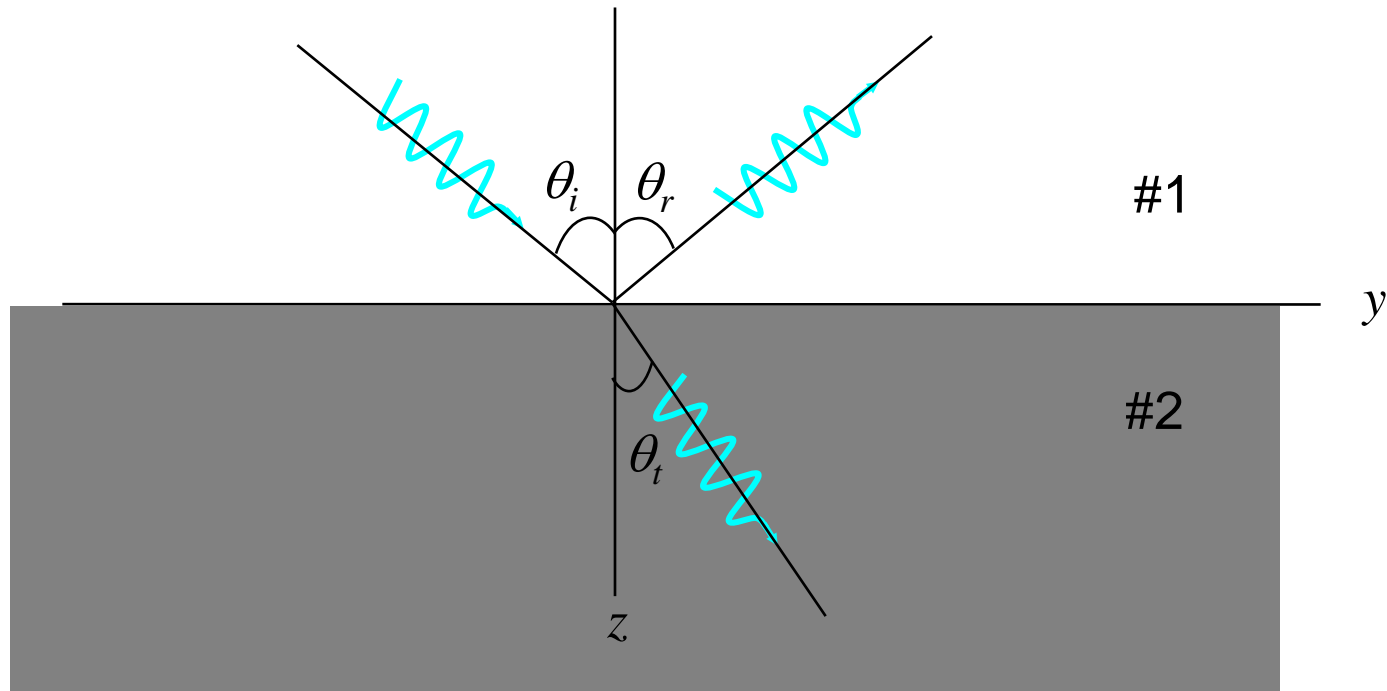
Reflection From Interface (cont.)



Phase matching condition: $\underbrace{k_{yi} = k_{yr}} = k_{yt} \Rightarrow k_1 \sin \theta_i = k_1 \sin \theta_r$

Hence $\theta_i = \theta_r$ (law of reflection)

Reflection From Interface (cont.)



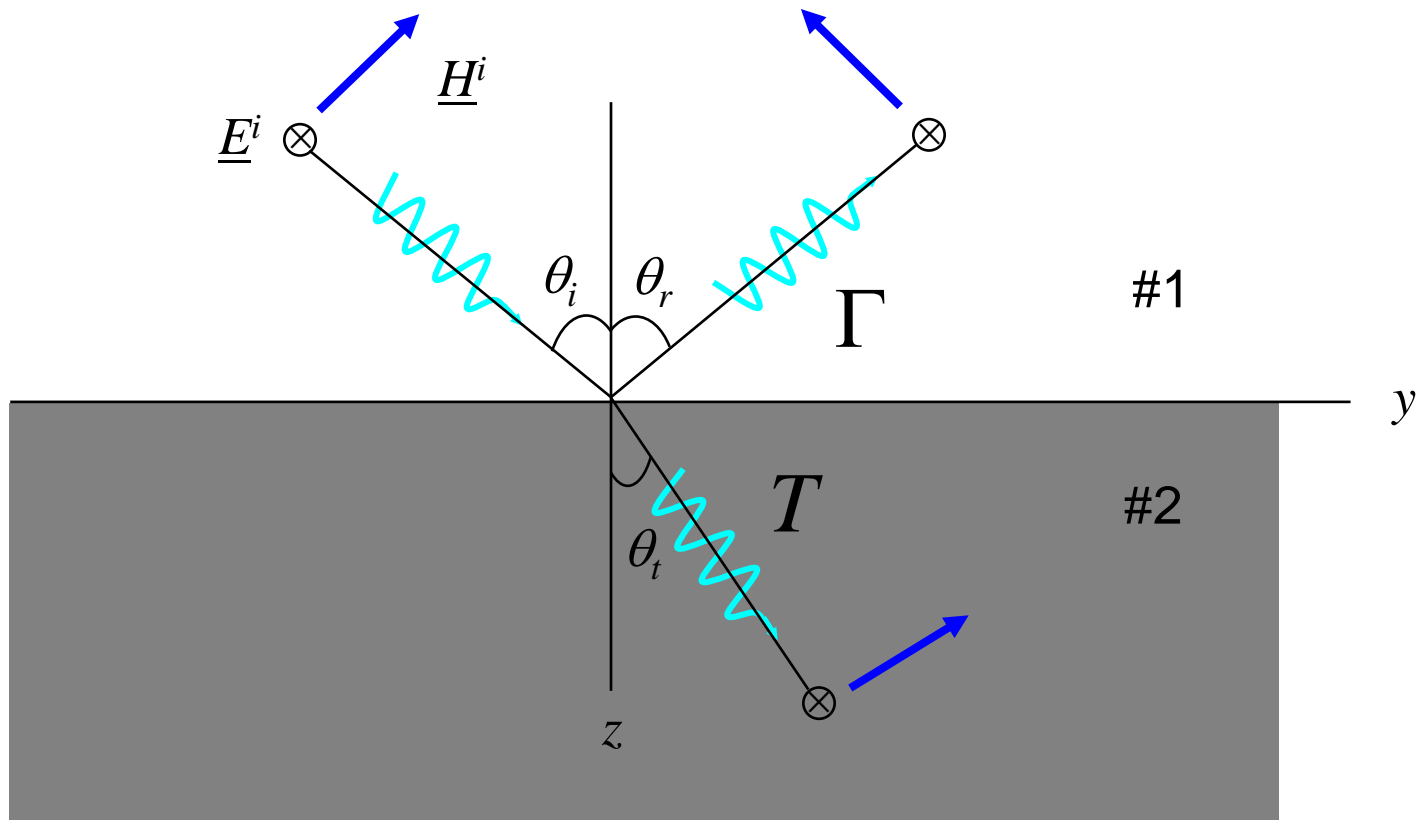
Phase matching condition: $k_{yi} = k_{yr} = k_{yt}$

Hence $k_1 \sin \theta_i = k_2 \sin \theta_t$ (Snell's law)

Reflection From Interface (cont.)

Determine reflection and transmission coefficients.

(Assume a TE_z wave)



Reflection From Interface (cont.)

Modeling equations:

$$E_x \leftrightarrow V$$

$$H_y \leftrightarrow I$$

$$E_x^i = V^i(z) e^{-jk_y y}$$

$$E_x^r = V^r(z) e^{-jk_y y}$$

$$E_x^t = V^t(z) e^{-jk_y y}$$

$$V^i(z) = e^{-jk_{z1} z}$$

$$V^r(z) = \Gamma e^{+jk_{z1} z}$$

$$V^t(z) = T e^{-jk_{z2} z}$$

$$k_y = k_1 \sin \theta_i$$

$$k_{z1} = \sqrt{k_1^2 - k_y^2} = k_1 \cos \theta_i$$

$$k_{z2} = \sqrt{k_2^2 - k_y^2} = k_2 \cos \theta_t$$

Reflection From Interface (cont.)

Practical note:

When dealing with lossy media, the wave in region 2 will be inhomogeneous. Therefore the transmitted angle will be complex. In this case it is usually easier to work with the separation equation (the square-root formula for k_{z2}) rather than the transmitted-angle formula.

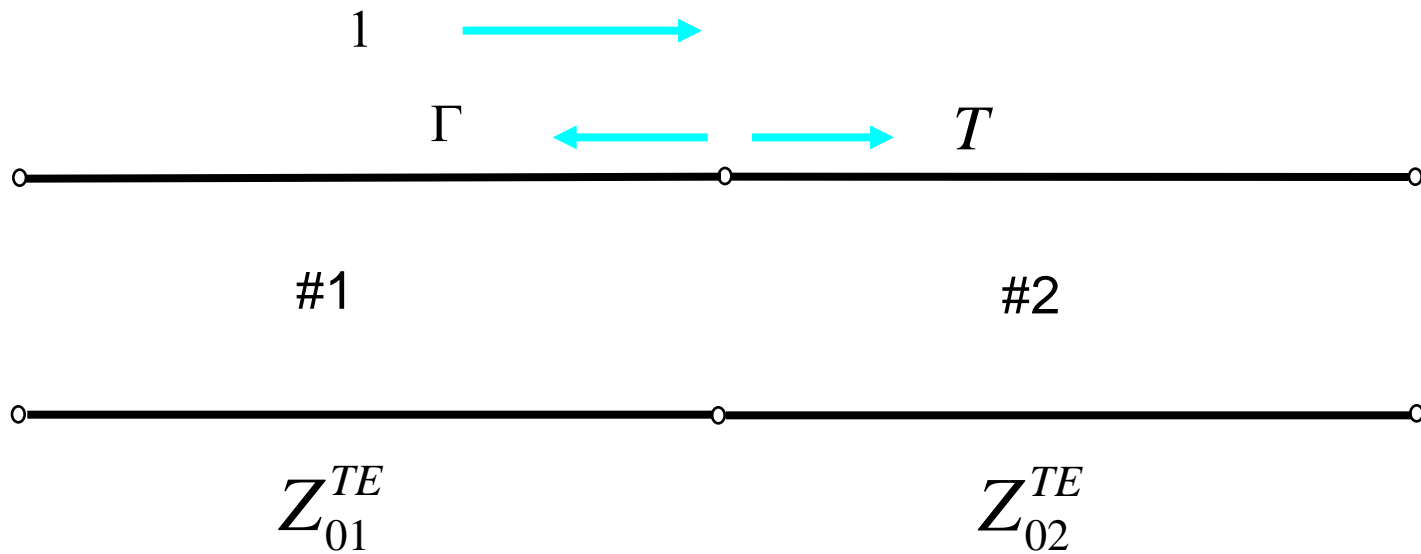
$$k_{y2} = k_2 \sin \theta_t = k_{y1} = k_1 \sin \theta_i = \text{real}$$

complex It is difficult to solve for θ_t

$$k_{z2} = \sqrt{k_2^2 - k_{y2}^2} = k_2 \cos \theta_t$$

Easy Hard (needs complex angle)

Reflection From Interface (cont.)



$$\Gamma = \frac{Z_{02}^{TE} - Z_{01}^{TE}}{Z_{02}^{TE} + Z_{01}^{TE}}$$

$$T = 1 + \Gamma = \frac{2Z_{02}^{TE}}{Z_{02}^{TE} + Z_{01}^{TE}}$$

Reflection From Interface (cont.)

Hence

$$\Gamma = \frac{\frac{\omega\mu_2}{k_{z2}} - \frac{\omega\mu_1}{k_{z1}}}{\frac{\omega\mu_2}{k_{z2}} + \frac{\omega\mu_1}{k_{z1}}}$$

or

$$\Gamma = \frac{\mu_2 k_{z1} - \mu_1 k_{z2}}{\mu_2 k_{z1} + \mu_1 k_{z2}} \quad T = \frac{2\mu_2 k_{z1}}{\mu_2 k_{z1} + \mu_1 k_{z2}}$$

Reflection From Interface (cont.)

Percent power reflected:

$$P_r^{\%} = 100 |\Gamma|^2$$

Percent power transmitted:

$$P_t^{\%} = 100 - P_r^{\%}$$