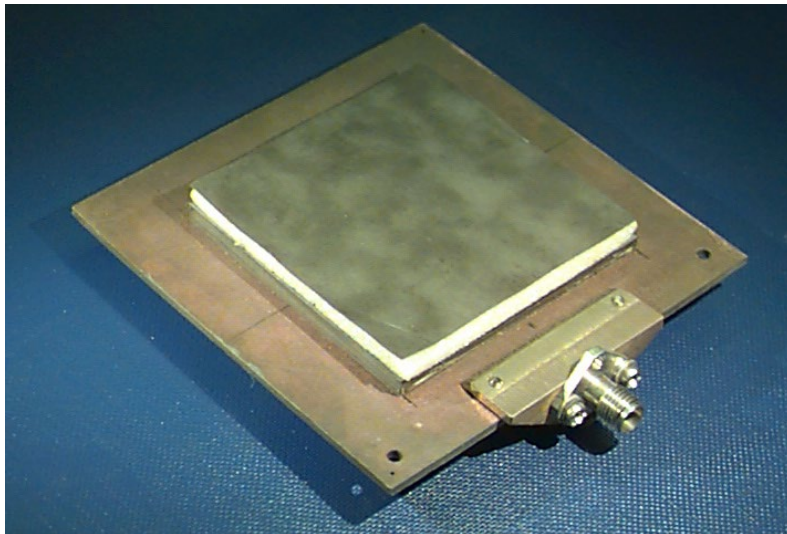


ECE 6345

Spring 2024

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ECE Dept.

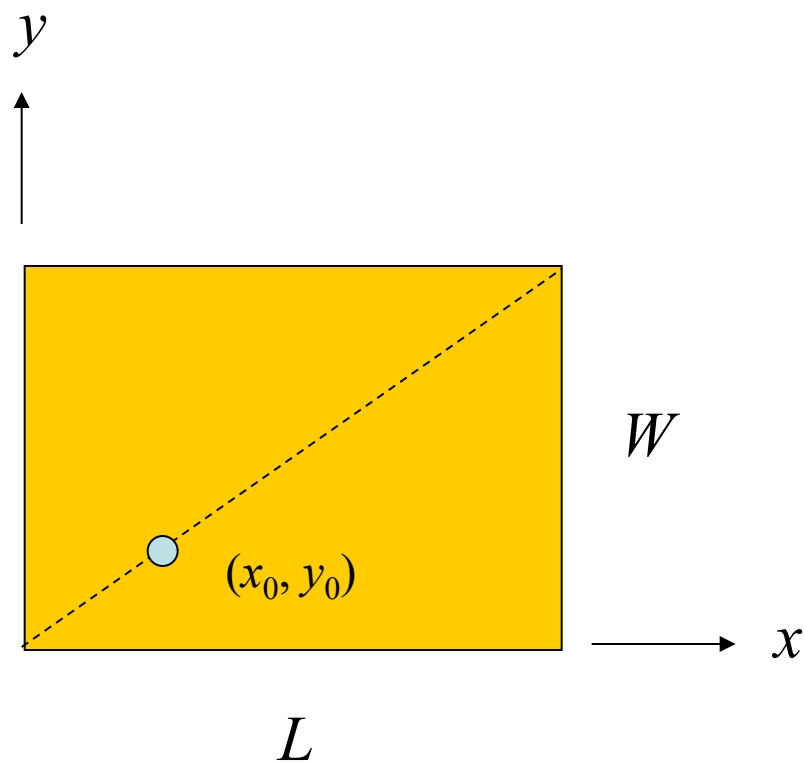


Notes 6

Overview

This set of notes treats **circular polarization**, obtained by using a **single feed**.

The nearly-square patch is fed with a coaxial feed along the diagonal.



$$L \approx W$$

$$y_0 \approx x_0$$

Overview

Goals:

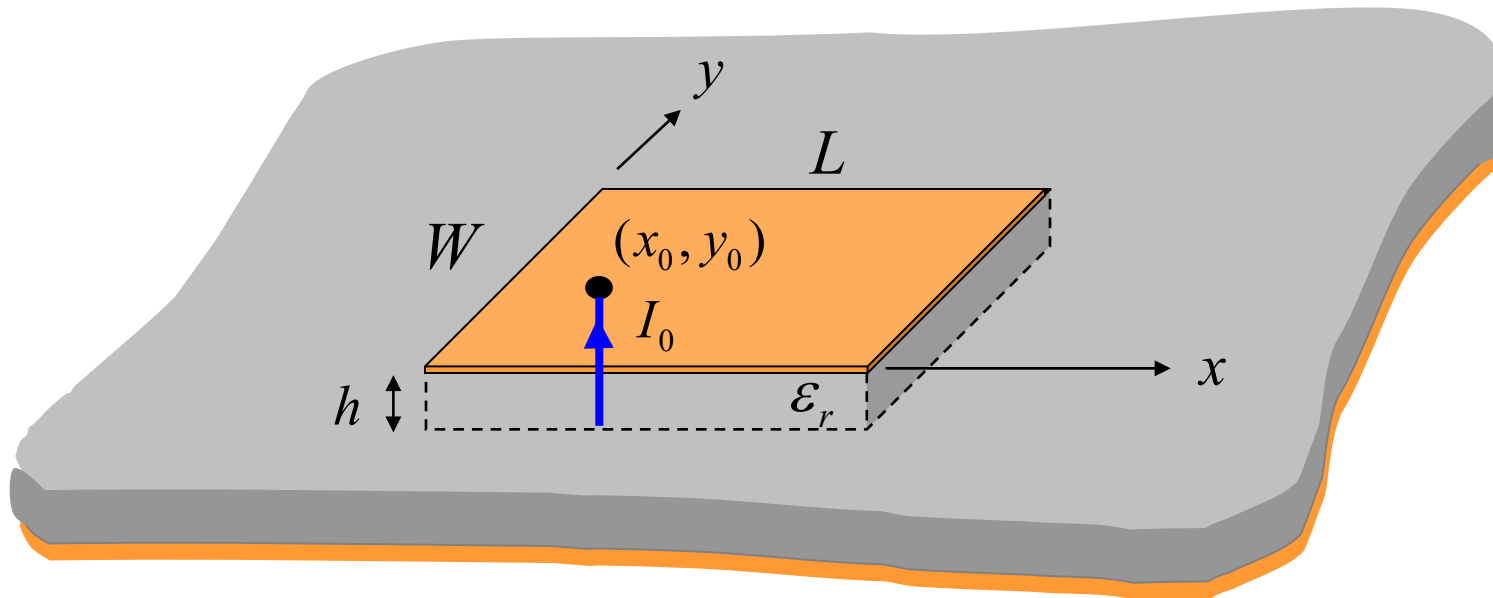
- Find the optimum dimensions of the CP patch
- Find the input impedance of the CP patch
- Find the pattern (axial-ratio) bandwidth of the CP patch
- Find the impedance bandwidth of the CP patch

We use the electric current model, and thus work with the currents on the patch.

Amplitude of Patch Currents

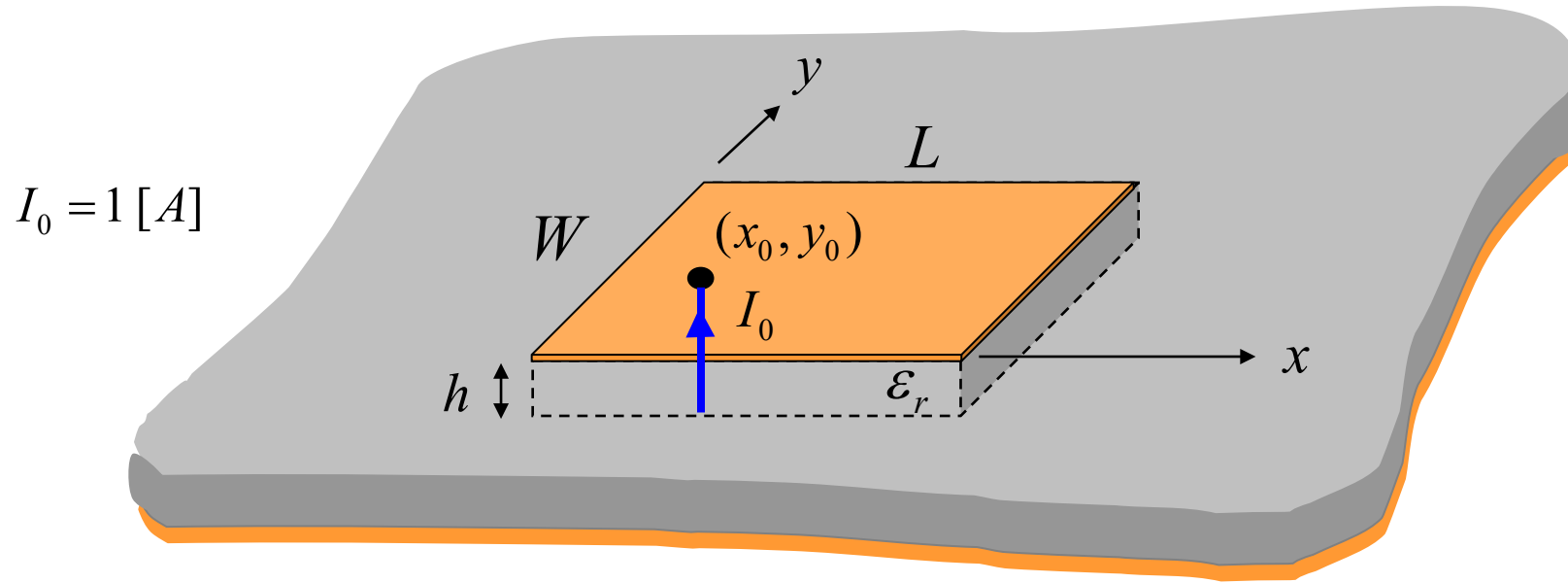
First Step:

Find the patch currents (x and y directions), and then relate them to the input impedance of the patch.



Assume: $I_0 = 1 [A]$

Amplitude of Patch Currents (cont.)



Patch currents:

x -directed current mode (TM_{10}):
$$\underline{J}_s^x = \hat{x} A_x \sin\left(\frac{\pi x}{L}\right)$$

y -directed current mode (TM_{01}):
$$\underline{J}_s^y = \hat{y} A_y \sin\left(\frac{\pi y}{W}\right)$$

Amplitude of Patch Currents (cont.)

The x mode (TM₁₀):

$$\underline{J}_s = \underline{\hat{n}} \times \underline{H} = -\underline{\hat{z}} \times \underline{H}$$

so $J_{sx} = H_y$

Assume: $\underline{H} = \underline{\hat{y}} A_x \sin\left(\frac{\pi x}{L}\right)$

To find \underline{E} , use $\nabla \times \underline{H} = j\omega\varepsilon \underline{E}$

$$E_z = \frac{1}{j\varepsilon_0\varepsilon_r} \left[\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right] = \frac{1}{j\varepsilon_0\varepsilon_r} A_x \left(\frac{\pi}{L} \right) \cos\left(\frac{\pi x}{L}\right)$$

Amplitude of Patch Currents (cont.)

$$Z_{\text{in}} = \frac{V}{I_0} = V = -hE_z = -h(E_z^x + E_z^y) = Z_{\text{in}}^x + Z_{\text{in}}^y$$

For the TM_{10} mode we have, from the last slide:

$$Z_{\text{in}}^x = \frac{jh}{\omega \epsilon_0 \epsilon_r} \left(\frac{\pi}{L} \right) A_x \cos \left(\frac{\pi x_0}{L} \right)$$

And thus, solving for A_x :

$$A_x = \frac{Z_{\text{in}}^x}{\cos \left(\frac{\pi x_0}{L} \right)} \left[\frac{\omega \epsilon_0 \epsilon_r L}{j\pi h} \right]$$

A similar derivation holds for the y mode (TM_{01}).

Amplitude of Patch Currents (cont.)

The y mode (TM_{01}):

$$A_y = \frac{Z_{\text{in}}^y}{\cos\left(\frac{\pi y_0}{W}\right)} \left[\frac{\omega \epsilon_0 \epsilon_r W}{j\pi h} \right]$$

The patch current amplitudes can then be written as:

$$A_x = A_1^x Z_{\text{in}}^x$$

$$A_y = A_1^y Z_{\text{in}}^y$$

$$\underline{J}_s^x = \underline{\hat{x}} A_x \sin\left(\frac{\pi x}{L}\right)$$

$$\underline{J}_s^y = \underline{\hat{y}} A_y \sin\left(\frac{\pi y}{W}\right)$$

where

$$A_1^x = \frac{1}{\cos\left(\frac{\pi x_0}{L}\right)} \left[\frac{\omega \epsilon_0 \epsilon_r L}{j\pi h} \right]$$

$$A_1^y = \frac{1}{\cos\left(\frac{\pi y_0}{W}\right)} \left[\frac{\omega \epsilon_0 \epsilon_r W}{j\pi h} \right]$$

Amplitude of Patch Currents (cont.)

$$A_1^x = \frac{1}{\cos\left(\frac{\pi x_0}{L}\right)} \left[\frac{\omega \varepsilon_0 \varepsilon_r L}{j\pi h} \right]$$

$$A_1^y = \frac{1}{\cos\left(\frac{\pi y_0}{W}\right)} \left[\frac{\omega \varepsilon_0 \varepsilon_r W}{j\pi h} \right]$$

Assume:

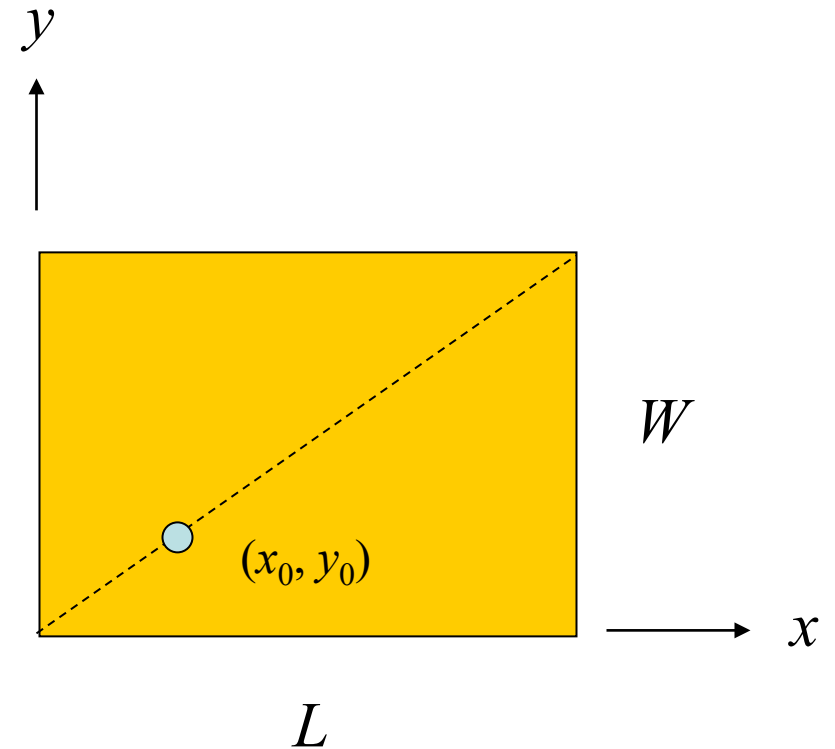
$$L \approx W$$
$$x_0 \approx y_0$$

Then $A_1^x \approx A_1^y \equiv A_1$

$$A_x = A_1 Z_{in}^x$$

$$A_y = A_1 Z_{in}^y$$

$$A_1 = \frac{1}{\cos\left(\frac{\pi x_0}{L}\right)} \left[\frac{\omega \varepsilon_0 \varepsilon_r L}{j\pi h} \right]$$



Amplitude of Patch Currents (cont.)

Because of the nearly equal dimensions and the feed along the diagonal, we also have:

$$R_{\text{in}}^x \approx R_{\text{in}}^y \equiv R \quad (\text{The resistor in the circuit model for either mode is the same.})$$

R_{in}^x = resonant input resistance of the mode x , when excited by itself, and similarly for R_{in}^y

(R_{in}^x only depends on x_0 , R_{in}^y only depends on y_0 .)

We then have:

$$A_x = A_2 \bar{Z}_{\text{in}}^x$$

$$A_y = A_2 \bar{Z}_{\text{in}}^y$$

where $A_2 \equiv A_1 R$

$$\Rightarrow A_2 = \frac{R}{\cos\left(\frac{\pi x_0}{L}\right)} \left[\frac{\omega \varepsilon_0 \varepsilon_r L}{j\pi h} \right]$$

Notation: The bar denotes impedances that are normalized by R .

Amplitude of Patch Currents (cont.)

The A_2 coefficient can be written as:

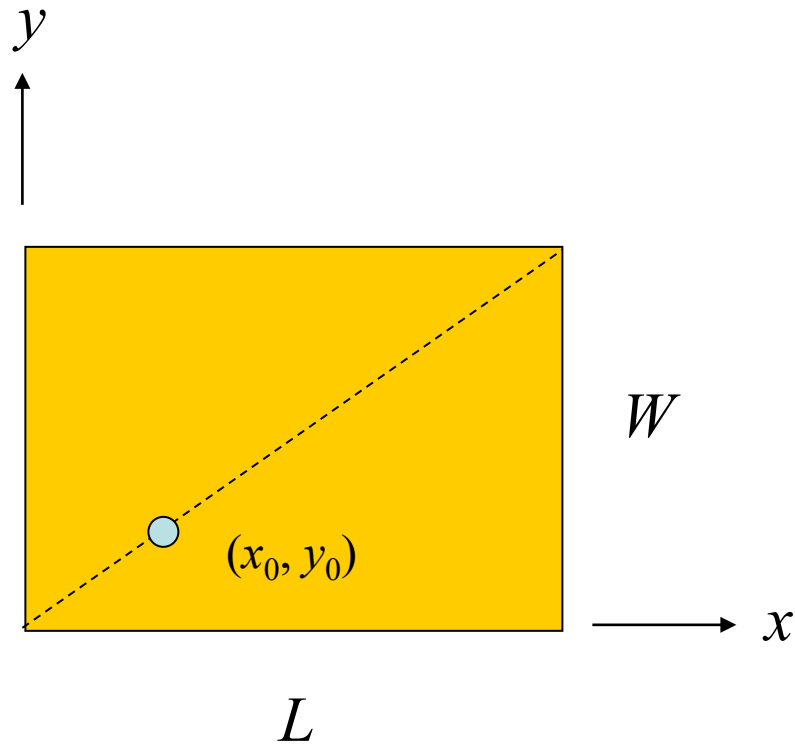
$$A_2 = \frac{R}{\cos\left(\frac{\pi x_0}{L}\right)} \left[\frac{\omega \varepsilon_0 \varepsilon_r L}{j\pi h} \right]$$
$$= \frac{R_{\text{edge}} \cos^2\left(\frac{\pi x_0}{L}\right)}{\cos\left(\frac{\pi x_0}{L}\right)} \left[\frac{\omega \varepsilon_0 \varepsilon_r L}{j\pi h} \right]$$

Note:
We can use the CAD formula for R_{edge} .

so

$$A_2 = R_{\text{edge}} \cos\left(\frac{\pi x_0}{L}\right) \left[\frac{\omega \varepsilon_0 \varepsilon_r L}{j\pi h} \right]$$

Circular Polarization Condition



Let

$$L = W(1 + \delta)$$

$$y_0 \approx x_0$$

The CP condition is

$$\frac{A_y}{A_x} = \pm j$$

A_x = amplitude of x mode current (TM_{10})

A_y = amplitude of y mode current (TM_{01})

- for RHCP
+ for LHCP

Circular Polarization Condition (cont.)

The frequency f_{CP} is defined as the frequency for which we get CP at broadside.

We have, from the CAD circuit model:

$$\bar{Z}_{in} = \frac{1}{1 + j2Q(f_r - 1)}$$

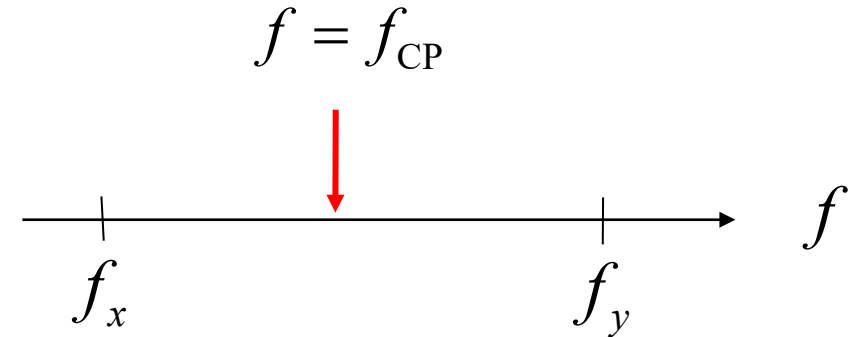
$$A_x = A_2 \bar{Z}_{in}^x$$

$$A_y = A_2 \bar{Z}_{in}^y$$

So, for each mode we have:

$$A_x = \frac{A_2}{1 + j2Q(f_{rx} - 1)}$$

$$A_y = \frac{A_2}{1 + j2Q(f_{ry} - 1)}$$



where

$$f_{rx} \equiv \frac{f}{f_x} \quad f_{ry} \equiv \frac{f}{f_y}$$

f_x = resonance frequency of TM_{10} mode

f_y = resonance frequency of TM_{01} mode

Circular Polarization Condition (cont.)

LHCP:

$$\frac{A_y}{A_x} = +j$$

at $f = f_{CP}$

Choose:

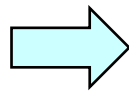
$$f_{rx} - 1 = \frac{1}{2Q}$$

$$f_{ry} - 1 = -\frac{1}{2Q}$$

Then we have:

$$A_x = \frac{A}{1+j}$$

$$A_y = \frac{A}{1-j}$$



$$\frac{A_y}{A_x} = \frac{1+j}{1-j} = \frac{\sqrt{2}e^{j\frac{\pi}{4}}}{\sqrt{2}e^{-j\frac{\pi}{4}}} = e^{j\frac{\pi}{2}} = j$$

(LHCP)

Circular Polarization Condition (cont.)

The frequency conditions for $f = f_{\text{CP}}$ can be written as:

$$f_{rx} - 1 = \frac{1}{2Q} \quad \Rightarrow \quad \frac{f_{\text{CP}}}{f_x} = 1 + \frac{1}{2Q} \quad \Rightarrow \quad \frac{f_x}{f_{\text{CP}}} \approx 1 - \frac{1}{2Q}$$

$$f_{ry} - 1 = -\frac{1}{2Q} \quad \Rightarrow \quad \frac{f_{\text{CP}}}{f_y} = 1 - \frac{1}{2Q} \quad \Rightarrow \quad \frac{f_y}{f_{\text{CP}}} \approx 1 + \frac{1}{2Q}$$

Note:

$$\frac{1}{1+x} \approx 1-x$$

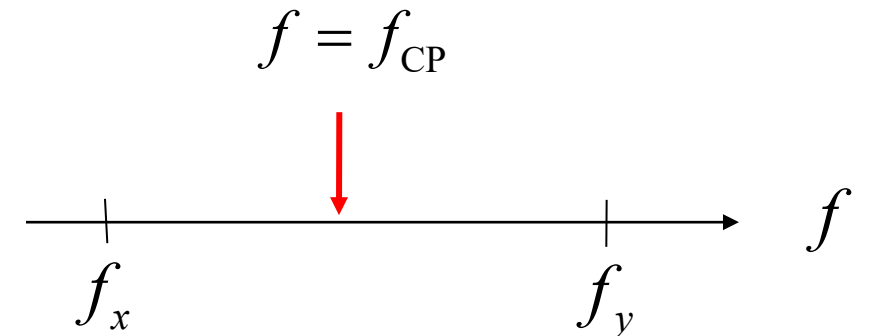
$$\frac{1}{1-x} \approx 1+x$$

so

$$\frac{f_x + f_y}{f_{\text{CP}}} = 2$$

or

$$f_{\text{CP}} = \frac{1}{2}(f_x + f_y)$$



Circular Polarization Condition (cont.)

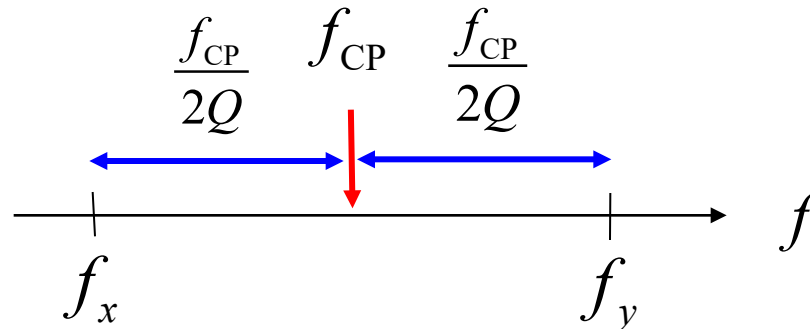
Also,
$$\frac{f_y}{f_{\text{CP}}} - \frac{f_x}{f_{\text{CP}}} = \left(1 + \frac{1}{2Q}\right) - \left(1 - \frac{1}{2Q}\right) = \frac{1}{Q}$$

Let $\Delta f \equiv f_y - f_x$

Then we have:

$$\frac{\Delta f}{f_{\text{CP}}} = \frac{1}{Q}$$

Note:
For RHCP, the two frequencies are reversed.



Circular Polarization Condition (cont.)

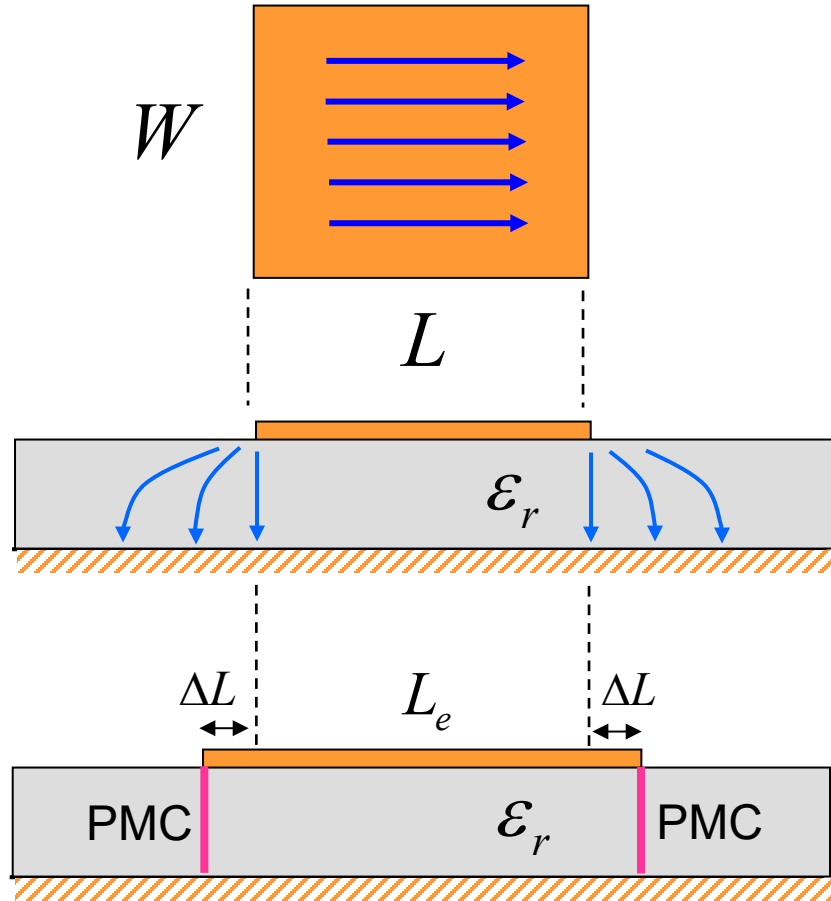
Summary of frequencies

$$f_x = f_{\text{CP}} \left[1 - \frac{1}{2Q} \right] \quad f_y = f_{\text{CP}} \left[1 + \frac{1}{2Q} \right] \quad (\text{LHCP})$$

$$f_x = f_{\text{CP}} \left[1 + \frac{1}{2Q} \right] \quad f_y = f_{\text{CP}} \left[1 - \frac{1}{2Q} \right] \quad (\text{RHCP})$$

f_{CP} = frequency for which we get CP at broadside.

Patch Dimensions for CP



$$\Delta L = 0.412h \left[\frac{\left(\frac{W}{h}\right) + 0.262}{\left(\frac{W}{h}\right) + 0.813} \right] \left[\frac{\epsilon_r^{\text{eff}} + 0.300}{\epsilon_r^{\text{eff}} - 0.258} \right]$$

$$\epsilon_r^{\text{eff}} = \left(\frac{\epsilon_r + 1}{2}\right) + \left(\frac{\epsilon_r - 1}{2}\right) \frac{1}{\sqrt{1 + 12\left(\frac{h}{W}\right)}}$$

$$L_e = L + 2\Delta L$$

$$\Delta L = f(h, \epsilon_r, W) \quad (\text{Hammerstad formula})$$

Patch Dimensions for CP (cont.)

$$f_x = f_{\text{CP}} \left[1 \mp \frac{1}{2Q} \right] \quad f_y = f_{\text{CP}} \left[1 \pm \frac{1}{2Q} \right]$$

Top sign: LHCP
Bottom sign: RHCP

Define:

$$k_{0x} \equiv 2\pi f_x \sqrt{\mu_0 \epsilon_0}$$

(These are known wavenumbers, since the frequencies are known.)

$$k_{0y} \equiv 2\pi f_y \sqrt{\mu_0 \epsilon_0}$$

$$\left. \begin{aligned} k_{0x} L_e \sqrt{\epsilon_r} &= \pi \\ k_{0y} W_e \sqrt{\epsilon_r} &= \pi \end{aligned} \right\} \text{(resonance condition)}$$

$$L_e = \frac{\pi}{k_{0x} \sqrt{\epsilon_r}}$$

$$W_e = \frac{\pi}{k_{0y} \sqrt{\epsilon_r}}$$

Now we know the effective dimensions of the patch.

Patch Dimensions for CP (cont.)

For the physical patch dimensions we have:

$$L = L_e - 2\Delta L(h, \epsilon_r, W)$$

$$W = W_e - 2\Delta W(h, \epsilon_r, L)$$

Note: For ΔW , we use the same formula as ΔL , but replace $W \rightarrow L$.

$$\Delta L = 0.412h \left[\frac{\left(\frac{W}{h}\right) + 0.262}{\left(\frac{W}{h}\right) + 0.813} \right] \left[\frac{\epsilon_r^{\text{eff}} + 0.300}{\epsilon_r^{\text{eff}} - 0.258} \right]$$

$$\Delta W = 0.412h \left[\frac{\left(\frac{L}{h}\right) + 0.262}{\left(\frac{L}{h}\right) + 0.813} \right] \left[\frac{\epsilon_r^{\text{eff}} + 0.300}{\epsilon_r^{\text{eff}} - 0.258} \right]$$

Physical Dimensions for CP (cont.)

Summary of Design Equations

Patch Dimensions

$$L = L_e - 2\Delta L$$

$$W = W_e - 2\Delta W$$

where

$$L_e = \frac{\pi}{k_{0x} \sqrt{\epsilon_r}}$$

$$k_{0x} \equiv 2\pi f_x \sqrt{\mu_0 \epsilon_0}$$

$$W_e = \frac{\pi}{k_{0y} \sqrt{\epsilon_r}}$$

$$k_{0y} \equiv 2\pi f_y \sqrt{\mu_0 \epsilon_0}$$

$$f_x = f_{\text{CP}} \left[1 \mp \frac{1}{2Q} \right] \quad f_y = f_{\text{CP}} \left[1 \pm \frac{1}{2Q} \right]$$

Top sign: LHCP, Bottom sign: RHCP

Note: The fringing length ΔL depends on W , and the fringing length ΔW depends on L .

Physical Dimensions for CP (cont.)

Iterative Design Process:

- Neglect fringing extensions, and take $L = L_e$ and $W = W_e$.
- Solve for the fringing extensions ΔL and ΔW using the Hammerstad formula.
- Solve for $L = L_e - 2\Delta L$ and $W = W_e - 2\Delta W$.
- Repeat from step 2 if necessary.

Input Impedance of CP Patch

$$Z_{\text{in}}^{\text{CP}}(f) = Z_{\text{in}}^x(f) + Z_{\text{in}}^y(f) = \frac{R}{1 + j2Q(f_{rx} - 1)} + \frac{R}{1 + j2Q(f_{ry} - 1)} \quad (R_{\text{in}}^x \approx R_{\text{in}}^y \equiv R)$$

At f_{CP} : $f_{rx} - 1 = 1/(2Q)$ and $f_{ry} - 1 = -1/(2Q)$ (LHCP)

so

$$Z_{\text{in}}^{\text{CP}}(f_{\text{CP}}) = \frac{R}{1 + j} + \frac{R}{1 - j} = \frac{R(1 - j) + R(1 + j)}{(1 + j)(1 - j)} = \frac{R(1 - j) + R(1 + j)}{2} = R$$

or $Z_{\text{in}} = R$

Recall: $R = R^x = R$ for TM_{10} mode

The CP frequency f_{CP} is also the resonance frequency where the input impedance is real (if we neglect the probe inductance).

Input Impedance of CP Patch (cont.)

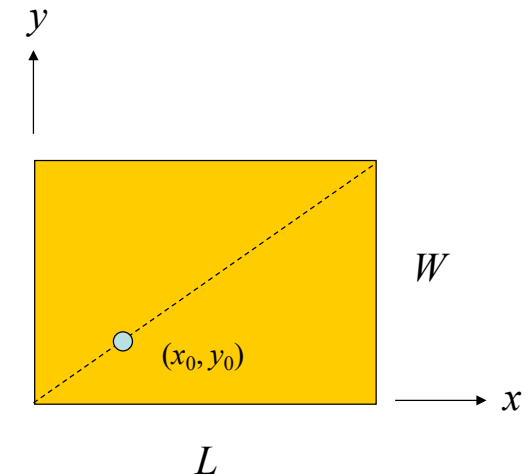
Hence, at the resonance (CP) frequency f_{CP} we have

$$Z_{in}^{CP} = R_{edge} \cos^2 \left(\frac{\pi x_0}{L} \right)$$

$$R_{edge} = R_{edge} \text{ for TM}_{10} \text{ mode}$$

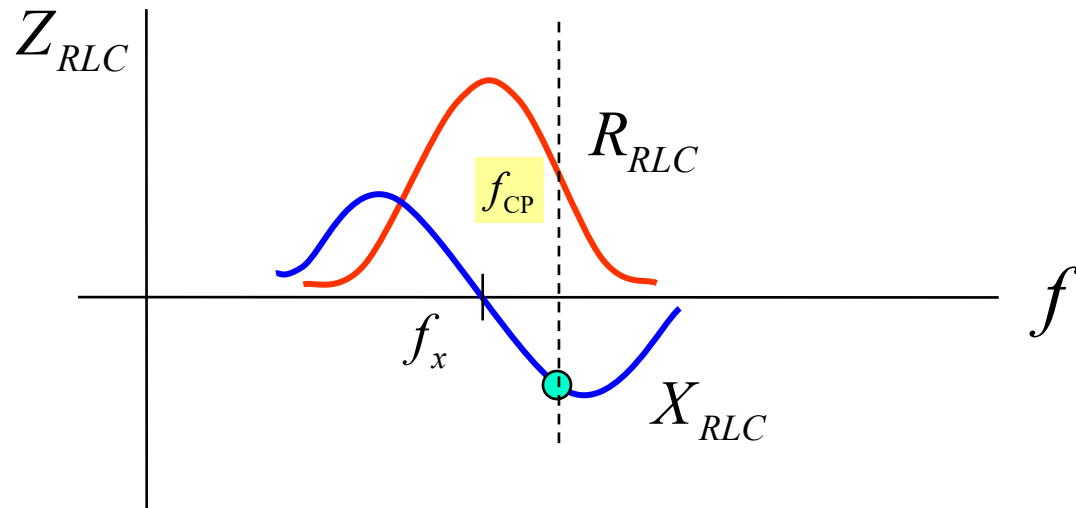
Note: We have a CAD formula for R_{edge} .

The feed position x_0 can be chosen to give the desired input resistance at the resonance frequency f_{CP} .



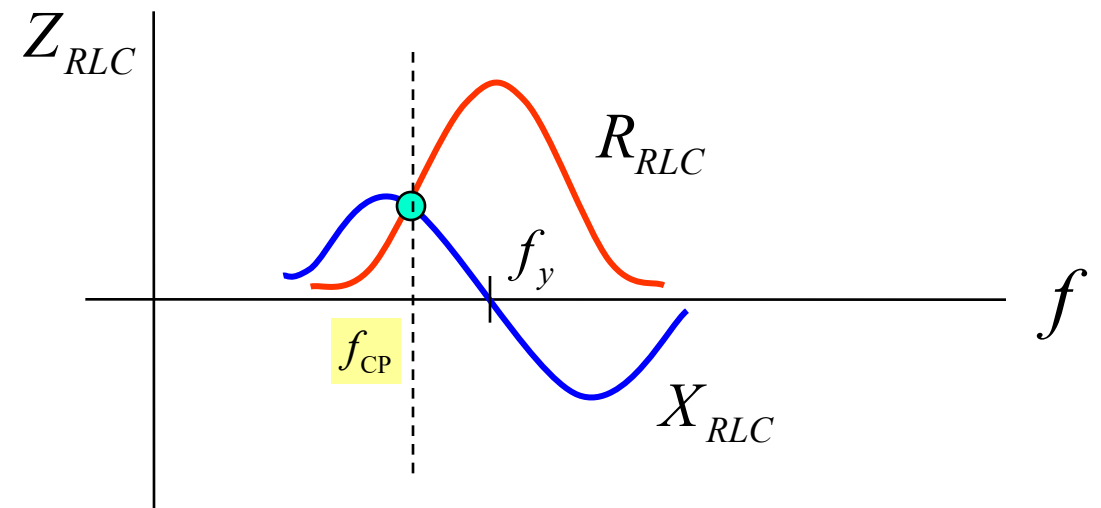
Input Impedance of CP Patch (cont.)

A physical picture of the frequencies (LHCP):



$$Z_{in}^x(f_{CP}) = \frac{R}{1+j}$$

$$X_{in}^x(f_{CP}) = -R_{in}^x(f_{CP})$$



$$Z_{in}^y(f_{CP}) = \frac{R}{1-j}$$

$$X_{in}^y(f_{CP}) = R_{in}^y(f_{CP})$$

CP (Axial Ratio) Bandwidth

We now examine the frequency dependence of the term A_y / A_x .

$$\begin{aligned} A_x &= \frac{A_2}{1 + j2Q(f_{rx} - 1)} \\ A_y &= \frac{A_2}{1 + j2Q(f_{ry} - 1)} \end{aligned} \quad \Rightarrow \quad \frac{A_y}{A_x} = \frac{1 + j2Q(f_{rx} - 1)}{1 + j2Q(f_{ry} - 1)}$$

where

$$f_{rx} \equiv \frac{f}{f_x} = \frac{f}{f_{CP} \left(1 - \frac{1}{2Q}\right)} = f_r \left(\frac{1}{1 - \frac{1}{2Q}} \right) \approx f_r \left(1 + \frac{1}{2Q} \right)$$

(LHCP)

$$f_r \equiv \frac{f}{f_{CP}}$$

This is the ratio of the operating frequency to the CP frequency.

CP Bandwidth (cont.)

Then

$$f_{rx} \approx f_r \left(1 + \frac{1}{2Q} \right)$$

$$f_{rx} = \frac{f}{f_x} = \frac{f}{f_{\text{CP}} \left(1 - \frac{1}{2Q} \right)} = f_r \left(\frac{1}{1 - \frac{1}{2Q}} \right) \approx f_r \left(1 + \frac{1}{2Q} \right)$$

Similarly,

$$f_{ry} \approx f_r \left(1 - \frac{1}{2Q} \right)$$

$$f_{ry} = \frac{f}{f_y} = \frac{f}{f_{\text{CP}} \left(1 + \frac{1}{2Q} \right)} = f_r \left(\frac{1}{1 + \frac{1}{2Q}} \right) \approx f_r \left(1 - \frac{1}{2Q} \right)$$

CP Bandwidth (cont.)

Hence, we have:

$$\frac{A_y}{A_x} = \frac{1 + j2Q(f_{rx} - 1)}{1 + j2Q(f_{ry} - 1)} = \frac{1 + j2Q\left(f_r\left(1 + \frac{1}{2Q}\right) - 1\right)}{1 + j2Q\left(f_r\left(1 - \frac{1}{2Q}\right) - 1\right)}$$

Note: $f_r = 1 \Rightarrow \frac{A_y}{A_x} = \frac{1+j}{1-j} = j$ (LHCP)

Let $x \equiv 2Q(f_r - 1)$ $f_r \equiv \frac{f}{f_{CP}}$

Then

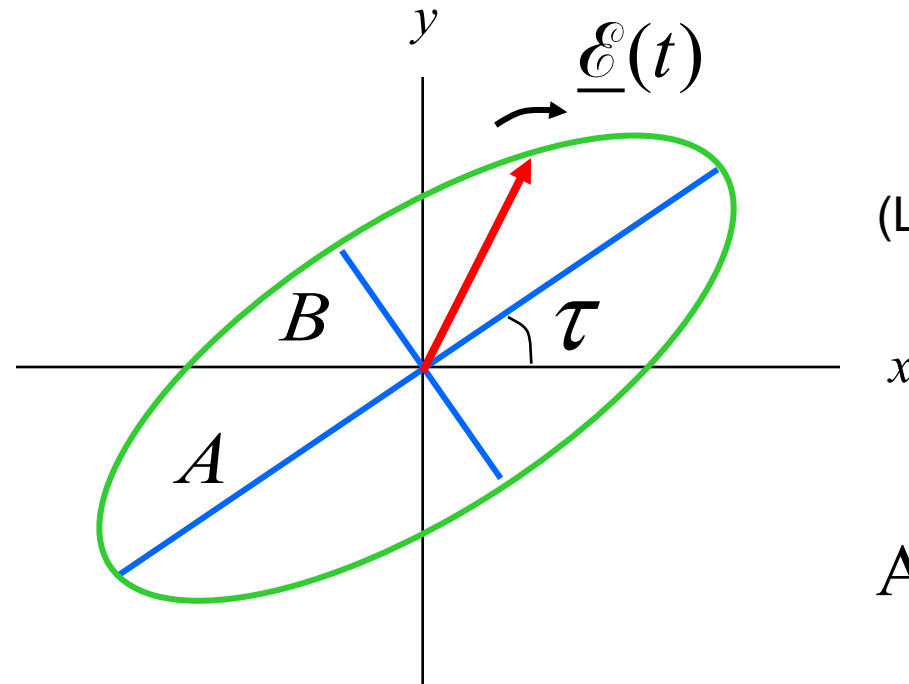
$$\frac{A_y}{A_x} = \frac{1 + j(f_r + x)}{1 - j(f_r - x)} \approx \frac{1 + j(1 + x)}{1 - j(1 - x)} \quad (f_r \approx 1)$$

CP Bandwidth (cont.)

$$\frac{A_y}{A_x} = \frac{1 + j(1+x)}{1 - j(1-x)}$$

$$x \equiv 2Q(f_r - 1)$$

$$f_r \equiv \frac{f}{f_{CP}}$$



$$AR \equiv \frac{A}{B}$$

CP Bandwidth (cont.)

From ECE 6340: $AR = \cot \xi$

where

$$\sin 2\xi = \sin(2\gamma) \sin \phi \quad \text{with} \quad -\frac{\pi}{2} \leq 2\xi \leq \frac{\pi}{2}$$

$$\gamma = \tan^{-1} \left| \frac{A_y}{A_x} \right|$$

$$\phi = \arg \left(\frac{A_y}{A_x} \right)$$

In our case:

$$\gamma = \tan^{-1} \sqrt{\frac{1 + (1+x)^2}{1 + (1-x)^2}}$$

$$\phi = \tan^{-1}(1+x) - \tan^{-1}(-(1-x)) = \tan^{-1}(1+x) + \tan^{-1}(1-x)$$

CP Bandwidth (cont.)

Set $AR = \sqrt{2}$ ($AR_{\text{dB}} = 20 \log_{10}(AR) = 3 \text{ dB}$)

From a numerical solution: $x = \pm 0.348$

CP Bandwidth (cont.)

Hence, we have: $2Q(f_r - 1) = \pm 0.348$

$$f_r = 1 \pm \frac{0.348}{2Q}$$

so

$$f_r^+ = 1 + \frac{0.348}{2Q}$$

$$f_r^- = 1 - \frac{0.348}{2Q}$$

Therefore, $\Delta f_r = f_r^+ - f_r^- = \frac{0.348}{Q}$

$$\text{BW}^{\text{AR CP}} \equiv \frac{f^+ - f^-}{f_{\text{CP}}} = f_r^+ - f_r^-$$

Hence, we have:

$$\text{BW}^{\text{AR CP}} = \frac{0.348}{Q}$$

Impedance Bandwidth

$$\begin{aligned} Z_{\text{in}}(f) &= \frac{R}{1 + j2Q(f_{rx} - 1)} + \frac{R}{1 + j2Q(f_{ry} - 1)} \\ &= \frac{R}{1 + j2Q\left(f_r\left(1 + \frac{1}{2Q}\right) - 1\right)} + \frac{R}{1 + j2Q\left(f_r\left(1 - \frac{1}{2Q}\right) - 1\right)} \quad (\text{LHCP}) \\ &= \frac{R}{1 + j(f_r + x)} + \frac{R}{1 - j(f_r - x)} \\ &\approx \frac{R}{1 + j(1 + x)} + \frac{R}{1 - j(1 - x)} \quad (\text{for } f_r \approx 1) \end{aligned}$$

where $x \equiv 2Q(f_r - 1)$

Note: At $x = 0$ we have $Z_{\text{in}} = \frac{R}{1 + j} + \frac{R}{1 - j} = R$

Impedance Bandwidth (cont.)

$$\bar{Z}_{\text{in}} \equiv \frac{Z_{\text{in}}}{R} = \frac{1}{1 + j(1+x)} + \frac{1}{1 - j(1-x)}$$

$$\Gamma = \frac{Z_{\text{in}} - Z_0}{Z_{\text{in}} + Z_0} = \frac{Z_{\text{in}} - R}{Z_{\text{in}} + R} = \frac{\bar{Z}_{\text{in}} - 1}{\bar{Z}_{\text{in}} + 1}$$

$$S = \text{SWR} = \frac{1 + |\Gamma|}{1 - |\Gamma|}$$

Set $S = S_0 = 2$ (bandwidth limits) $\Rightarrow |\Gamma| = \frac{1}{3}$ (-9.5 dB)

$$\Rightarrow \left| \frac{\frac{1}{1 + j(1+x)} + \frac{1}{1 - j(1-x)} - 1}{\frac{1}{1 + j(1+x)} + \frac{1}{1 - j(1-x)} + 1} \right| = \frac{1}{3} \Rightarrow x_0 = \pm\sqrt{2} \quad (\text{derivation omitted})$$

Impedance Bandwidth (cont.)

Hence $x \equiv 2Q(f_r - 1) = x_0 = \pm\sqrt{2}$

so $f_r = 1 \pm \frac{\sqrt{2}}{2Q}$

The band edges (in normalized frequency) are then

$$f_r^+ = 1 + \frac{1}{\sqrt{2}Q}$$

$$f_r^- = 1 - \frac{1}{\sqrt{2}Q}$$

Impedance Bandwidth (cont.)

$$\text{BW}^{\text{Imp CP}} \equiv \frac{f^+ - f^-}{f_{\text{CP}}} = f_r^+ - f_r^-$$

Hence

$$\text{BW}^{\text{Imp CP}} = 2 \left(\frac{1}{\sqrt{2}Q} \right)$$

Hence

$$\text{BW}^{\text{Imp CP}} = \frac{\sqrt{2}}{Q}$$

Summary

CP antenna

$$BW^{\text{AR CP}} = \frac{0.348}{Q}$$

$$BW^{\text{Imp CP}} = \frac{\sqrt{2}}{Q} = \frac{1.414}{Q}$$

Linear antenna

$$BW^{\text{Imp Lin}} = \frac{1}{\sqrt{2}Q} = \frac{0.707}{Q}$$