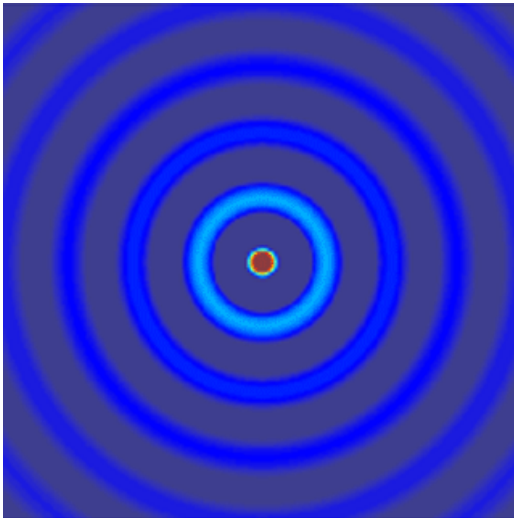


ECE 6341

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Notes 44

Overview

In this set of notes we derive the SDI formulation using a more mathematical, but more general, approach (we directly Fourier transform Maxwell's equations).

- ❖ This allows for **all possible types of sources** (horizontal, vertical, electric, and magnetic) to be treated in one derivation.

General SDI Method

Start with Ampere's law:

$$\nabla \times \underline{H} = \underline{J}^i = j\omega\varepsilon \underline{E}$$

$$\nabla = \nabla_t + \hat{\underline{z}} \frac{\partial}{\partial z}$$

where

$$\nabla_t = \hat{\underline{x}} \frac{\partial}{\partial x} + \hat{\underline{y}} \frac{\partial}{\partial y}$$

Assume a 2D spatial transform:

$$\begin{aligned} \tilde{\nabla}_t &= \hat{\underline{x}}(-jk_x) + \hat{\underline{y}}(-jk_y) \\ &= -j(\hat{\underline{x}}k_x + \hat{\underline{y}}k_y) \\ &= -j\underline{k}_t \\ &= -jk_t \hat{\underline{u}} \end{aligned}$$

General SDI Method (cont.)

Hence we have
$$\left(-jk_t \underline{\hat{u}} + \underline{\hat{z}} \frac{\partial}{\partial z} \right) \times \tilde{\underline{H}} = \tilde{\underline{J}}^i + j\omega\varepsilon \tilde{\underline{E}}$$

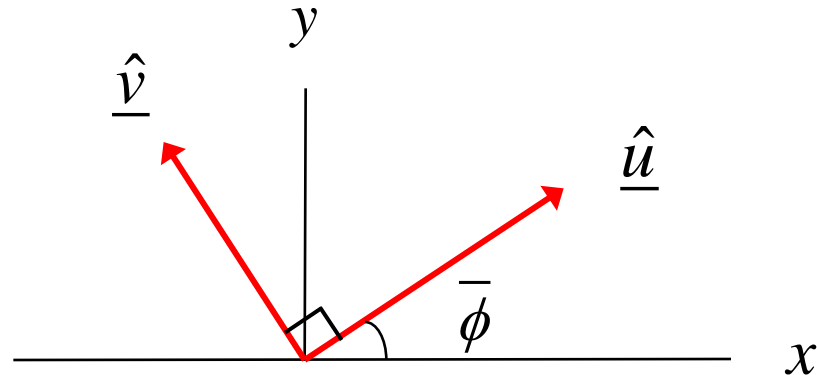
Next, represent the field as
$$\begin{aligned} \tilde{\underline{H}}_t &= \underline{\hat{u}} \tilde{H}_u + \underline{\hat{v}} \tilde{H}_v \\ &= \underline{\hat{u}} (\tilde{\underline{H}} \cdot \underline{\hat{u}}) + \underline{\hat{v}} (\tilde{\underline{H}} \cdot \underline{\hat{v}}) \end{aligned}$$

Note that

$$\underline{\hat{u}} \times \underline{\hat{v}} = \underline{\hat{z}}$$

$$\underline{\hat{z}} \times \underline{\hat{u}} = \underline{\hat{v}}$$

$$\underline{\hat{z}} \times \underline{\hat{v}} = -\underline{\hat{u}}$$



Take the $\underline{\hat{z}}$, $\underline{\hat{u}}$, $\underline{\hat{v}}$ components of the transformed Ampere's equation:

General SDI Method (cont.)

$$\underline{\hat{z}}) -jk_t \tilde{H}_v = \tilde{J}_z^i + j\omega\varepsilon \tilde{E}_z$$

$$\underline{\hat{u}}) -\frac{\partial \tilde{H}_v}{\partial z} = \tilde{J}_u^i + j\omega\varepsilon \tilde{E}_u$$

$$\underline{\hat{v}}) jk_t \tilde{H}_z + \frac{\partial \tilde{H}_u}{\partial z} = \tilde{J}_v^i + j\omega\varepsilon \tilde{E}_v$$

Examine the **TM_z** field: $(\tilde{E}_u, \tilde{H}_v, \tilde{E}_z)$ (Ignore $\underline{\hat{v}}$ equation)

$$-jk_t \tilde{H}_v = \tilde{J}_z^i + j\omega\varepsilon \tilde{E}_z \quad (1)$$

$$-\frac{\partial \tilde{H}_v}{\partial z} = \tilde{J}_u^i + j\omega\varepsilon \tilde{E}_u \quad (2)$$

TM_z Fields

We wish to eliminate \tilde{E}_z from Eq. (1). To do this, use Faraday's law:

$$\nabla \times \underline{E} = -\underline{M}^i - j\omega\mu\underline{H}$$

→ $\left(-jk_t \underline{\hat{u}} + \underline{\hat{z}} \frac{\partial}{\partial z}\right) \times \underline{\tilde{E}} = -\underline{\tilde{M}}^i - j\omega\mu\underline{\tilde{H}}$

Take the $\underline{\hat{v}}$ component of the transformed Faraday's Law:

Recall:

$$\underline{\hat{u}} \times \underline{\hat{v}} = \underline{\hat{z}}$$

$$\underline{\hat{z}} \times \underline{\hat{u}} = \underline{\hat{v}}$$

$$\underline{\hat{z}} \times \underline{\hat{v}} = -\underline{\hat{u}}$$

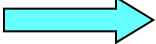
$$jk_t \tilde{E}_z + \frac{\partial \tilde{E}_u}{\partial z} = -\tilde{M}_v^i - j\omega\mu \tilde{H}_v \quad (3)$$

TM_z Fields (cont.)

Substitute \tilde{E}_z from (1) into (3) to eliminate \tilde{E}_z

$$-jk_t \tilde{H}_v = \tilde{J}_z^i + j\omega\varepsilon \tilde{E}_z \quad (1)$$

$$jk_t \tilde{E}_z + \frac{\partial \tilde{E}_u}{\partial z} = -\tilde{M}_v^i - j\omega\mu \tilde{H}_v \quad (3)$$


$$jk_t \left[\frac{1}{j\omega\varepsilon} (-\tilde{J}_z^i - jk_t \tilde{H}_v) \right] + \frac{\partial \tilde{E}_u}{\partial z} = -\tilde{M}_v^i - j\omega\mu \tilde{H}_v$$

or

$$\frac{\partial \tilde{E}_u}{\partial z} + \frac{k_t^2}{j\omega\varepsilon} \tilde{H}_v + j\omega\mu \tilde{H}_v = -\tilde{M}_v^i + \frac{k_t}{\omega\varepsilon} \tilde{J}_z^i$$

TM_z Fields (cont.)

Note that

$$\begin{aligned}\frac{k_t^2}{j\omega\epsilon} + j\omega\mu &= \frac{1}{j\omega\epsilon} (k_t^2 - \omega^2\mu\epsilon) \\ &= \frac{1}{j\omega\epsilon} (k_t^2 - k^2) \\ &= \frac{-1}{j\omega\epsilon} (k^2 - k_t^2) \\ &= \frac{-k_z^2}{j\omega\epsilon}\end{aligned}$$

Hence

$$\frac{\partial \tilde{E}_u}{\partial z} - \left(\frac{k_z^2}{j\omega\epsilon} \right) \tilde{H}_v = -\tilde{M}_v^i + \left(\frac{k_t}{\omega\epsilon} \right) \tilde{J}_z^i \quad (4)$$

TM_z Fields (cont.)

Equations (2) and (4) are the final TM_z field modeling equations:

$$\frac{\partial \tilde{H}_v}{\partial z} = -\tilde{J}_u^i - j\omega\epsilon\tilde{E}_u$$
$$\frac{\partial \tilde{E}_u}{\partial z} = -\tilde{M}_v^i + \left(\frac{k_t}{\omega\epsilon}\right)\tilde{J}_z^i + \left(\frac{k_z^2}{j\omega\epsilon}\right)\tilde{H}_v$$

TM_z Fields (cont.)

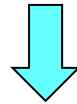
Define TEN modeling equations:

$$V^{TM}(z) \equiv \tilde{E}_u(k_x, k_y, z)$$

$$I^{TM}(z) \equiv \tilde{H}_v(k_x, k_y, z)$$

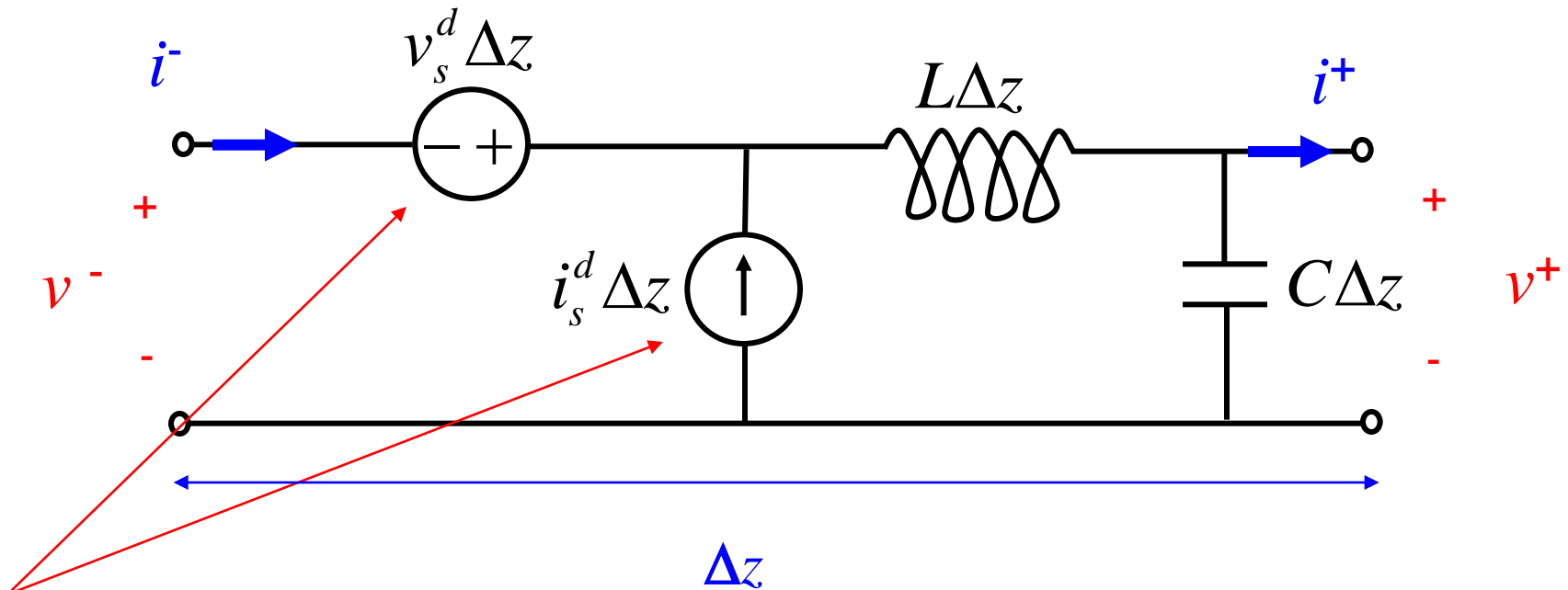
We then have

$$\begin{aligned} \frac{\partial \tilde{H}_v}{\partial z} &= -\tilde{J}_u^i - j\omega\epsilon\tilde{E}_u \\ \frac{\partial \tilde{E}_u}{\partial z} &= -\tilde{M}_v^i + \left(\frac{k_t}{\omega\epsilon}\right)\tilde{J}_z^i + \left(\frac{k_z^2}{j\omega\epsilon}\right)\tilde{H}_v \end{aligned}$$



$$\begin{aligned} \frac{\partial I^{TM}}{\partial z} &= -j\omega\epsilon V^{TM} - \tilde{J}_u^i \\ \frac{\partial V^{TM}}{\partial z} &= \left(\frac{k_z^2}{j\omega\epsilon}\right)I^{TM} + \left[-\tilde{M}_v^i + \left(\frac{k_t}{\omega\epsilon}\right)\tilde{J}_z^i\right] \end{aligned}$$

Telegrapher's Equation



Allow for distributed sources

$$v^+ - v^- = v_s^d \Delta z - (L\Delta z) \frac{\partial}{\partial t} (i^- + i_s^d \Delta z)$$

$$\Delta z \rightarrow 0 \quad \text{so} \quad \frac{\partial v}{\partial z} = v_s^d - L \frac{\partial i}{\partial t}$$

Telegrapher's Equation (cont.)

Hence, in the phasor domain,

$$\frac{\partial V}{\partial z} = -j\omega LI + V_s^d$$

Also, $i^+ - i^- = i_s^d \Delta z - (C\Delta z) \frac{\partial v^+}{\partial t}$

$\Delta z \rightarrow 0$ so $\frac{\partial i}{\partial z} = i_s^d - C \frac{\partial v}{\partial t}$

Hence, we have

$$\frac{\partial I}{\partial z} = -j\omega CV + I_s^d$$

Telegrapher's Equation (cont.)

Compare field equations for TM_z fields with TL equations:

$$\frac{\partial I^{TM}}{\partial z} = -j\omega\epsilon V^{TM} - \tilde{J}_u^i$$

$$\frac{\partial V^{TM}}{\partial z} = \left(\frac{k_z^2}{j\omega\epsilon} \right) I^{TM} + \left[-\tilde{M}_v^i + \left(\frac{k_t}{\omega\epsilon} \right) \tilde{J}_z^i \right]$$

$$\frac{\partial I}{\partial z} = -j\omega CV + I_s^d$$

$$\frac{\partial V}{\partial z} = -j\omega LI + V_s^d$$

Telegrapher's Equation (cont.)

We then make the following identifications:

$$C = \varepsilon$$
$$\omega L = \frac{k_z^2}{\omega \varepsilon}$$

Hence

$$k_z^{TL} = \omega \sqrt{LC} = \omega \sqrt{\left(\frac{k_z^2}{\omega^2 \varepsilon}\right) \varepsilon} = k_z$$

$$Z_0^{TL} = \sqrt{\frac{L}{C}} = \sqrt{\left(\frac{k_z^2}{\omega^2 \varepsilon}\right) \frac{1}{\varepsilon}} = \frac{k_z}{\omega \varepsilon}$$

so

$$k_z^{TL} = k_z$$
$$Z_0^{TL} = \frac{k_z}{\omega \varepsilon}$$

Sources: TM_z

For the sources we have, for the TM_z case:

$$I_s^{dTM} = -\tilde{J}_u^i$$
$$V_s^{dTM} = -\tilde{M}_v^i + \left(\frac{k_t}{\omega \epsilon} \right) \tilde{J}_z^i$$

Sources: TM_z (cont.)

Special case: **planar** horizontal surface current sources:

Assume

$$\underline{J}^i(x, y, z) = \underline{J}_s^i(x, y) \delta(z)$$
$$\underline{M}^i(x, y, z) = \underline{M}_s^i(x, y) \delta(z)$$

Then we have

$$I_s^{dTM} = -\tilde{J}_u^i = -\tilde{J}_{su}^i \delta(z)$$
$$V_s^{dTM} = -\tilde{M}_v^i = -\tilde{M}_{sv}^i \delta(z)$$

These correspond to **lumped** current and voltage sources:

$$I_s^{TM} = -\tilde{J}_{su}^i$$

lumped parallel current source

$$V_s^{TM} = -\tilde{M}_{sv}^i$$

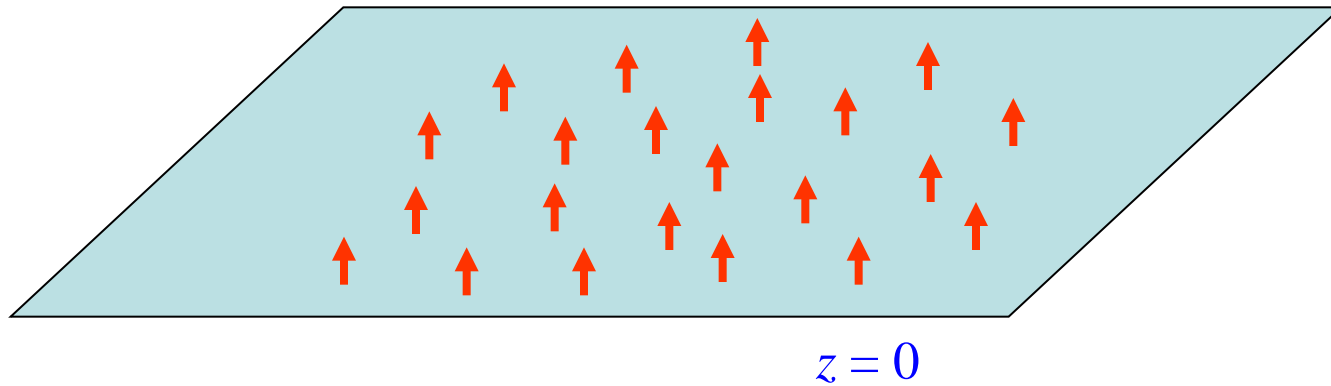
lumped series voltage source

Sources: TM_z (cont.)

For a vertical electric current:

Assume $J_z^i(x, y, z) = J_{sz}^i(x, y) \delta(z)$

“planar vertical current distribution”



Sources: TM_z (cont.)

$$J_z^i(x, y, z) = J_{sz}^i(x, y) \delta(z)$$

$$V_s^{dTM} = \left(\frac{k_t}{\omega \epsilon} \right) \tilde{J}_z^i(k_x, k_y) = \left(\frac{k_t}{\omega \epsilon} \right) \tilde{J}_{sz}^i(k_x, k_y) \delta(z)$$

This corresponds to a *lumped series voltage source*:

$$V_s^{TM} = \left(\frac{k_t}{\omega \epsilon} \right) \tilde{J}_{sz}^i(k_x, k_y)$$

Sources: TM_z (cont.)

Special case: vertical electric dipole

$$J_z^i(x, y, z) = \delta(x) \delta(y) \delta(z)$$

$$J_{sz}^i(x, y) = \delta(x) \delta(y)$$

$$\Rightarrow \tilde{J}_{sz}^i(k_x, k_y) = 1$$

Hence

$$V_s^{TM} = \left(\frac{k_t}{\omega \epsilon} \right)$$

TE_z Fields

Use duality:

$$\underline{E} \rightarrow \underline{H}$$

$$\underline{H} \rightarrow -\underline{E}$$

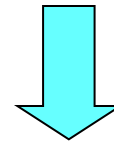
$$\underline{J}^i \rightarrow \underline{M}^i$$

$$\underline{M}^i \rightarrow -\underline{J}^i$$

$$\varepsilon \rightleftharpoons \mu$$

$$\frac{\partial \tilde{H}_v}{\partial z} = -\tilde{J}_u^i - j\omega\varepsilon\tilde{E}_u$$

$$\frac{\partial \tilde{E}_u}{\partial z} = -\tilde{M}_v^i + \left(\frac{k_t}{\omega\varepsilon}\right)\tilde{J}_z^i + \left(\frac{k_z^2}{j\omega\varepsilon}\right)\tilde{H}_v$$



$$-\frac{\partial \tilde{E}_v}{\partial z} = -\tilde{M}_u^i - j\omega\mu\tilde{H}_u$$

$$\frac{\partial \tilde{H}_u}{\partial z} = +\tilde{J}_v^i + \left(\frac{k_t}{\omega\mu}\right)\tilde{M}_z^i - \left(\frac{k_z^2}{j\omega\mu}\right)\tilde{E}_v$$

TM_z

TE_z

TE_z (cont.)

Define:

$$V^{TE}(z) \equiv -\tilde{E}_v(k_x, k_y, z)$$

$$I^{TE}(z) \equiv \tilde{H}_u(k_x, k_y, z)$$

$$\frac{\partial V^{TE}}{\partial z} = -j\omega\mu I^{TE} - \tilde{M}_u^i$$

$$\frac{\partial I^{TE}}{\partial z} = \left(\frac{k_z^2}{j\omega\mu} \right) V^{TE} + \tilde{J}_v^i + \left(\frac{k_t}{\omega\mu} \right) \tilde{M}_z^i$$

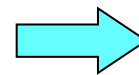
$$\frac{\partial V}{\partial z} = -j\omega LI + V_s^d$$

$$\frac{\partial I}{\partial z} = -j\omega CV + I_s^d$$

We then identify

$$L = \mu$$

$$\omega C = \frac{k_z^2}{\omega\mu}$$



$$k_z^{TL} = k_z$$
$$Z_0^{TE} = \frac{\omega\mu}{k_z}$$

TE_z (cont.)

For the sources, we have

$$V_s^{dTE} = -\tilde{M}_u^i$$
$$I_s^{dTE} = \tilde{J}_v^i + \left(\frac{k_t}{\omega\mu} \right) \tilde{M}_z^i$$

Special case of planar horizontal surface currents:

$$V_s^{TE} = -\tilde{M}_{su}^i$$

lumped series voltage source

$$I_s^{TE} = +\tilde{J}_{sv}^i$$

lumped parallel current source

Sources: TE_z (cont.)

For a vertical magnetic current:

Assume $M_z^i(x, y, z) = M_{sz}^i(x, y) \delta(z)$

This corresponds to a *lumped* parallel current source:

Then we have
$$I_s^{TM} = \left(\frac{k_t}{\omega u} \right) \tilde{M}_{sz}^i(k_x, k_y)$$

Sources: TE_z (cont.)

Special case: vertical magnetic dipole

$$M_z^i(x, y, z) = \delta(x)\delta(y)\delta(z)$$

$$M_{sz}^i(x, y) = \delta(x)\delta(y)$$

$$\Rightarrow \tilde{M}_{sz}^i(k_x, k_y) = 1$$

Hence

$$I_s^{TM} = \left(\frac{k_t}{\omega u} \right)$$

Summary

Results for 3D (volumetric) sources

$$\begin{aligned}V^{TM} &= \tilde{E}_u \\I^{TM} &= \tilde{H}_v \\V^{TE} &= -\tilde{E}_v \\I^{TE} &= \tilde{H}_u\end{aligned}$$

Horizontal

$$\begin{aligned}I_s^{dTM} &= -\tilde{J}_u^i \\V_s^{dTM} &= -\tilde{M}_v^i\end{aligned}$$

$$\begin{aligned}I_s^{dTE} &= +\tilde{J}_v^i \\V_s^{dTE} &= -\tilde{M}_u^i\end{aligned}$$

Vertical

$$V_s^{dTM} = \left(\frac{k_t}{\omega \epsilon} \right) \tilde{J}_z^i$$

$$I_s^{dTE} = \left(\frac{k_t}{\omega \mu} \right) \tilde{M}_z^i$$

*Distributed sources:
either parallel current sources
or series voltage sources*

Summary

Results for 2D (planar) sources

$$\begin{aligned}V^{TM} &= \tilde{E}_u \\ I^{TM} &= \tilde{H}_v \\ V^{TE} &= -\tilde{E}_v \\ I^{TE} &= \tilde{H}_u\end{aligned}$$

Horizontal

$$\begin{aligned}I_s^{TM} &= -\tilde{J}_{su}^i \\ V_s^{TM} &= -\tilde{M}_{sv}^i\end{aligned}$$

$$\begin{aligned}I_s^{TE} &= +\tilde{J}_{sv}^i \\ V_s^{TE} &= -\tilde{M}_{su}^i\end{aligned}$$

Vertical

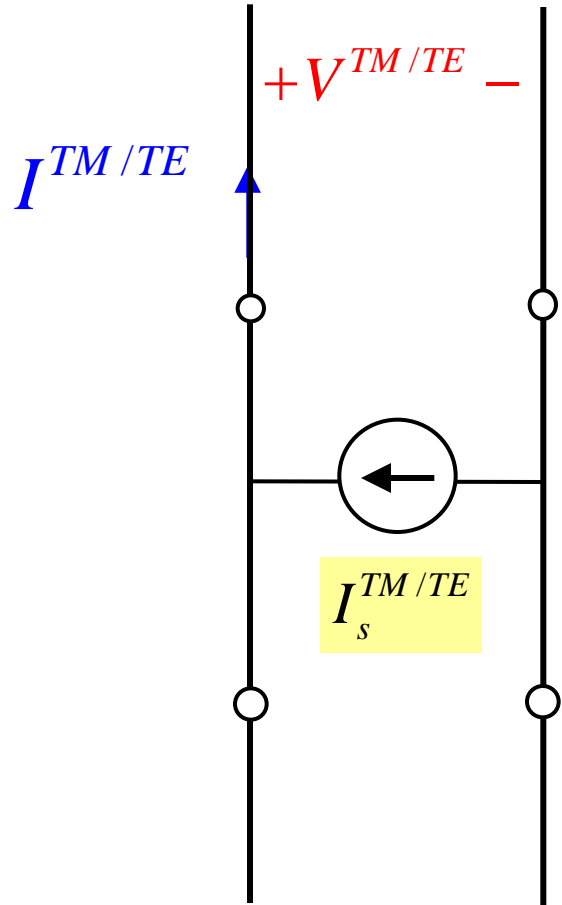
$$V_s^{TM} = \left(\frac{k_t}{\omega \epsilon} \right) \tilde{J}_{sz}^i$$

$$I_s^{TE} = \left(\frac{k_t}{\omega \mu} \right) \tilde{M}_{sz}^i$$

*Lumped sources:
either parallel current sources
or series voltage sources*

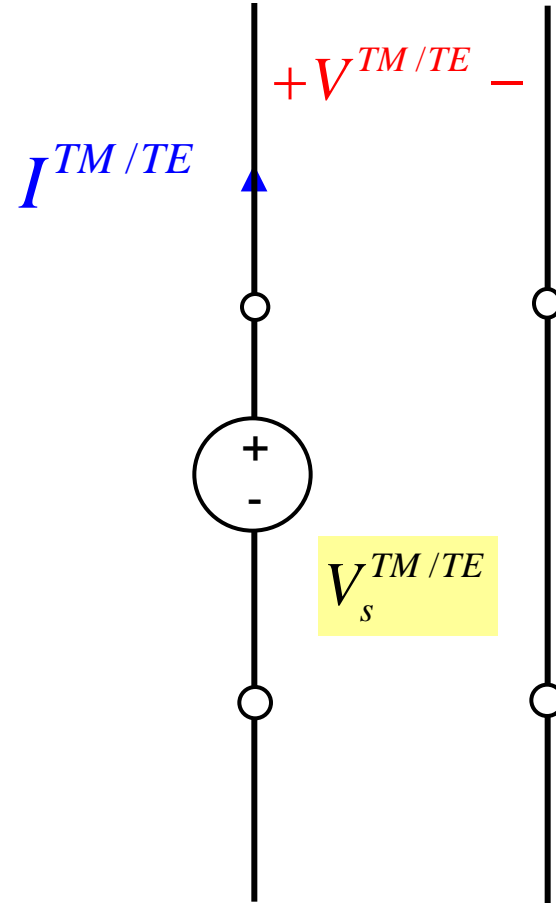
Summary (cont.)

TEN Models



$$V^{TM/TE}(z) = V_i^{TM/TE}(z) I_s^{TM/TE}$$

$$I^{TM/TE}(z) = I_i^{TM/TE}(z) I_s^{TM/TE}$$



$$V^{TM/TE}(z) = V_v^{TM/TE}(z) V_s^{TM/TE}$$

$$I^{TM/TE}(z) = I_v^{TM/TE}(z) V_s^{TM/TE}$$

Summary (cont.)

Michalski functions

