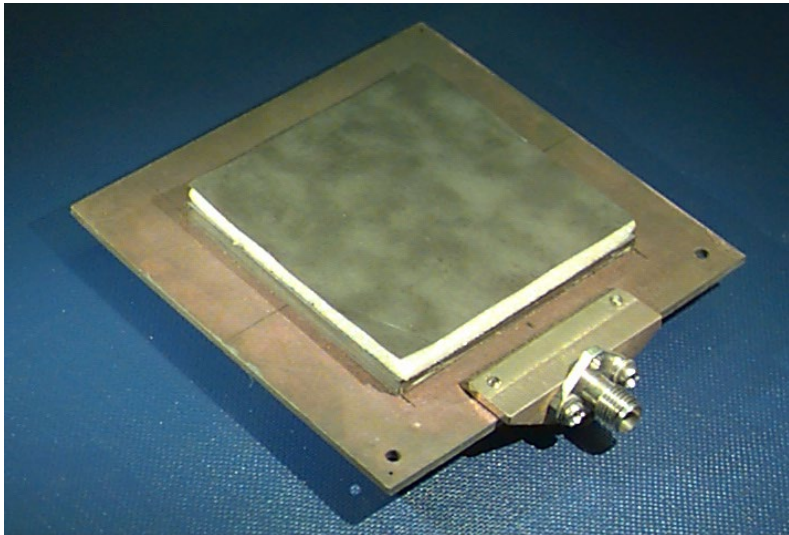


ECE 6345

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ECE Dept.



Notes 10

Overview

In this set of notes we calculate the far field of a rectangular patch using the **magnetic current model**.

The analysis assumes an infinite substrate, but for a truncated substrate we can use the same final result, setting the substrate permittivity to that of air (please see the discussion in Notes 7).

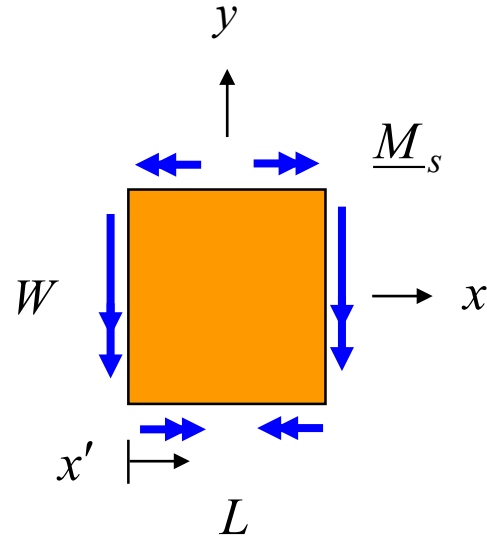
Magnetic Current Model

Assume:

$$E_z^{1,0} = -\sin\left(\frac{\pi x}{L}\right)$$

Measured from lower left corner:

$$E_z^{1,0} = \cos\left(\frac{\pi x'}{L}\right)$$

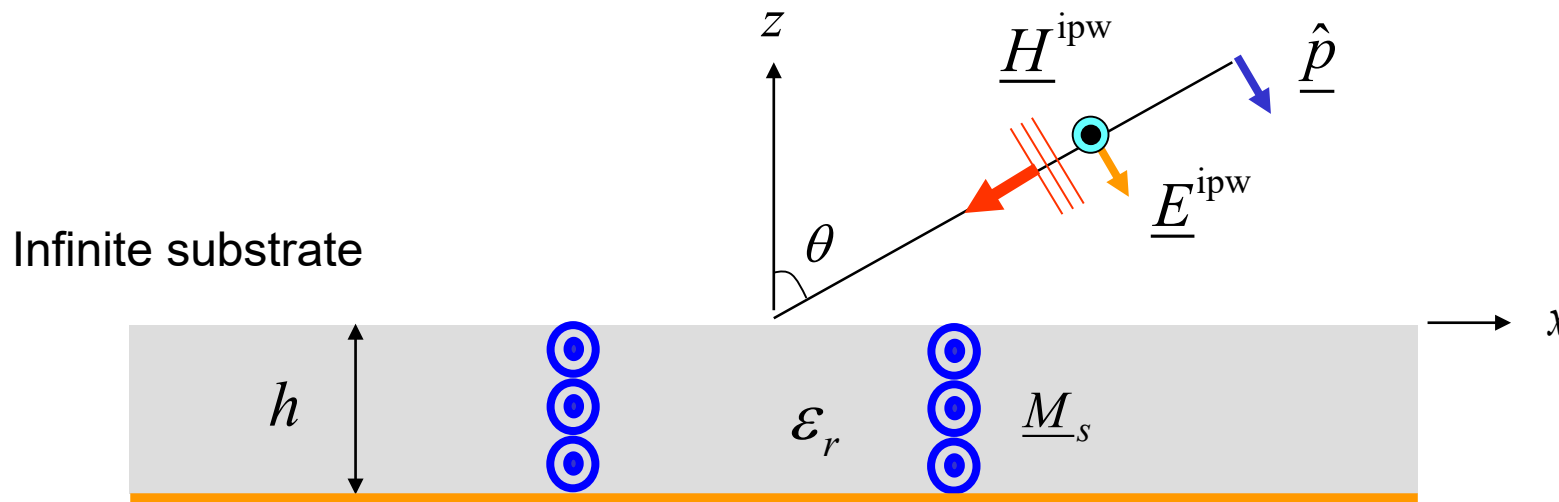


Only the radiating edges contribute to the E- and H-plane patterns, so we will ignore the non-radiating edges.

Radiating edges:

$$\underline{M}_s = -\hat{y}$$

($\hat{p} = \hat{\theta}$ or $\hat{\phi}$)
(shown for $\hat{p} = \hat{\theta}$)

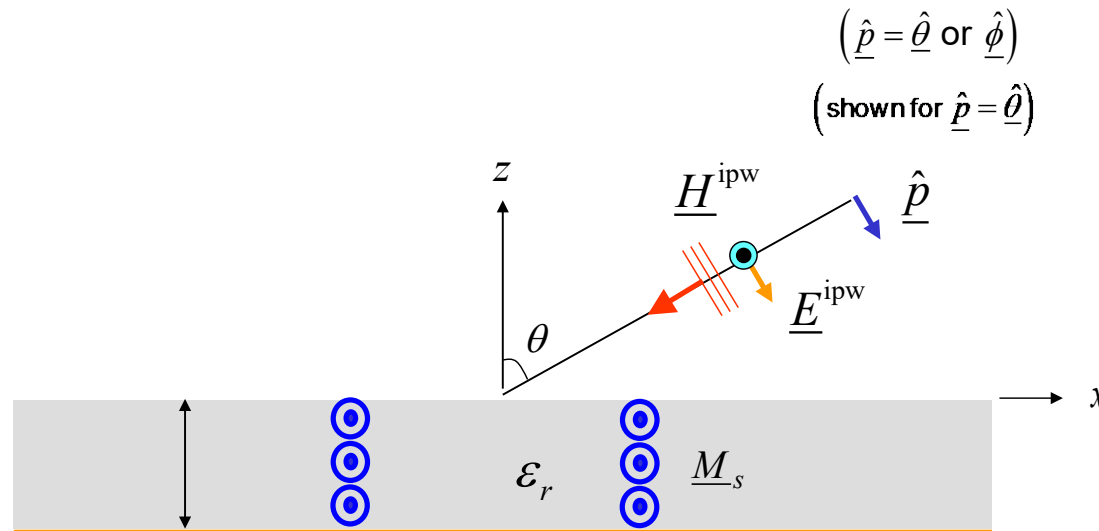


$$\underline{M}_s = \begin{cases} -\hat{y} & x = L \\ -\hat{y} & x = 0 \\ -\hat{x} \cos\left(\frac{\pi x}{L}\right) & y = W \\ \hat{x} \cos\left(\frac{\pi x}{L}\right) & y = 0 \end{cases}$$

Magnetic Current Model (cont.)

From reciprocity:

$$\begin{aligned}
 E_p^{\text{FF}}(r, \theta, \phi) &= \langle a, b \rangle \\
 &= \langle b, a \rangle \\
 &= - \int_S \left(\underline{H}^{\text{pw}} \cdot \underline{M}_s^a \right) dS \\
 &= - \int_S H_y^{\text{pw}} M_{sy}^a dS
 \end{aligned}$$



a = radiating magnetic current
 b = testing dipole ($Il = 1$)

$$S = S_L + S_R \quad (\text{left+right edges})$$

$$M_{sy} = -1$$

Magnetic Current Model (cont.)

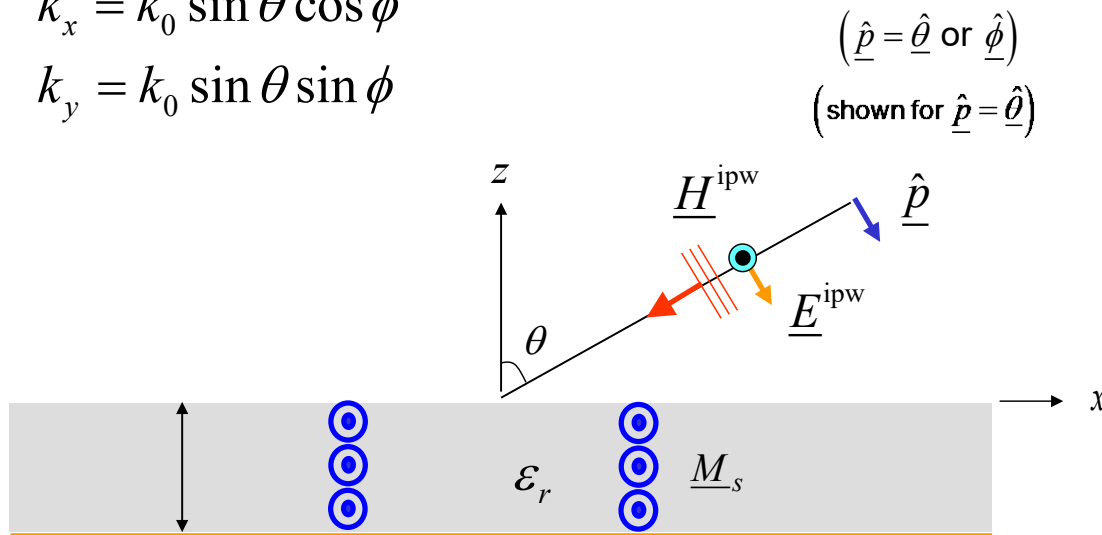
TEN modeling equation:

$$H_y^{\text{pw}}(x, y, z) = \psi_t(x, y) I(z)$$

$$\psi_t(x, y) = e^{j(k_x x + k_y y)}$$

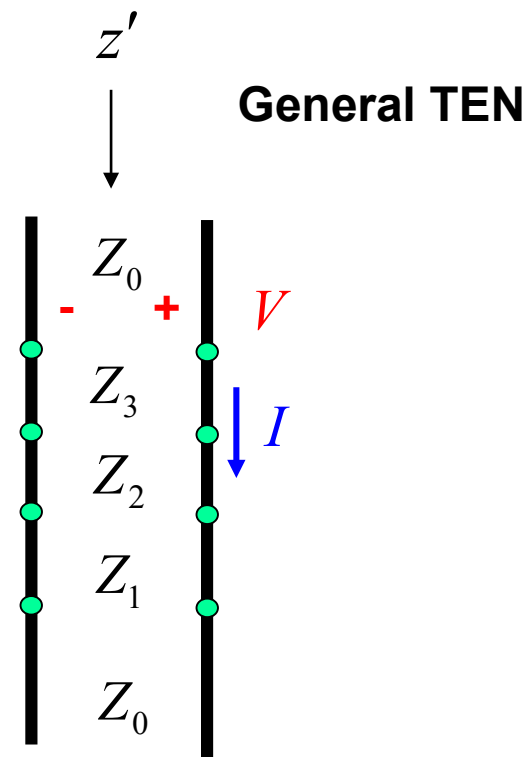
$$k_x = k_0 \sin \theta \cos \phi$$

$$k_y = k_0 \sin \theta \sin \phi$$



$$E_x(x, y, z) = \psi_t(x, y) V(z)$$

$$-H_y(x, y, z) = \psi_t(x, y) I(z)$$



Magnetic Current Model (cont.)

$$\begin{aligned}
 -H_y^{\text{pw}}(x, y, z) &= \psi_t(x, y) I(z) = \psi_t(x, y) I(0) \left[\frac{\cos(k_{z1}(z+h))}{\cos(k_{z1}h)} \right] \\
 &= -H_y^{\text{pw}}(x, y, 0) \cos(k_{z1}(z+h)) \sec(k_{z1}h)
 \end{aligned}$$

$$k_{z1} = k_0 N_1(\theta)$$

$$N_1(\theta) \equiv \sqrt{n_1^2 - \sin^2(\theta)}$$

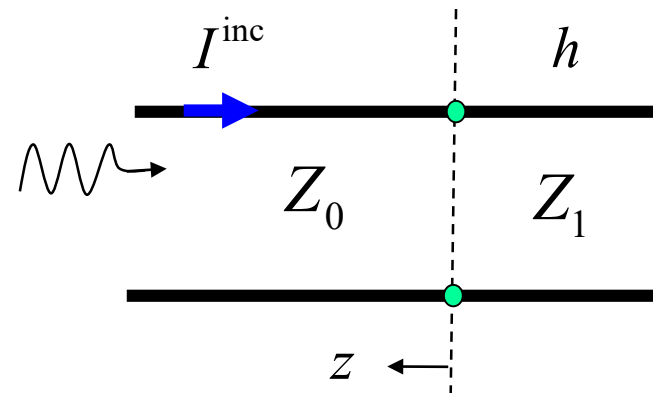
$$n_1 = \sqrt{\epsilon_r \mu_r}$$

$$H_y^{\text{pw}}(x, y, z) = H_y^{\text{pw}}(x, y, 0) \cos(k_{z1}(z+h)) \sec(k_{z1}h)$$

Also,

$$H_y^{\text{pw}}(x, y, 0) = H_y^{\text{ipw}}(x, y, 0) (1 - \Gamma^{\text{TM/TE}})$$

$$I(0) = I^{\text{inc}}(0) (1 - \Gamma^{\text{TM/TE}})$$



$$\text{TM}_z : \underline{\hat{p}} = \underline{\hat{\theta}}$$

$$\text{TE}_z : \underline{\hat{p}} = \underline{\hat{\phi}}$$

Magnetic Current Model (cont.)

$$\underline{\hat{p}} = \underline{\hat{\theta}} :$$

$$\underline{H}^{\text{ipw}} = -\underline{\hat{\phi}} \frac{E_{\theta}^{\text{ipw}}}{\eta_0} = -\underline{\hat{\phi}} \left(\frac{E_0}{\eta_0} \right) \psi(x, y, z) = -\underline{\hat{\phi}} \left(\frac{E_0}{\eta_0} \right) \psi_t(x, y) e^{jk_z z}$$

$$H_y^{\text{ipw}}(x, y, 0) = \left(\frac{E_0}{\eta_0} \right) (-\cos \phi) \psi_t(x, y)$$

$$(\underline{\hat{\phi}} \cdot \underline{\hat{y}} = \cos \phi)$$

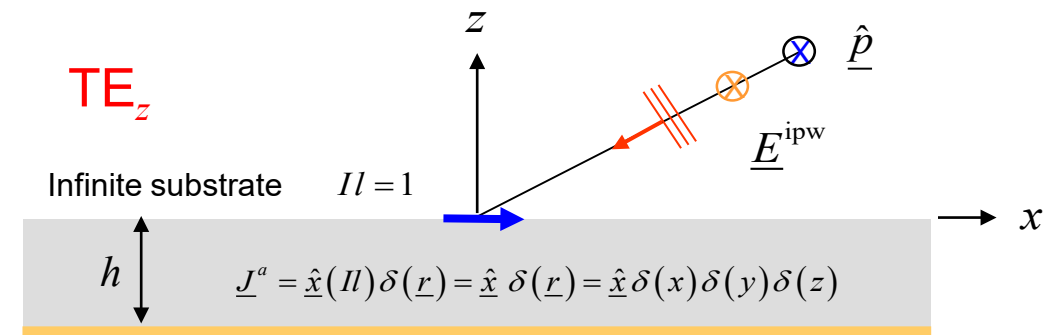
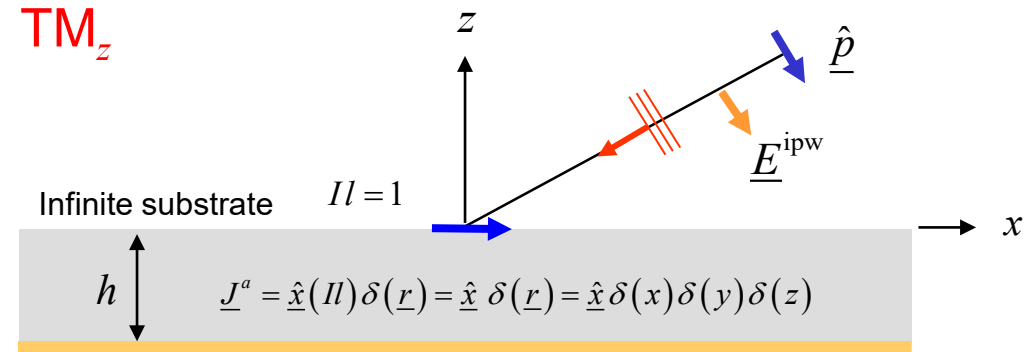
$$\underline{\hat{p}} = \underline{\hat{\phi}} :$$

$$\underline{H}^{\text{ipw}} = \underline{\hat{\theta}} \frac{E_{\phi}^{\text{ipw}}}{\eta_0} = \underline{\hat{\theta}} \left(\frac{E_0}{\eta_0} \right) \psi(x, y, z) = \underline{\hat{\theta}} \left(\frac{E_0}{\eta_0} \right) \psi_t(x, y) e^{jk_z z}$$

$$H_y^{\text{ipw}}(x, y, 0) = \left(\frac{E_0}{\eta_0} \right) (\cos \theta \sin \phi) \psi_t(x, y)$$

$$(\underline{\hat{\theta}} \cdot \underline{\hat{y}} = \cos \theta \sin \phi)$$

$$E_0 = -\left(\frac{j\omega\mu_0}{4\pi r} \right) e^{-jk_0 r}$$

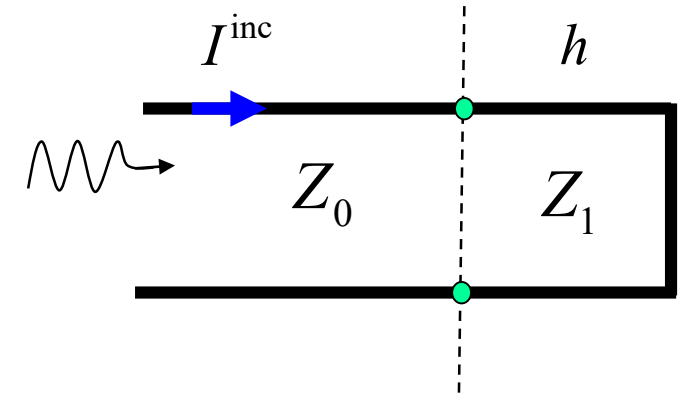


Magnetic Current Model (cont.)

$$\underline{\hat{p}} = \underline{\hat{\theta}} :$$

$$H_y^{\text{ipw}}(x, y, 0) = \left(\frac{E_0}{\eta_0} \right) (-\cos \phi) \psi_t(x, y)$$

$$I^{\text{inc}} = \left(\frac{E_0}{\eta_0} \right) (\cos \phi), E_\theta \text{ (TM}_z\text{)}$$



$$\underline{\hat{p}} = \underline{\hat{\phi}} :$$

$$H_y^{\text{ipw}}(x, y, 0) = \left(\frac{E_0}{\eta_0} \right) (\cos \theta \sin \phi) \psi_t(x, y)$$

$$I^{\text{inc}} = - \left(\frac{E_0}{\eta_0} \right) (\cos \theta \sin \phi), E_\phi \text{ (TE}_z\text{)}$$

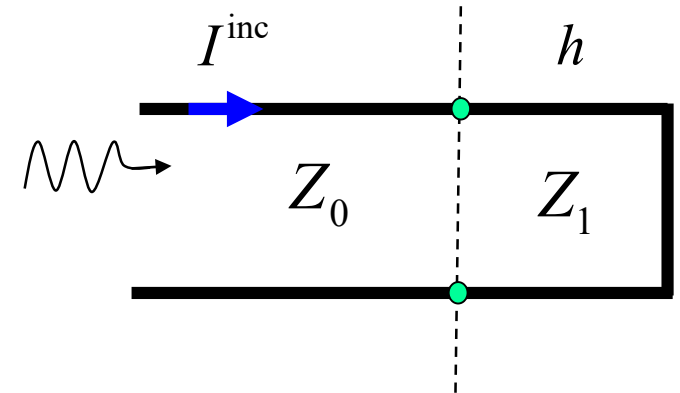
Magnetic Current Model (cont.)

Therefore, we have:

$$-H_y^{\text{pw}}(x, y, z) = e^{j(k_x x + k_y y)} I^{\text{inc}} (1 - \Gamma^{\text{TM/TE}}) \sec(k_{z1} h) \cos k_{z1} (z + h)$$

$$I^{\text{inc}} = \left(\frac{E_0}{\eta_0} \right) (\cos \phi), E_\theta \text{ (TM}_z \text{)}$$

$$I^{\text{inc}} = - \left(\frac{E_0}{\eta_0} \right) (\cos \theta \sin \phi), E_\phi \text{ (TE}_z \text{)}$$



We can thus now consider

$$E_p^{\text{FF}}(r, \theta, \phi) = - \int_{S_L + S_R} H_y^{\text{pw}}(x, y, z) (-1) dS$$

$M_{sy} = -1$

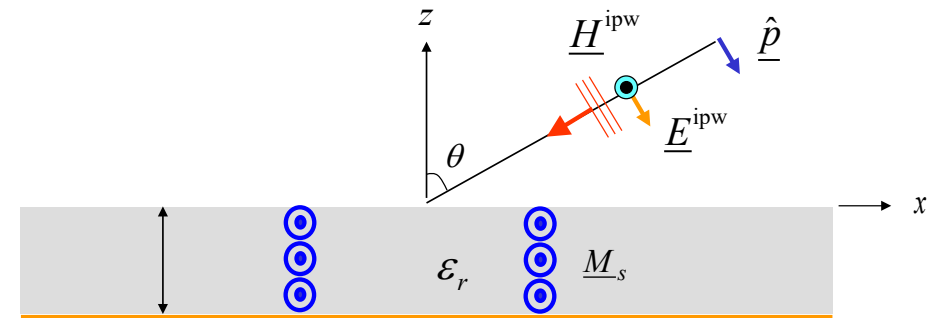
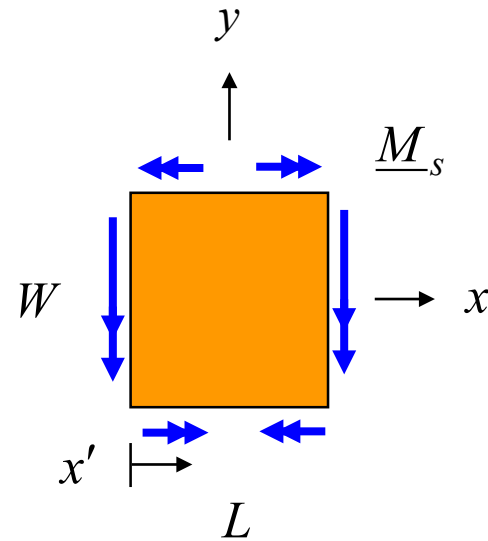
Magnetic Current Model (cont.)

Consider the right edge S_R :

$$x = \frac{L}{2}$$

$$y \in \left(-\frac{W}{2}, \frac{W}{2}\right)$$

$$z \in (-h, 0)$$



Therefore, we have:

$$E_p^R(r, \theta, \phi) = - \int_{-h}^0 \int_{-\frac{W}{2}}^{\frac{W}{2}} e^{j\left(k_x \frac{L}{2}\right)} e^{jk_y y} I^{\text{inc}} \left(1 - \Gamma^{\text{TM/TE}}\right) \sec(k_{z1} h) \cos k_{z1} (z + h) dy dz$$

Magnetic Current Model (cont.)

$$E_p^R(r, \theta, \phi) = -I^{\text{inc}} (1 - \Gamma^{\text{TM/TE}}) \sec(k_{z1}h) e^{j\left(k_x \frac{L}{2}\right)} \underbrace{\int_{-\frac{W}{2}}^{+\frac{W}{2}} e^{jk_y y} dy}_{\text{integral1}} \underbrace{\int_{-h}^0 \cos k_{z1}(z+h) dz}_{\text{integral2}}$$

Integral 1: $W \operatorname{sinc}\left(\frac{k_y W}{2}\right)$

Integral 2: $h \operatorname{sinc}(k_{z1}h)$

Hence, we have:

$$E_p^R(r, \theta, \phi) = -I^{\text{inc}} (1 - \Gamma^{\text{TM/TE}}) e^{jk_x \frac{L}{2}} (Wh) \operatorname{sinc}\left(\frac{k_y W}{2}\right) \operatorname{tanc}(k_{z1}h)$$

where

$$\operatorname{tanc}(x) \equiv \frac{\tan x}{x} \quad (\sec(x) \operatorname{sinc}(x) = \operatorname{tanc}(x))$$

Magnetic Current Model (cont.)

For S_L : Replace $\frac{L}{2} \rightarrow -\frac{L}{2}$ $\Rightarrow e^{jk_x \frac{L}{2}} \rightarrow e^{-jk_x \frac{L}{2}}$

For $S_L + S_R$: $e^{jk_x \frac{L}{2}} + e^{-jk_x \frac{L}{2}} \Rightarrow 2 \cos\left(k_x \frac{L}{2}\right)$

Hence, we have:

$$E_p^{\text{FF}}(r, \theta, \phi) = -I^{\text{inc}} (1 - \Gamma^{\text{TM/TE}}) 2 \cos\left(k_x \frac{L}{2}\right) (Wh) \text{sinc}\left(k_y \frac{W}{2}\right) \text{tanc}(k_{z1} h)$$

$$I^{\text{inc}} = \left(\frac{E_0}{\eta_0}\right) (\cos \phi), E_\theta \text{ (TM}_z\text{)}$$

$$I^{\text{inc}} = -\left(\frac{E_0}{\eta_0}\right) (\cos \theta \sin \phi), E_\phi \text{ (TE}_z\text{)}$$

Magnetic Current Model (cont.)

Substituting for J^{inc} in the TM_z and TE_z cases, we have:

$$E_{\theta}^{\text{FF}}(r, \theta, \phi) = -2Wh \left(\frac{E_0}{\eta_0} \right) \cos \phi (1 - \Gamma^{\text{TM}}(\theta)) \cos \left(k_x \frac{L}{2} \right) \text{sinc} \left(k_y \frac{W}{2} \right) \text{tanc}(k_{z1}h)$$

$$E_{\phi}^{\text{FF}}(r, \theta, \phi) = 2Wh \left(\frac{E_0}{\eta_0} \right) (\cos \theta \sin \phi) (1 - \Gamma^{\text{TE}}(\theta)) \cos \left(k_x \frac{L}{2} \right) \text{sinc} \left(k_y \frac{W}{2} \right) \text{tanc}(k_{z1}h)$$

$$E_0 = \left(\frac{-j\omega\mu_0}{4\pi r} \right) e^{-jk_0 r} \quad \begin{aligned} k_x &= k_0 \sin \theta \cos \phi \\ k_y &= k_0 \sin \theta \sin \phi \end{aligned}$$

Magnetic Current Model (cont.)

Examine the reflection coefficient terms:

$$1 - \Gamma = 1 - \left(\frac{Z_{\text{in}} - Z_0}{Z_{\text{in}} + Z_0} \right) \\ = \frac{2Z_0}{Z_{\text{in}} + Z_0}$$

$$\Gamma = \Gamma^{\text{TM}} \quad \text{for} \quad \underline{\hat{p}} = \underline{\hat{\theta}}$$

$$\Gamma = \Gamma^{\text{TE}} \quad \text{for} \quad \underline{\hat{p}} = \underline{\hat{\phi}}$$

$$Z_{\text{in}}^{\text{TM/TE}} = jZ_1^{\text{TM/TE}} \tan(k_{z1}h) \\ = jZ_1^{\text{TM/TE}} \tan(k_0 N_1(\theta)h)$$

$$N_1(\theta) \equiv \sqrt{n_1^2 - \sin^2 \theta} = \sqrt{\epsilon_r \mu_r - \sin^2 \theta}$$

$$Z_1^{\text{TM}} = \frac{\eta_0}{\epsilon_r} N_1(\theta)$$

$$Z_1^{\text{TE}} = \frac{\eta_0 \mu_r}{N_1(\theta)}$$

$$Z_0^{\text{TM}} = \eta_0 \cos \theta$$

$$Z_0^{\text{TE}} = \frac{\eta_0}{\cos \theta}$$

Magnetic Current Model (cont.)

Substituting the impedance formulas into the formulas for the reflection coefficients, and then simplifying the results, we have:

$$1 - \Gamma^{\text{TM}}(\theta) = \frac{2}{1 + j \left(\frac{N_1(\theta) \sec \theta}{\epsilon_r} \right) \tan(k_0 h N_1(\theta))}$$

$$1 - \Gamma^{\text{TE}}(\theta) = \frac{2}{1 + j \left(\frac{\mu_r \cos \theta}{N_1(\theta)} \right) \tan(k_0 h N_1(\theta))}$$

Summary

$$E_{\theta}^{\text{FF}}(r, \theta, \phi) = -2Wh \left(\frac{E_0}{\eta_0} \right) \cos \phi (1 - \Gamma^{\text{TM}}(\theta)) \cos \left(k_x \frac{L}{2} \right) \text{sinc} \left(k_y \frac{W}{2} \right) \text{tanc}(k_{z1}h)$$

$$E_{\phi}^{\text{FF}}(r, \theta, \phi) = 2Wh \left(\frac{E_0}{\eta_0} \right) (\cos \theta \sin \phi) (1 - \Gamma^{\text{TE}}(\theta)) \cos \left(k_x \frac{L}{2} \right) \text{sinc} \left(k_y \frac{W}{2} \right) \text{tanc}(k_{z1}h)$$

$$1 - \Gamma^{\text{TM}}(\theta) = \frac{2}{1 + j \left(\frac{N_1(\theta) \sec \theta}{\epsilon_r} \right) \tan(k_0 h N_1(\theta))}$$

$$1 - \Gamma^{\text{TE}}(\theta) = \frac{2}{1 + j \left(\frac{\mu_r \cos \theta}{N_1(\theta)} \right) \tan(k_0 h N_1(\theta))}$$

$$E_0 = \left(\frac{-j\omega\mu_0}{4\pi r} \right) e^{-jk_0 r}$$

$$k_x = k_0 \sin \theta \cos \phi$$

$$k_y = k_0 \sin \theta \sin \phi$$

$$k_{z1} = k_0 N_1(\theta)$$

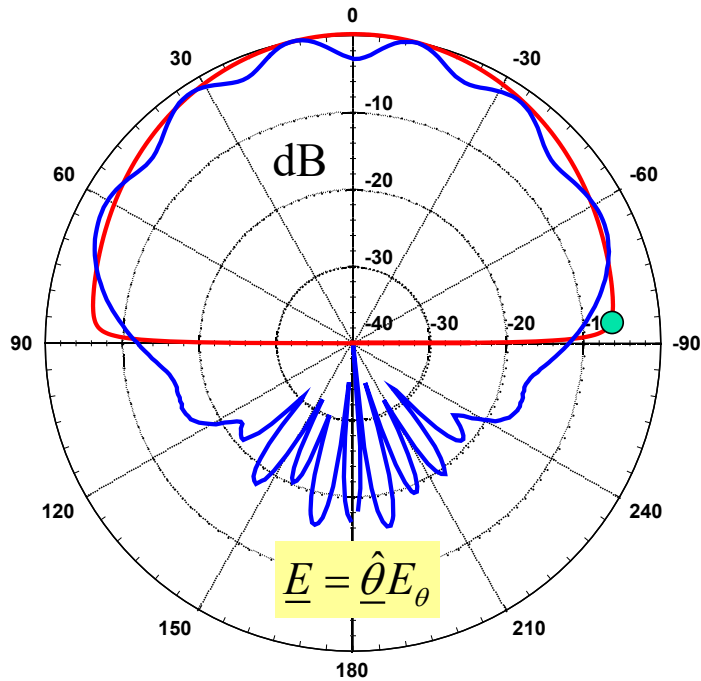
$$N_1(\theta) \equiv \sqrt{n_1^2 - \sin^2 \theta}$$

$$n_1 = \sqrt{\mu_r \epsilon_r}$$

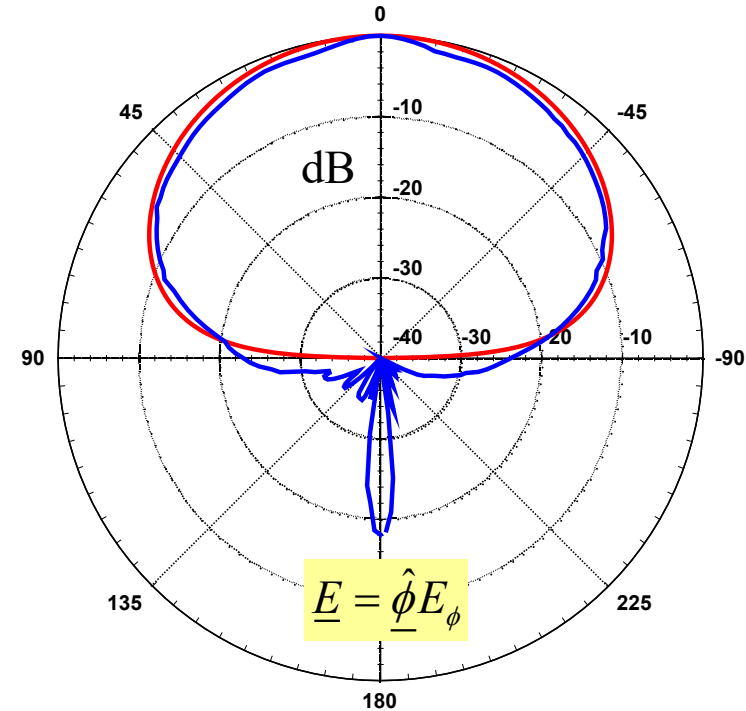
Assumption: $E_z^{1,0} = -\sin \left(\frac{\pi x}{L} \right)$

Summary

E-plane pattern ($\phi = 0^\circ$)



H-plane pattern ($\phi = 90^\circ$)



Comments:

Red: infinite substrate and ground plane

Blue: 1 meter diameter ground plane

- The E plane is broader than the H plane.
- The E-plane pattern “tucks in” and tends to zero at the horizon due to the presence of the infinite substrate (green dot). (As the substrate gets thinner, the tuck-in point approaches 90° .)